

New search strategies for well tempered neutralino dark matter

Bryan Ostdiek
University of Oregon

UC Davis High Energy Seminar
February 1, 2015

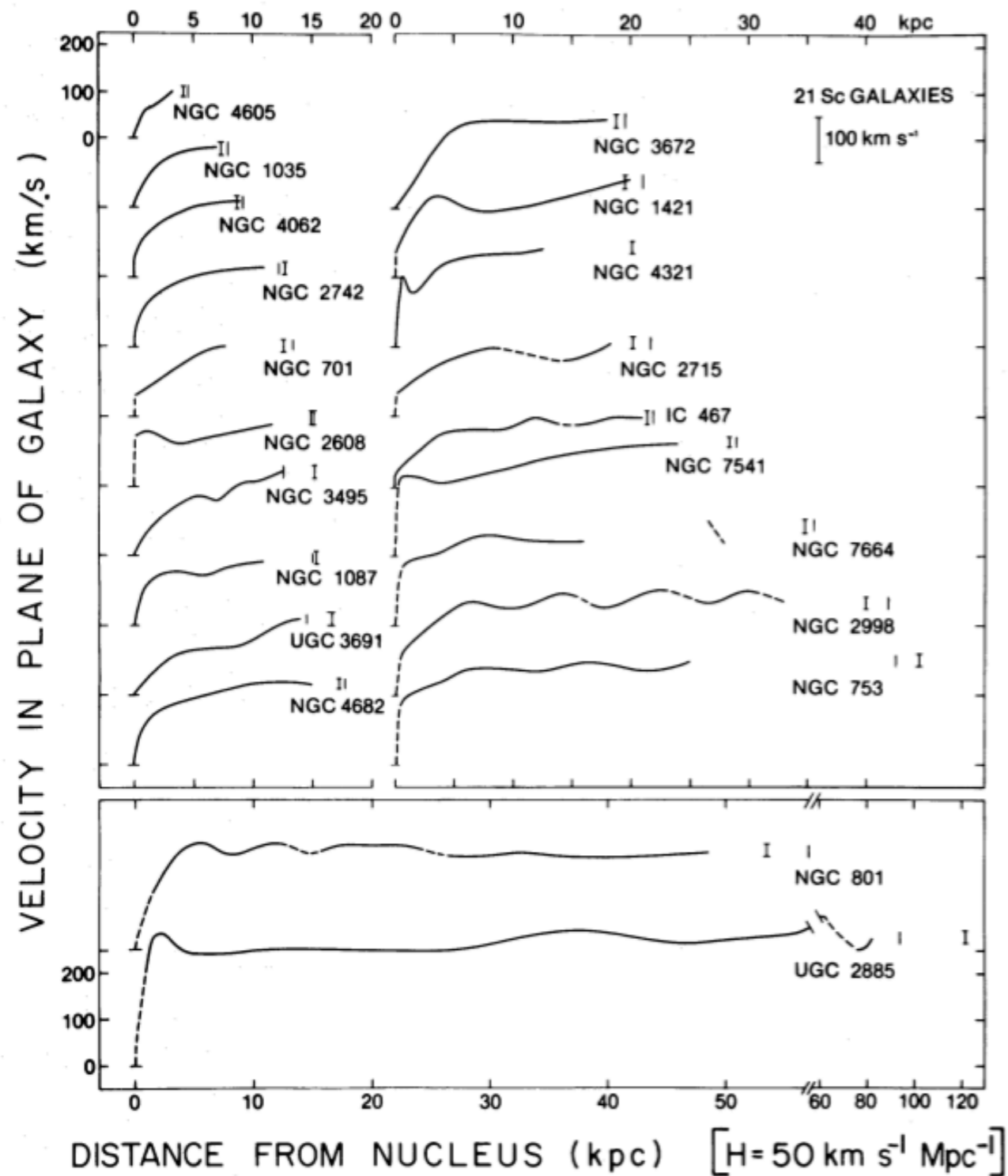
Roadmap of the talk

- Briefly review WIMP dark matter
- Explain what a well tempered neutralino is
- Astrophysical constraints from Indirect/Direct Detection
- Discuss why well tempered neutralinos are hard to find
- Show example of new strategies
- Parameter space covered with strategies

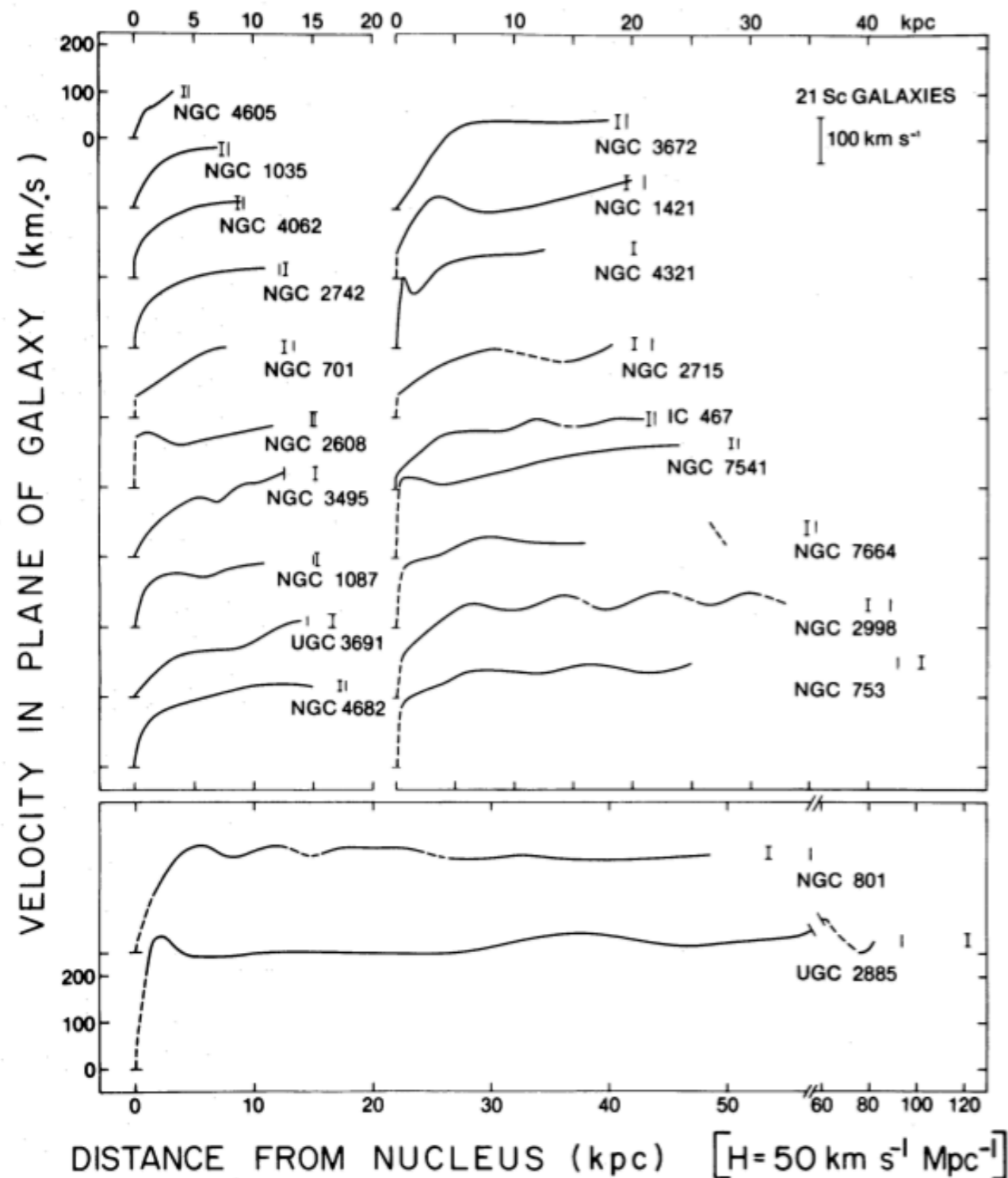
Dark Matter Overview

Evidence for DM

- Rotation curves
- Gravitational lensing
- Ia supernovae
- Cosmological nucleosynthesis
- CMB anisotropies



Dark Matter Overview



Evidence for DM

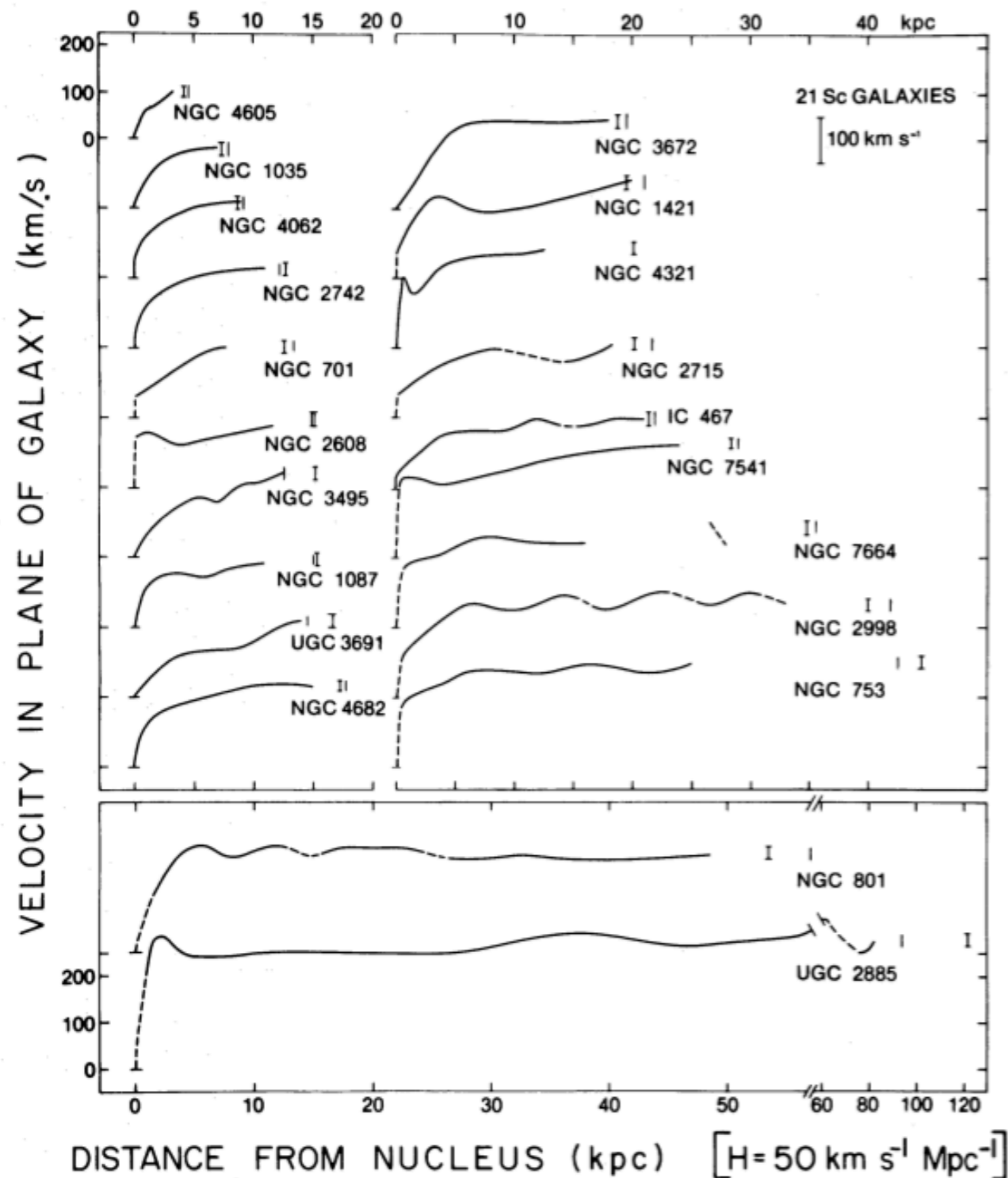
- Rotation curves
- Gravitational lensing
- Ia supernovae
- Cosmological nucleosynthesis
- CMB anisotropies

What do we Know?

- Density

$$\Omega_c = \rho \frac{8\pi G}{3H_0^2} = 0.2568$$
- Interacts with gravity, not photons

Dark Matter Overview



Evidence for DM

- Rotation curves
- Gravitational lensing
- Ia supernovae
- Cosmological nucleosynthesis
- CMB anisotropies

What do we Know?

- Density

$$\Omega_c = \rho \frac{8\pi G}{3H_0^2} = 0.2568$$
- Interacts with gravity, not photons

Still need to know:

- Mass, spin, interactions...

WIMP Dark Matter

Assume that DM does interact with the SM by means beyond gravity

WIMP Dark Matter

Assume that DM does interact with the SM by means beyond gravity

- DM can be in thermal equilibrium with SM in the early universe
- Expansion leads to freeze-out, “WIMP miracle”

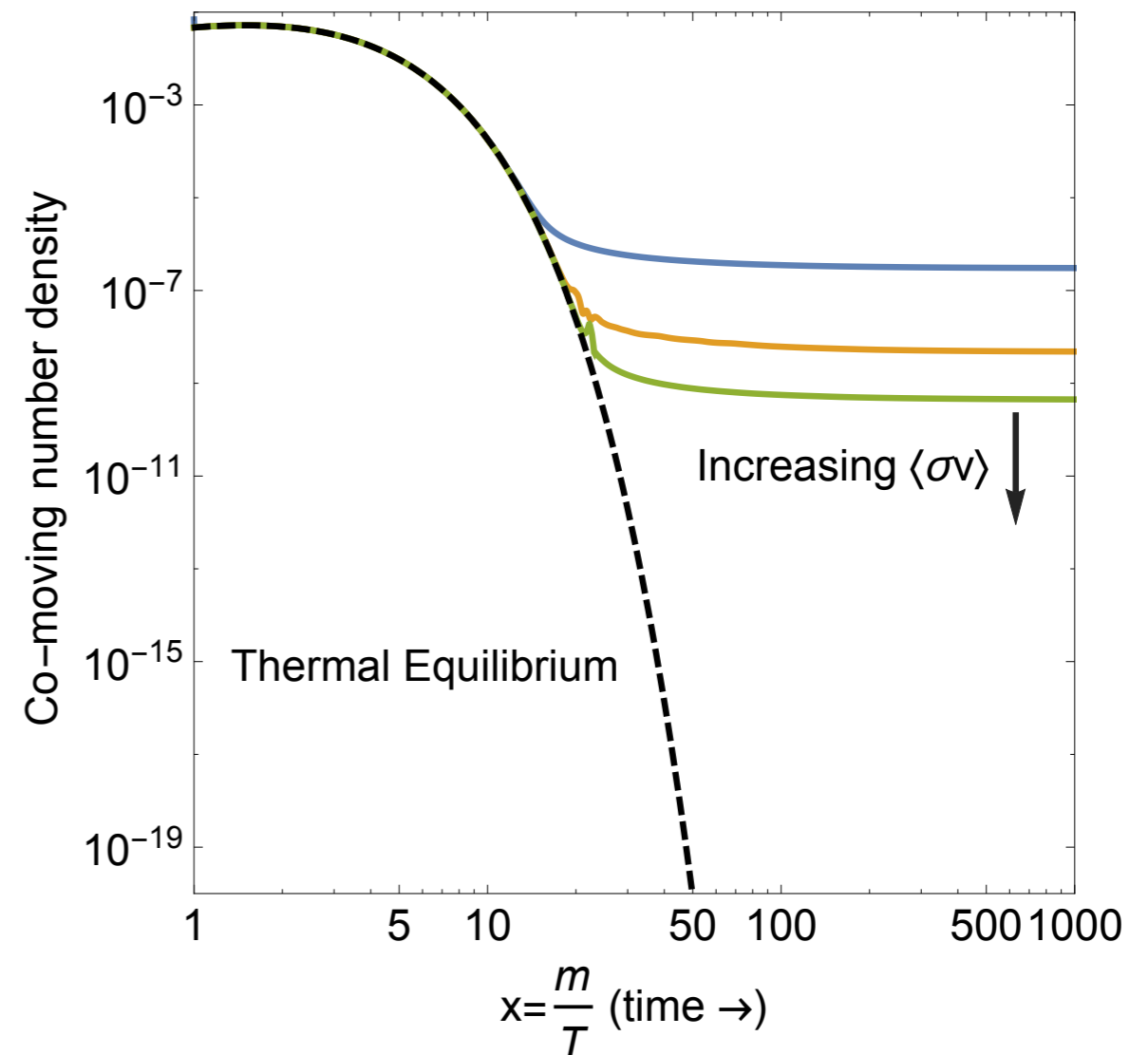
$$\frac{d(n_X a^3)}{dt} = -(n_X^2 - n_{X,eq}^2) a^3 \langle \sigma v \rangle$$

WIMP Dark Matter

Assume that DM does interact with the SM by means beyond gravity

- DM can be in thermal equilibrium with SM in the early universe
- Expansion leads to freeze-out, “WIMP miracle”

$$\frac{d(n_X a^3)}{dt} = -(n_X^2 - n_{X,eq}^2) a^3 \langle \sigma v \rangle$$



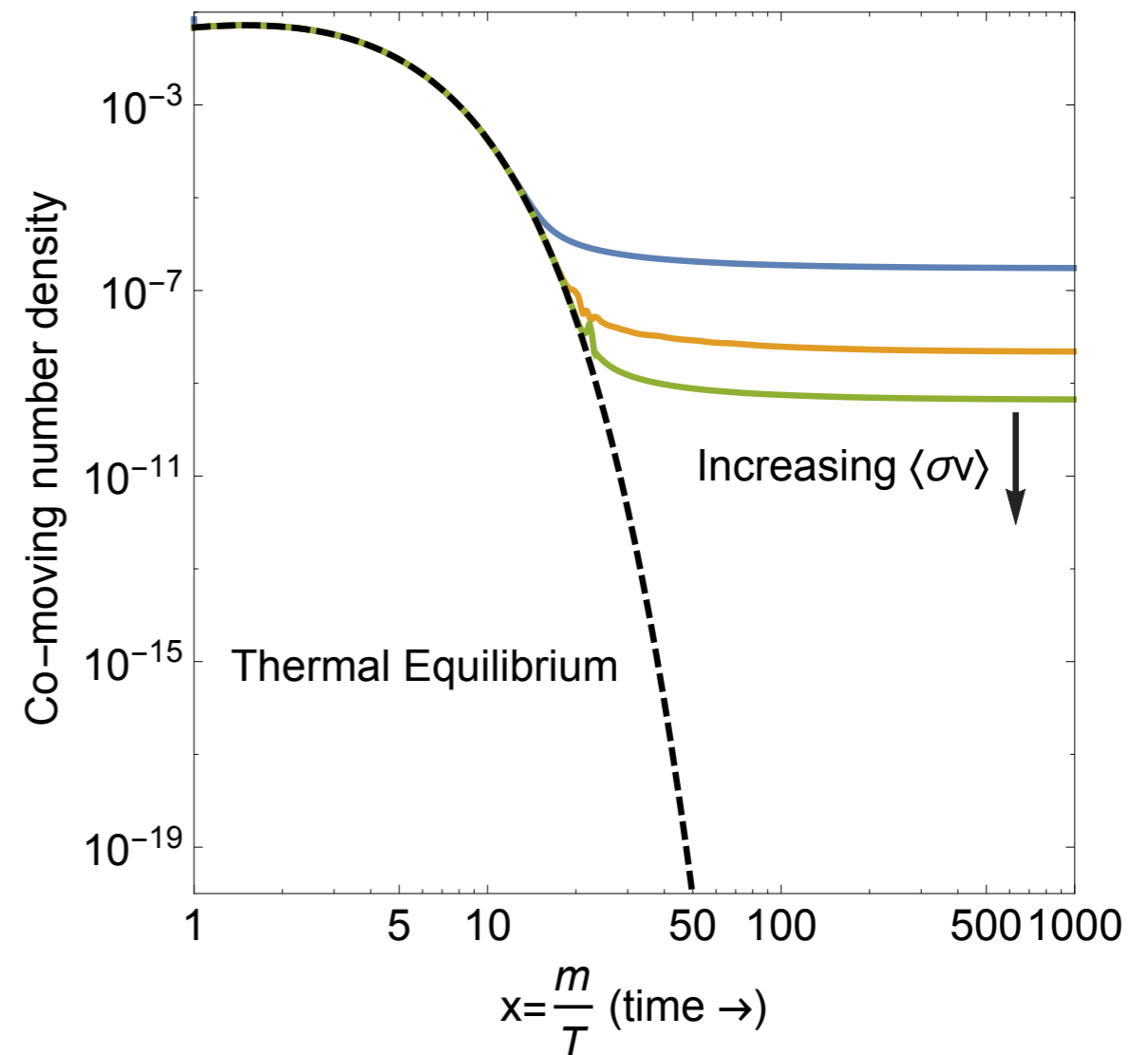
WIMP Dark Matter

Assume that DM does interact with the SM by means beyond gravity

- DM can be in thermal equilibrium with SM in the early universe
- Expansion leads to freeze-out, “WIMP miracle”

$$\frac{d(n_X a^3)}{dt} = -(n_X^2 - n_{X,eq}^2) a^3 \langle \sigma v \rangle$$

- Allows for search methods
 - Direct Detection
 - Indirect Detection
 - Production at colliders



WIMP Dark Matter

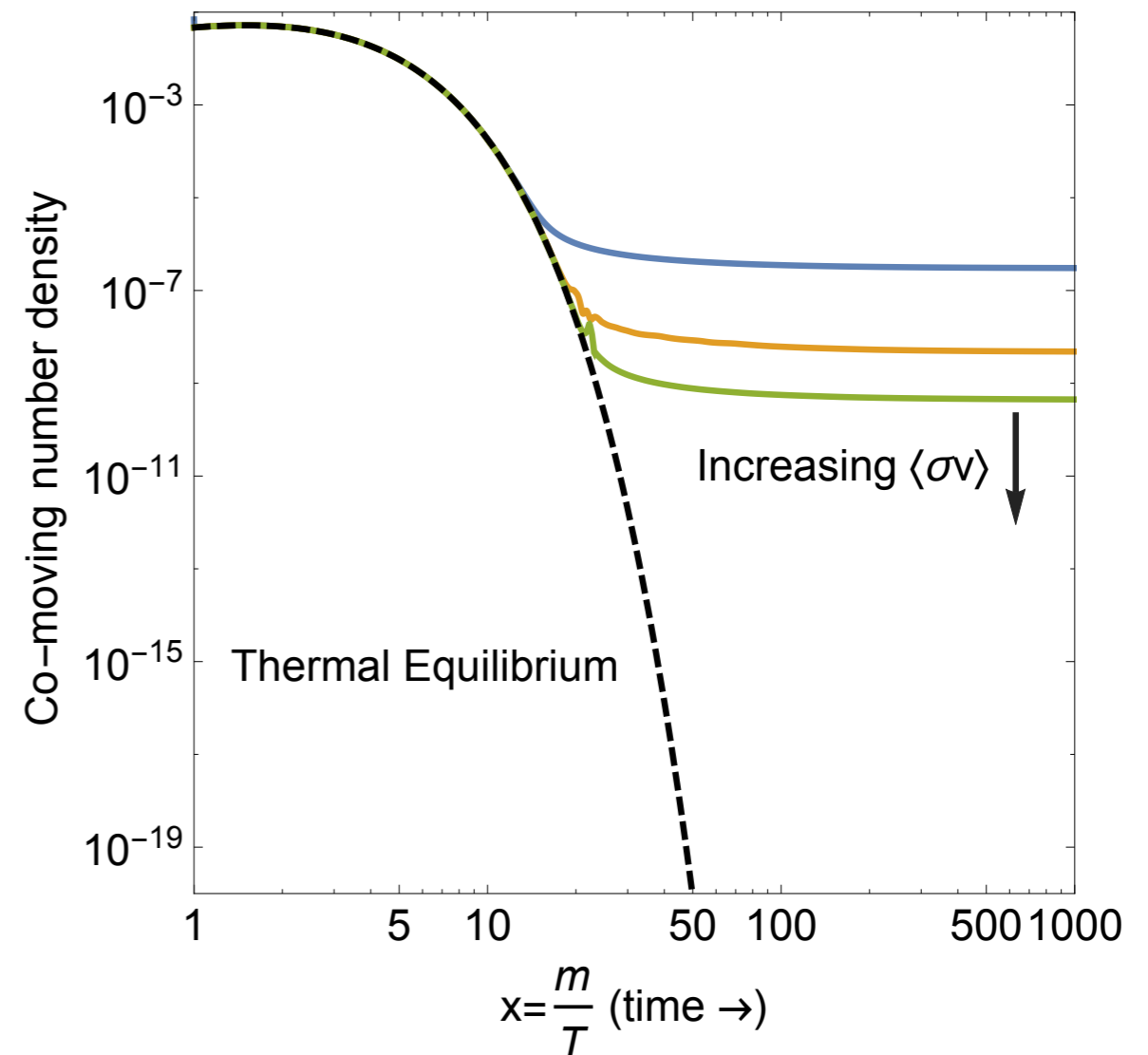
Assume that DM does interact with the SM by means beyond gravity

- DM can be in thermal equilibrium with SM in the early universe
- Expansion leads to freeze-out, “WIMP miracle”

$$\frac{d(n_X a^3)}{dt} = -(n_X^2 - n_{X,eq}^2) a^3 \langle \sigma v \rangle$$

- Allows for search methods
 - Direct Detection
 - Indirect Detection
- Production at colliders

Use supersymmetry as a WIMP model. Extra particles affect $\langle \sigma v \rangle$



Dark matter from supersymmetry

Spin 0	Spin 1/2	Spin 1
h		
	q, l	
		V_μ

Dark matter from supersymmetry

Spin 0	Spin 1/2	Spin 1
h	\tilde{h}	
	q, l	
		V_μ

Dark matter from supersymmetry

Spin 0	Spin 1/2	Spin 1
h	\tilde{h}	
$\tilde{q}, \tilde{\ell}$	q, l	
		V_μ

Dark matter from supersymmetry

Spin 0	Spin 1/2	Spin 1
h	\tilde{h}	
$\tilde{q}, \tilde{\ell}$	q, l	
	\tilde{V}_μ	V_μ

Dark matter from supersymmetry

Spin 0	Spin 1/2	Spin 1
h	\tilde{h}	
$\tilde{q}, \tilde{\ell}$	q, l	
	\tilde{V}_μ	V_μ

Higgsinos and Gauginos mix to form Neutralinos and Charginos

Dark matter from supersymmetry

Spin 0	Spin 1/2	Spin 1
h	\tilde{h}	
$\tilde{q}, \tilde{\ell}$	q, l	
	\tilde{V}_μ	V_μ

Higgsinos and Gauginos mix to form Neutralinos and Charginos

$$\begin{pmatrix} \tilde{B} & \tilde{W}^3 & \tilde{H}_d^0 & \tilde{H}_u^0 \end{pmatrix} \begin{pmatrix} M_1 & 0 & -c_\beta s_W m_Z & s_\beta s_W m_Z \\ 0 & M_2 & c_\beta c_W m_Z & -s_\beta c_W m_Z \\ -c_\beta s_W m_Z & c_\beta c_W m_Z & 0 & -\mu \\ s_\beta s_W m_Z & -s_\beta c_W m_Z & -\mu & 0 \end{pmatrix} \begin{pmatrix} \tilde{B} \\ \tilde{W}^3 \\ \tilde{H}_d^0 \\ \tilde{H}_u^0 \end{pmatrix}$$

Dark matter from supersymmetry

Spin 0	Spin 1/2	Spin 1
h	\tilde{h}	
$\tilde{q}, \tilde{\ell}$	q, l	
	\tilde{V}_μ	V_μ

Higgsinos and Gauginos mix to form Neutralinos and Charginos

$$\begin{pmatrix} \tilde{B} & \tilde{W}^3 & \tilde{H}_d^0 & \tilde{H}_u^0 \end{pmatrix} \begin{pmatrix} M_1 & 0 & -c_\beta s_W m_Z & s_\beta s_W m_Z \\ 0 & M_2 & c_\beta c_W m_Z & -s_\beta c_W m_Z \\ -c_\beta s_W m_Z & c_\beta c_W m_Z & 0 & -\mu \\ s_\beta s_W m_Z & -s_\beta c_W m_Z & -\mu & 0 \end{pmatrix} \begin{pmatrix} \tilde{B} \\ \tilde{W}^3 \\ \tilde{H}_d^0 \\ \tilde{H}_u^0 \end{pmatrix}$$

$\langle \sigma v \rangle$ **determined by** M_1, M_2, μ , and $\tan \beta$

R parity doesn't allow LSP to decay

Dark matter from supersymmetry

Spin 0	Spin 1/2	Spin 1
h	\tilde{h}	
$\tilde{q}, \tilde{\ell}$	q, l	
	\tilde{V}_μ	V_μ

Higgsinos and Gauginos mix to form Neutralinos and Charginos

$$\begin{pmatrix} \tilde{B} & \tilde{W}^3 & \tilde{H}_d^0 & \tilde{H}_u^0 \end{pmatrix} \begin{pmatrix} M_1 & 0 & -c_\beta s_W m_Z & s_\beta s_W m_Z \\ 0 & M_2 & c_\beta c_W m_Z & -s_\beta c_W m_Z \\ -c_\beta s_W m_Z & c_\beta c_W m_Z & 0 & -\mu \\ s_\beta s_W m_Z & -s_\beta c_W m_Z & -\mu & 0 \end{pmatrix} \begin{pmatrix} \tilde{B} \\ \tilde{W}^3 \\ \tilde{H}_d^0 \\ \tilde{H}_u^0 \end{pmatrix}$$

$\langle \sigma v \rangle$ **determined by** M_1, M_2, μ , and $\tan \beta$

R parity doesn't allow LSP to decay

Isolated EW-inos

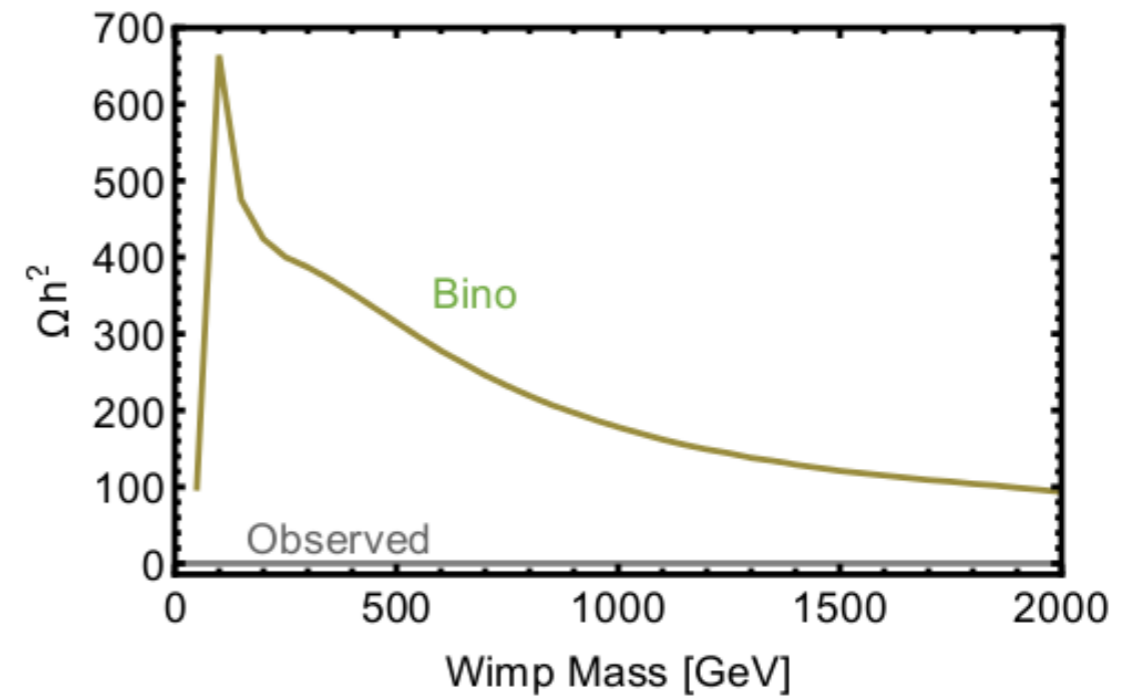
What mass is needed to quench relic abundance?

Isolated EW-inos

What mass is needed to quench relic abundance?

Bino

- Gauge singlet
- 1 Neutralino, 0 Charginos



Isolated EW-inos

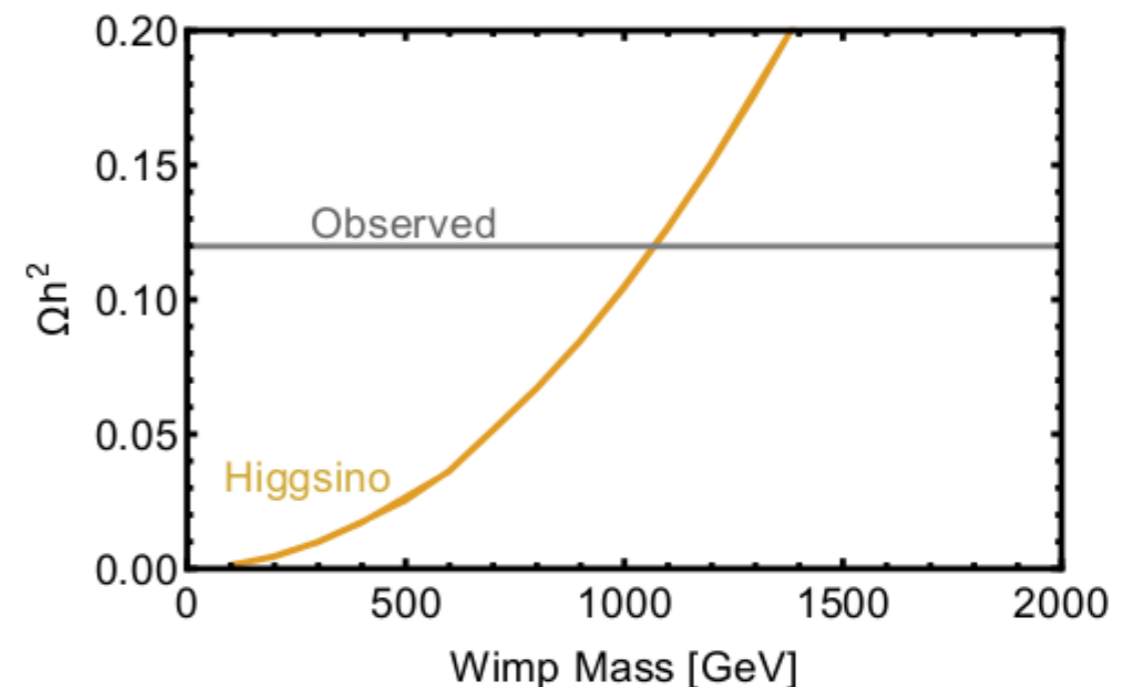
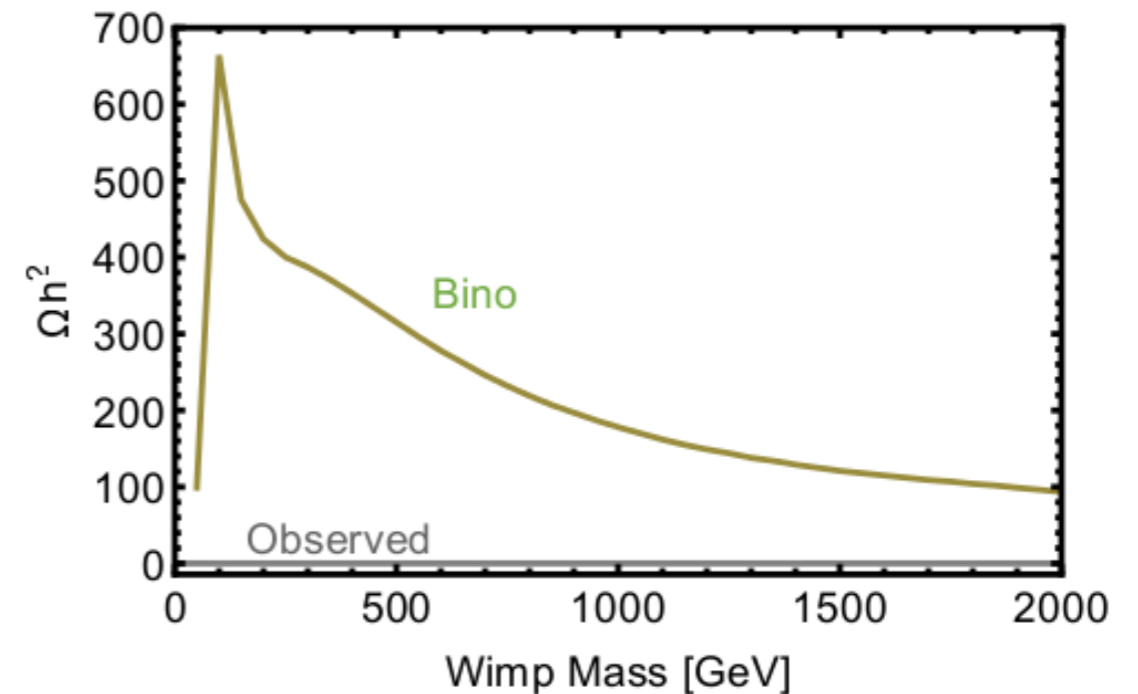
What mass is needed to quench relic abundance?

Bino

- Gauge singlet
- 1 Neutralino, 0 Charginos

Higgsinos

- 2 Gauge doublets
- 2 Neutralinos, 1 Charginos



Isolated EW-inos

What mass is needed to quench relic abundance?

Bino

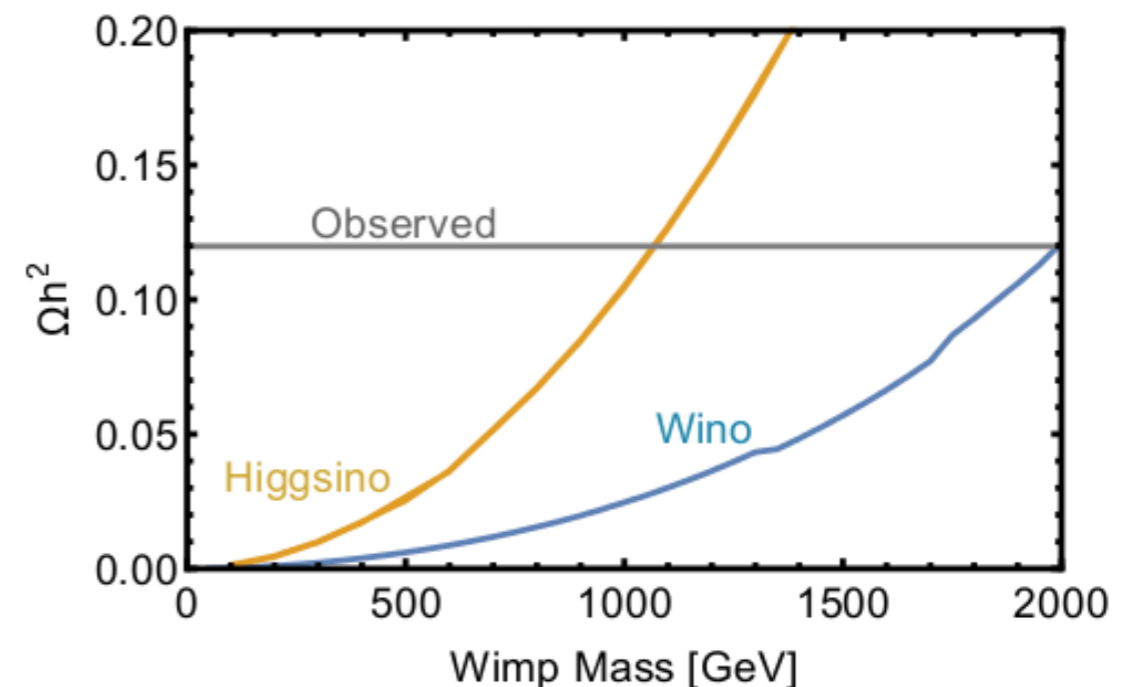
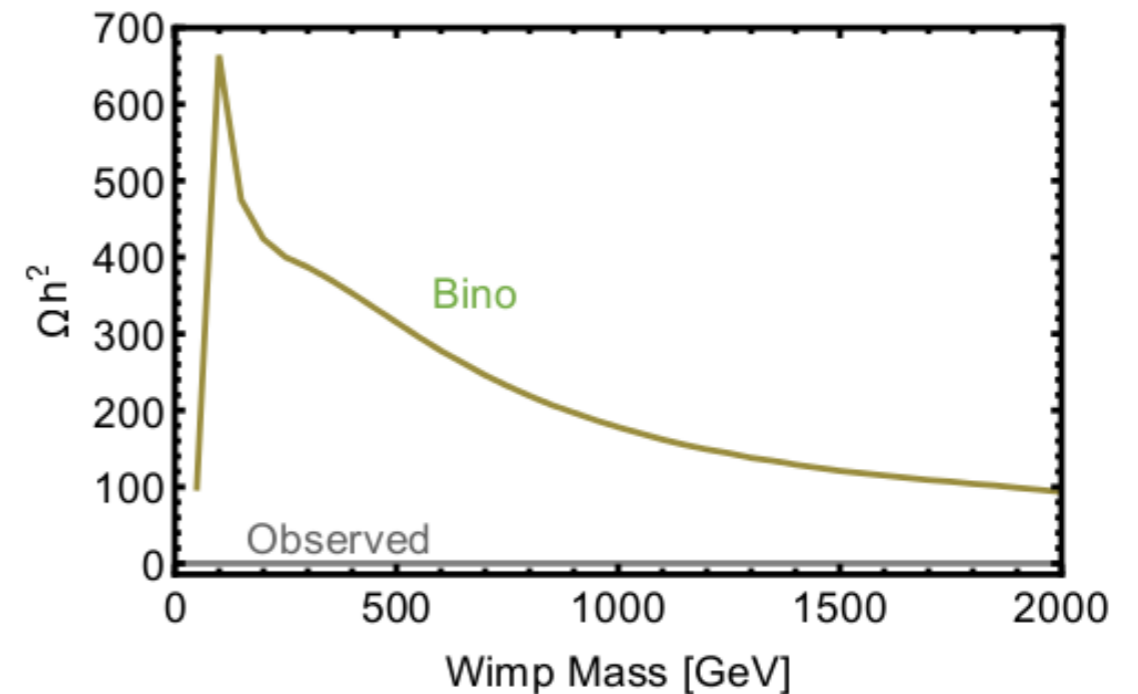
- Gauge singlet
- 1 Neutralino, 0 Charginos

Higgsinos

- 2 Gauge doublets
- 2 Neutralinos, 1 Charginos

Wino

- 1 Gauge triplet
- 1 Neutralino, 1 Charginos

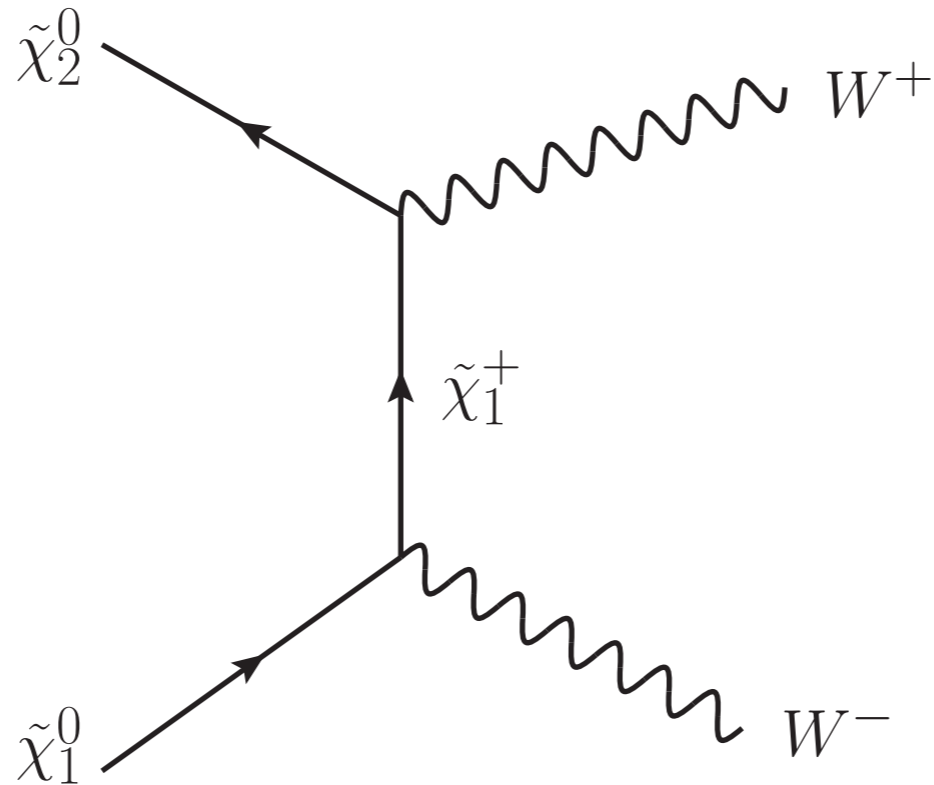


Co-annihilations

Why are the relic abundances different for the pure electroweakinos?

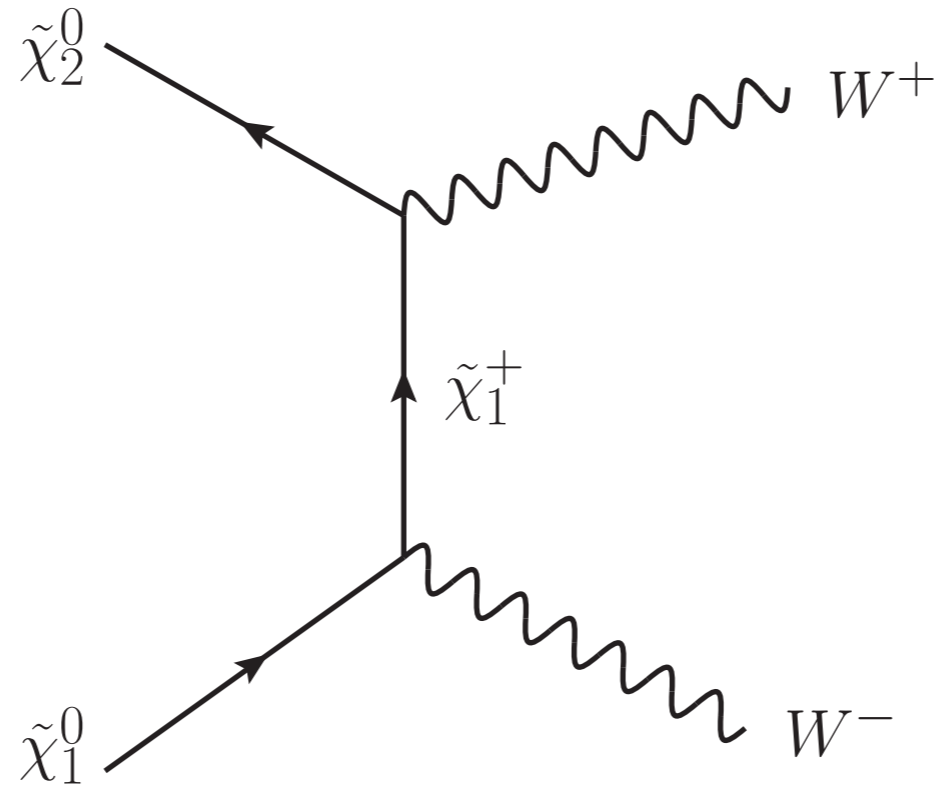
Co-annihilations

Why are the relic abundances different for the pure electroweakinos?



Co-annihilations

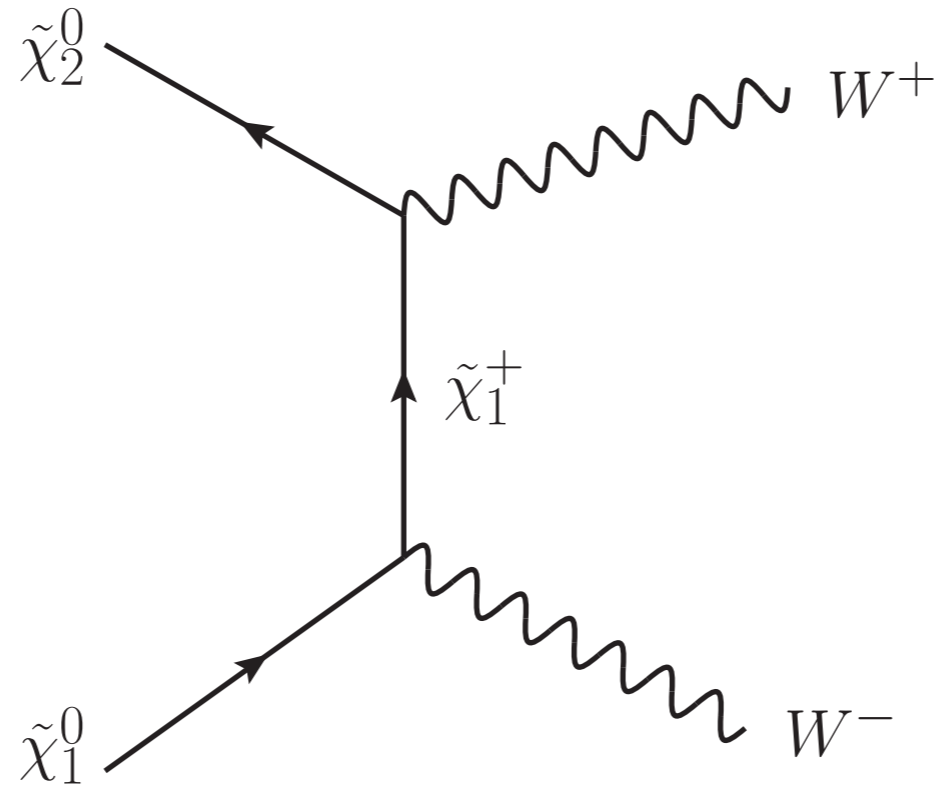
Why are the relic abundances different for the pure electroweakinos?



Annihilation cross section affected by particles near in mass

Co-annihilations

Why are the relic abundances different for the pure electroweakinos?



Annihilation cross section affected by particles near in mass

Well Tempering tunes the values of M_1 , M_2 , μ , and $\tan \beta$ to achieve observed relic abundance

The Well Tempered Surface

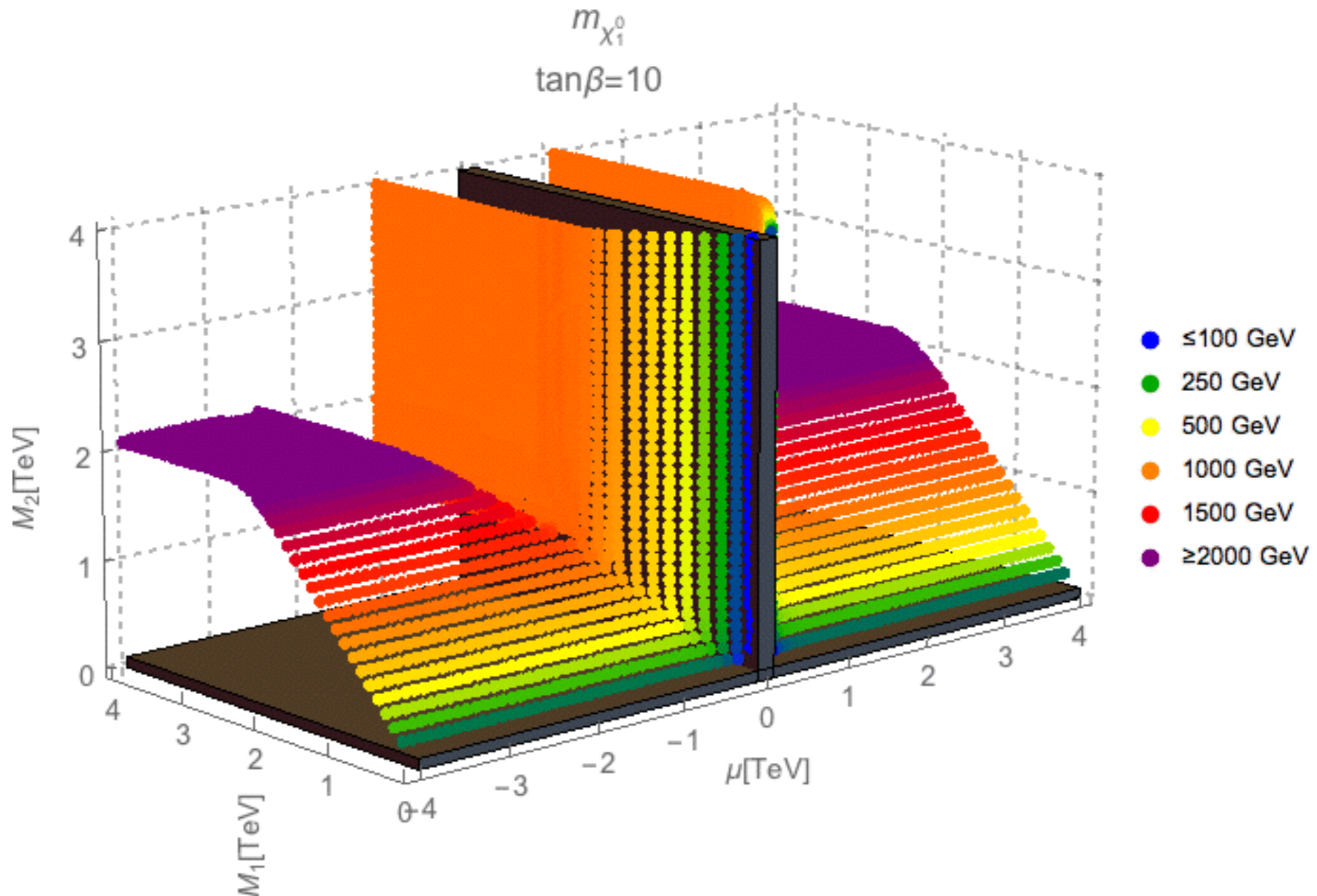
Well Tempering tunes the values of M_1 , M_2 , μ , and $\tan \beta$ to achieve observed relic abundance

- Decouple all supersymmetric scalars (heavy Higgs and sfermions)
- Chose a value for $\tan \beta$ (10)
- Scan over values of M_1 , M_2 , and μ
(Spectrum calculated with SUSPECT)
- Keep model point if $\Omega h^2 = 0.12$
(DM properties calculated with micrOMEGAs)
- Points left over define the **Well Tempered Surface**

Next set of plots were first presented in J. Bramante, P. J. Fox, A. Martin, BO, T. Plehn, T. Schell and M. Takeuchi, “Relic neutralino surface at a 100 TeV collider,” Phys. Rev. D 91, no. 5, 054015 (2015) [arXiv:1412.4789 [hep-ph]]

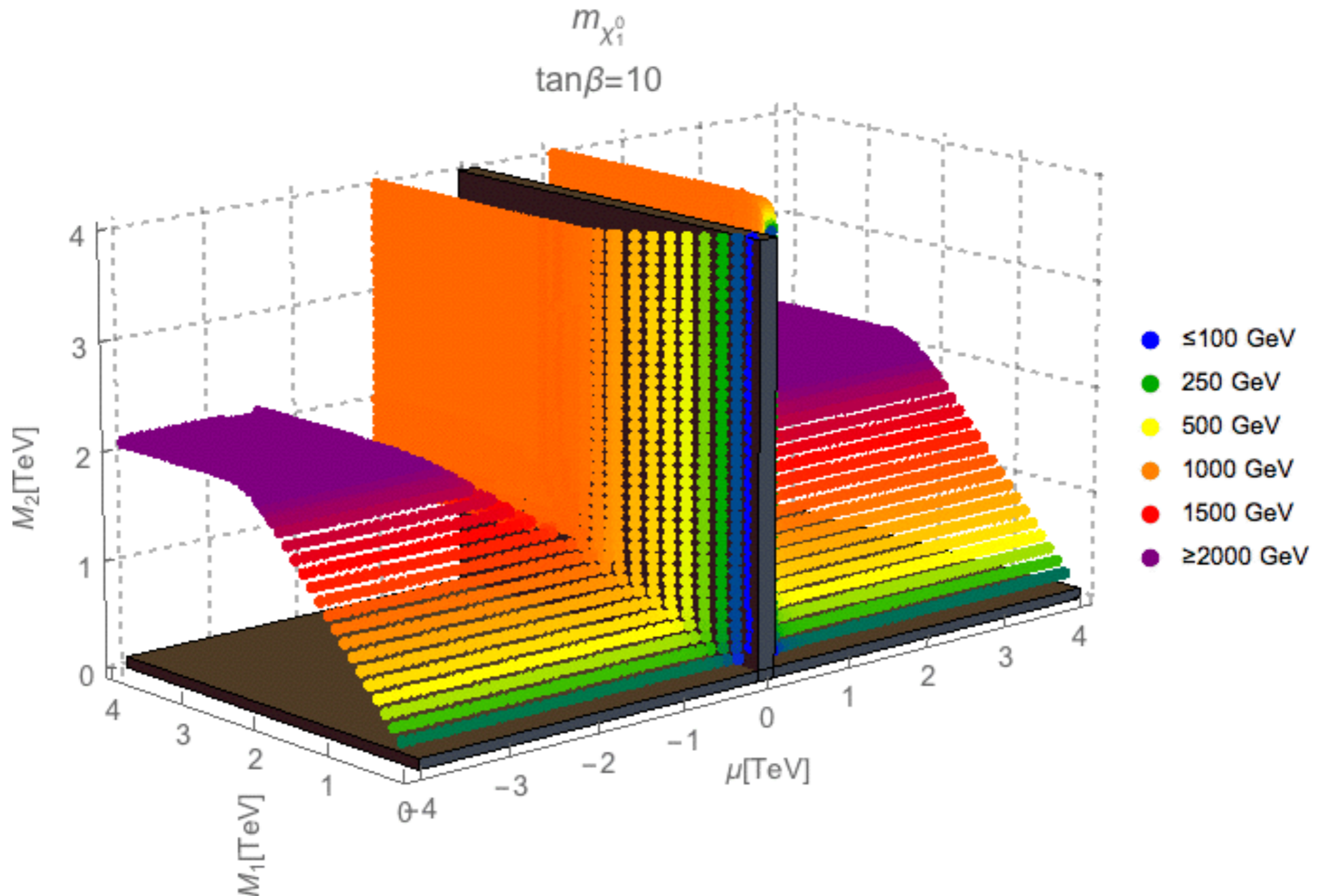
The Well Tempered Surface

Mass of the lightest neutralino



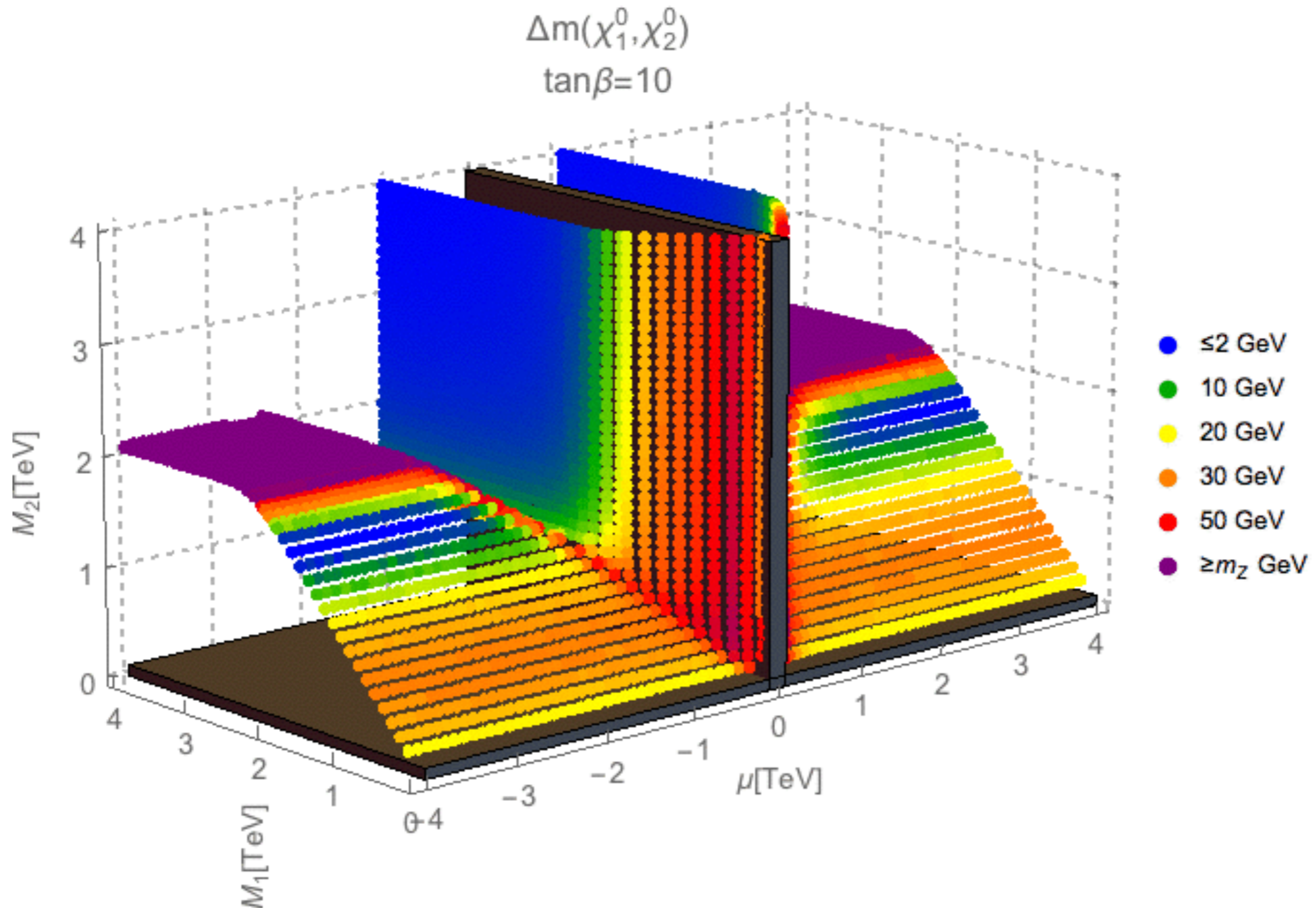
The Well Tempered Surface

Mass of the lightest neutralino



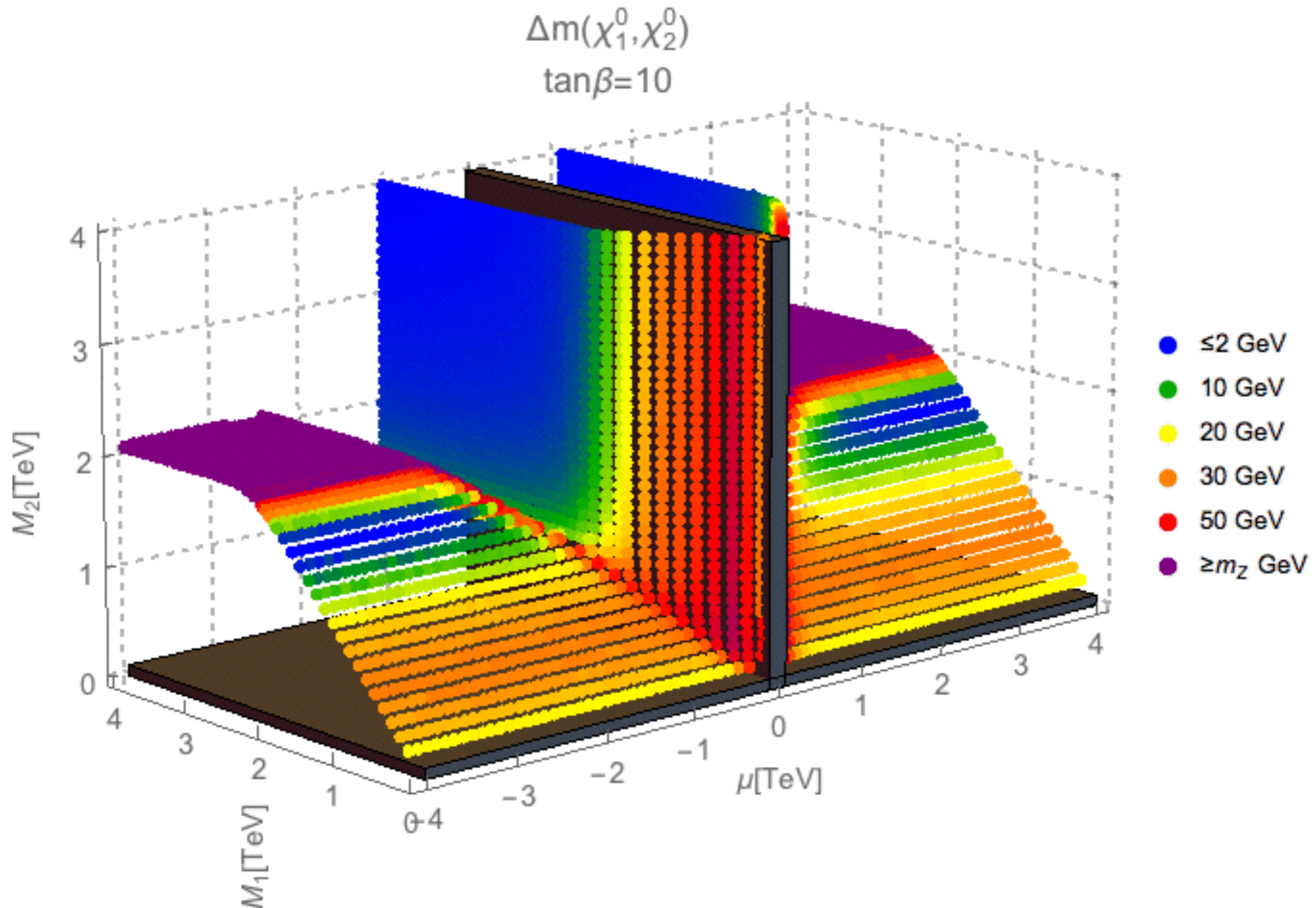
The Well Tempered Surface

Mass difference between the lightest two neutralinos



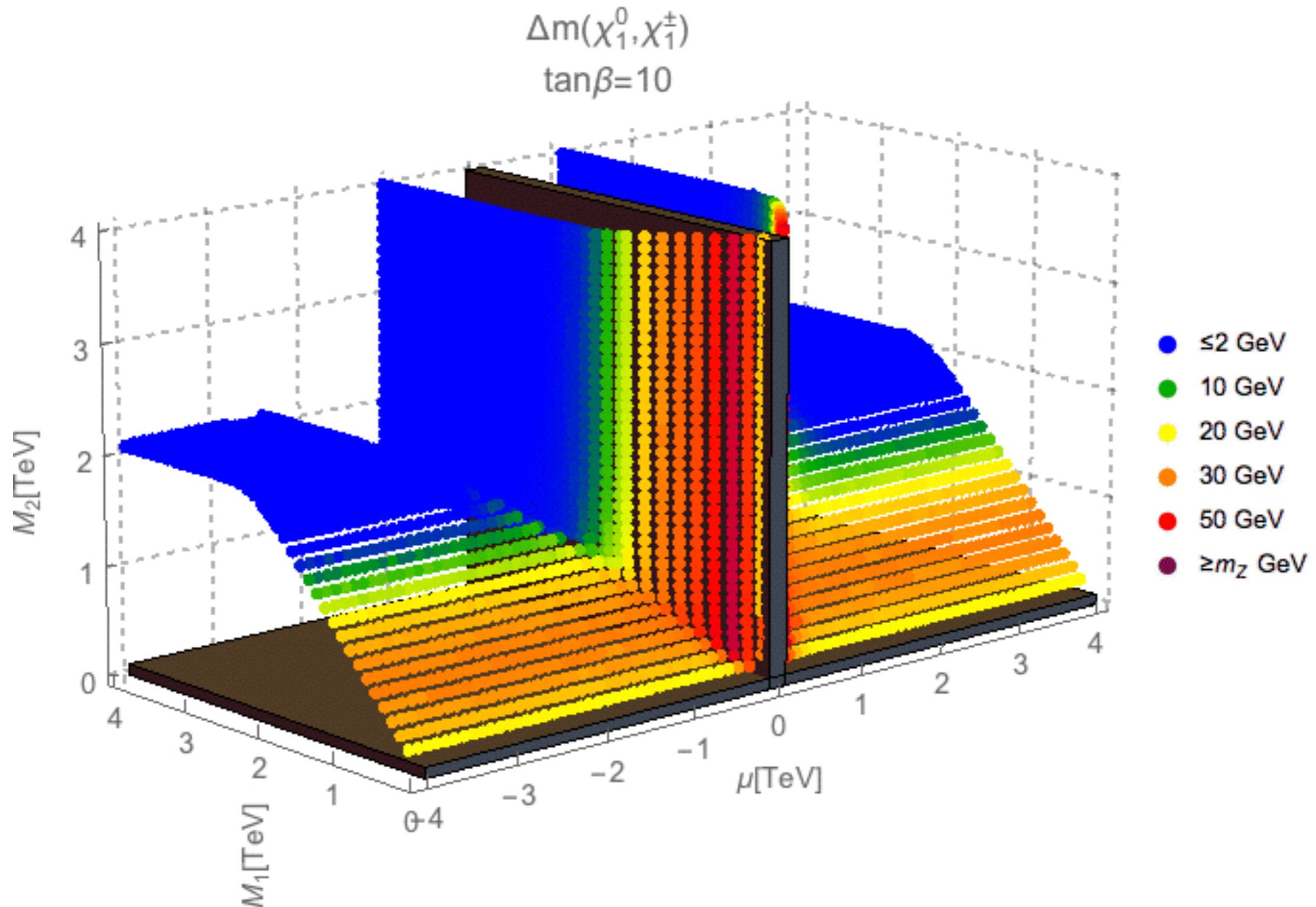
The Well Tempered Surface

Mass difference between the lightest two neutralinos



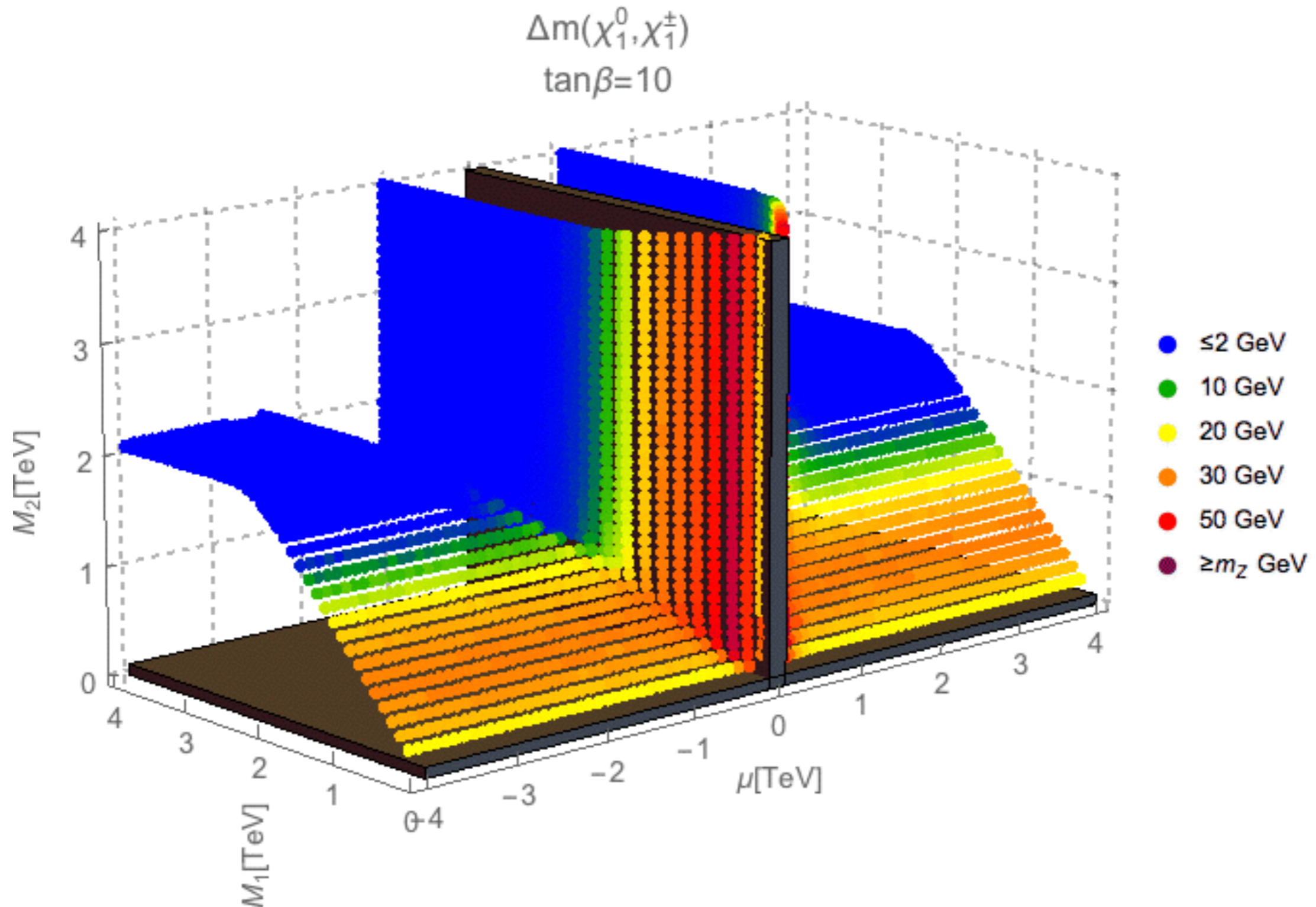
The Well Tempered Surface

Mass difference between the lightest neutralino and chargino



The Well Tempered Surface

Mass difference between the lightest neutralino and chargino



What have we learned?

- SUSY with R-parity provides a DM candidate
- Not all SUSY DM candidates give correct DM abundance
- Well tempering leads to small mass splittings

What have we learned?

- SUSY with R-parity provides a DM candidate
- Not all SUSY DM candidates give correct DM abundance
- Well tempering leads to small mass splittings

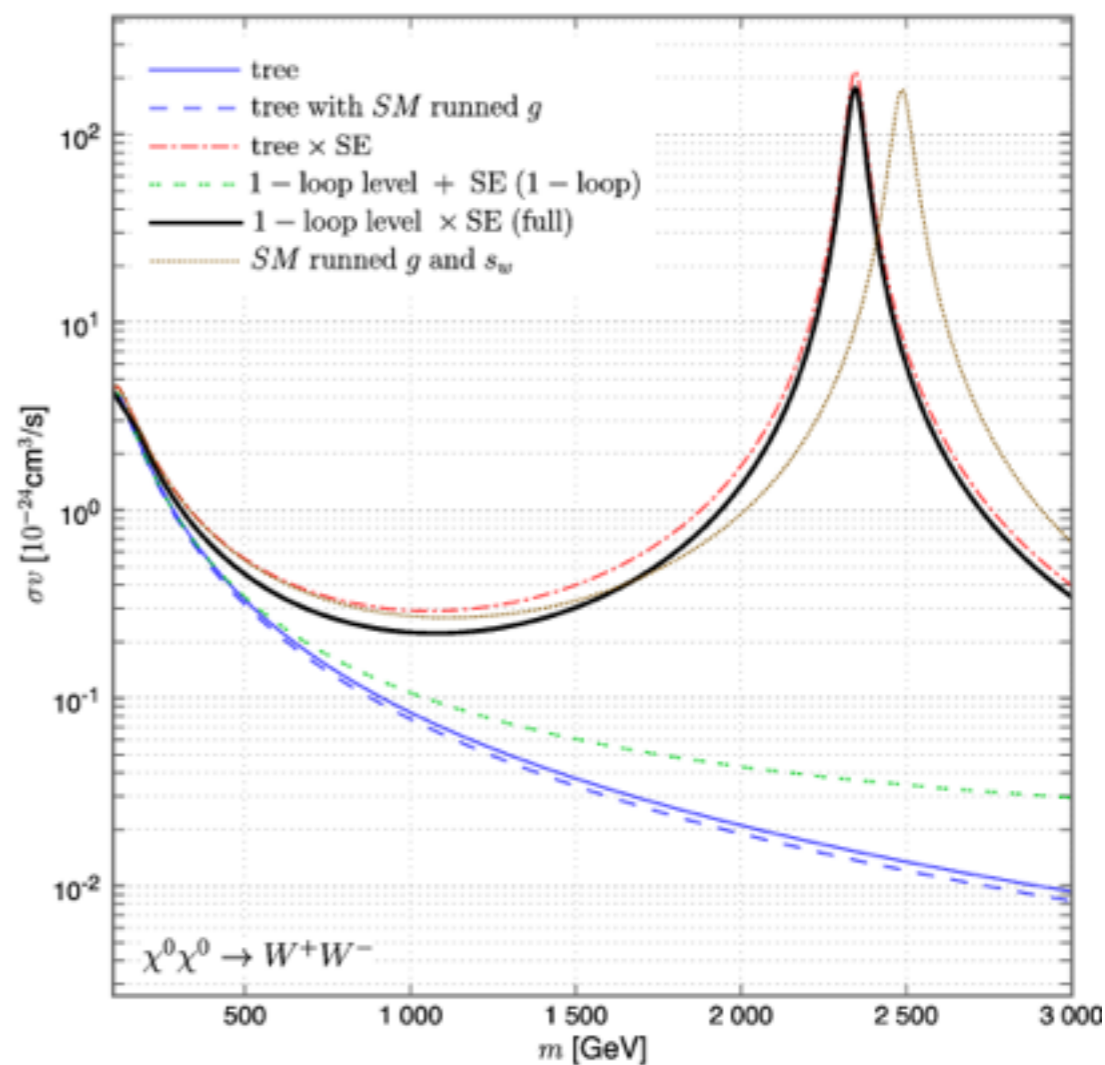
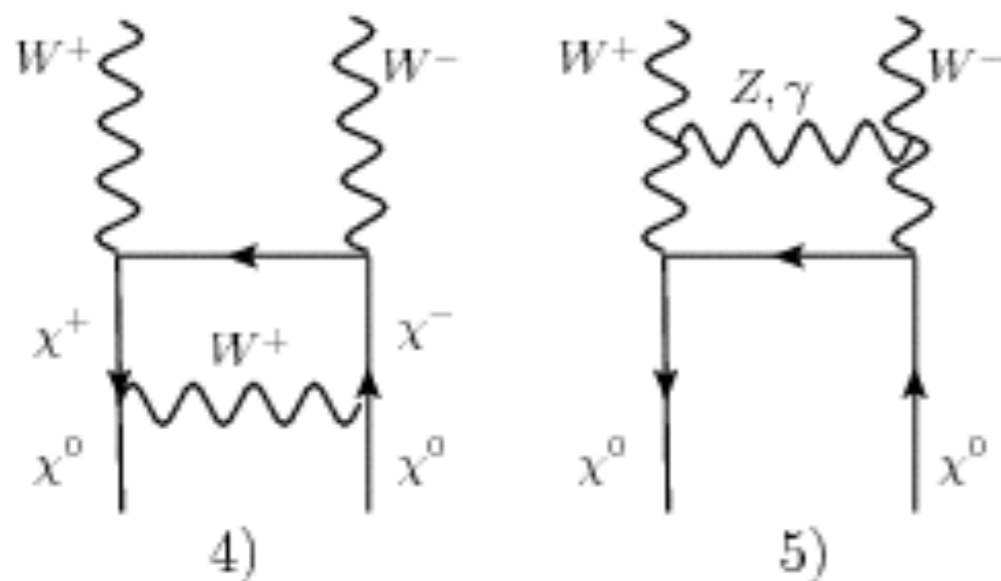
Modifications and improvements

- Sommerfeld enhancement substantially increases pure Wino annihilation cross section
- Some effect for pure Higgsino
- How is the surface affected?
 - Use DarkSE code by Hruczuk

J. Bramante, N. Desai, P. Fox, A. Martin, B. Ostdiek and T. Plehn, "Towards the Final Word on Neutralino Dark Matter," arXiv:1510.03460 [hep-ph]

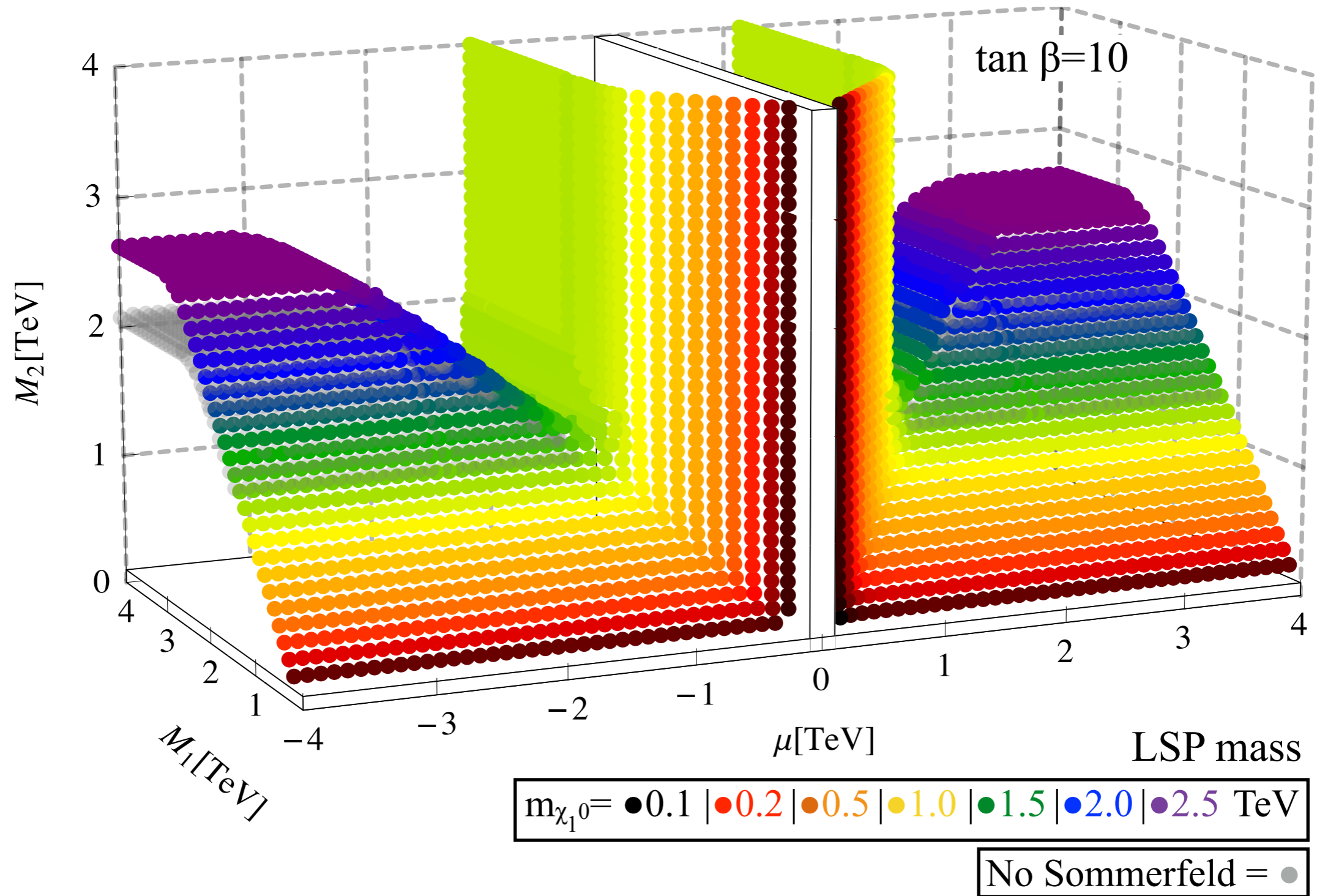
What is the Sommerfeld enhancement?

“The modification of the wave function of the incoming non-relativistic particles due to their mutual interaction”



A. Hryczuk and R. Iengo, “The one-loop and Sommerfeld electroweak corrections to the Wino dark matter annihilation,” *JHEP* **1201**, 163 (2012) [*JHEP* **1206**, 137 (2012)] doi:10.1007/JHEP01(2012)163, 10.1007/JHEP06(2012)137 [arXiv:1111.2916 [hep-ph]]

Sommerfelded Surface



Can entire surface be discovered with current/future experiments?

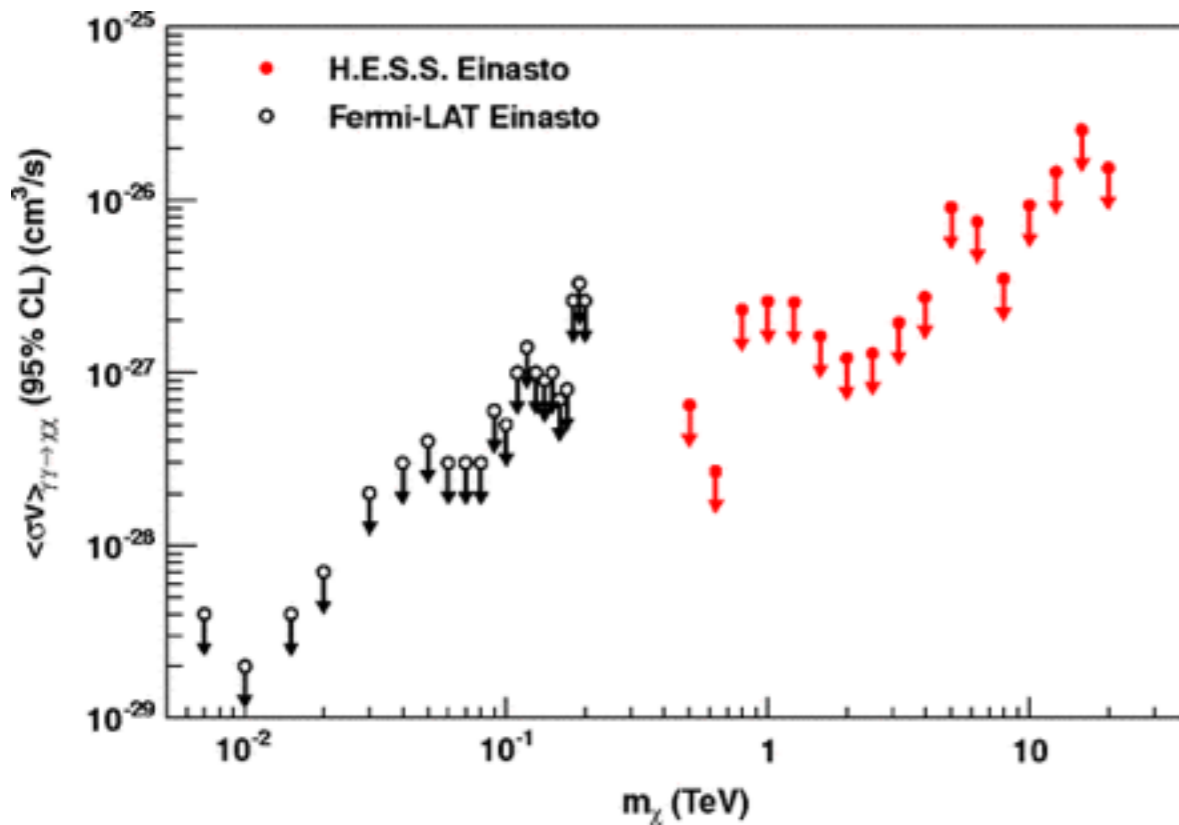
- Examine indirect/direct detection constraints
- Explore search techniques at high energy colliders
- Complementarity between experiments

Can entire surface be discovered with current/future experiments?

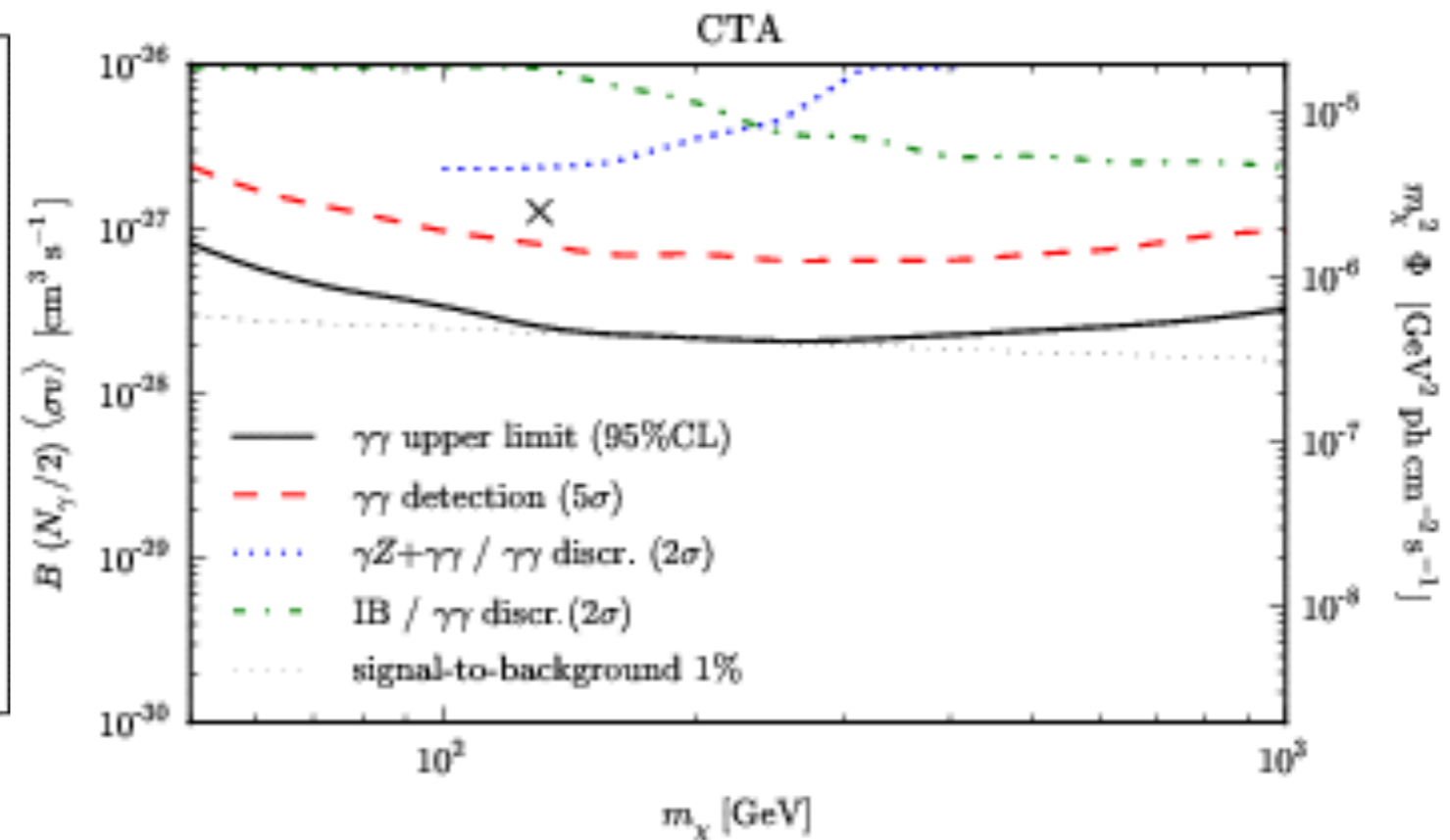
- Examine indirect/direct detection constraints
- Explore search techniques at high energy colliders
- Complementarity between experiments

Indirect Detection

- Annihilations still happen in dense regions
- Can do $\chi\chi \rightarrow \gamma\gamma$ ($E_\gamma = m_\chi$)
- Lack of signal leads to constraints



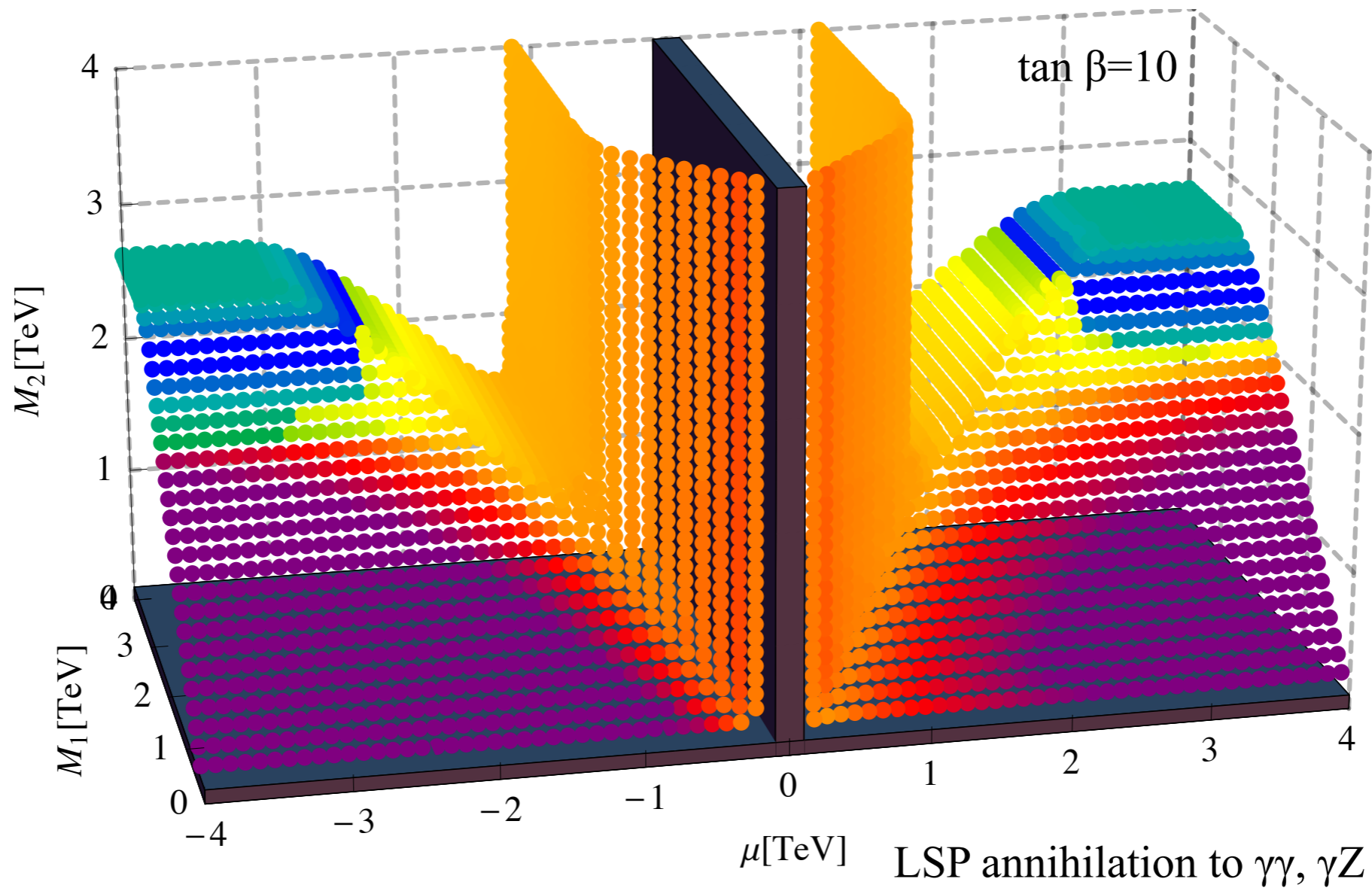
Current



Projection

Indirect Detection

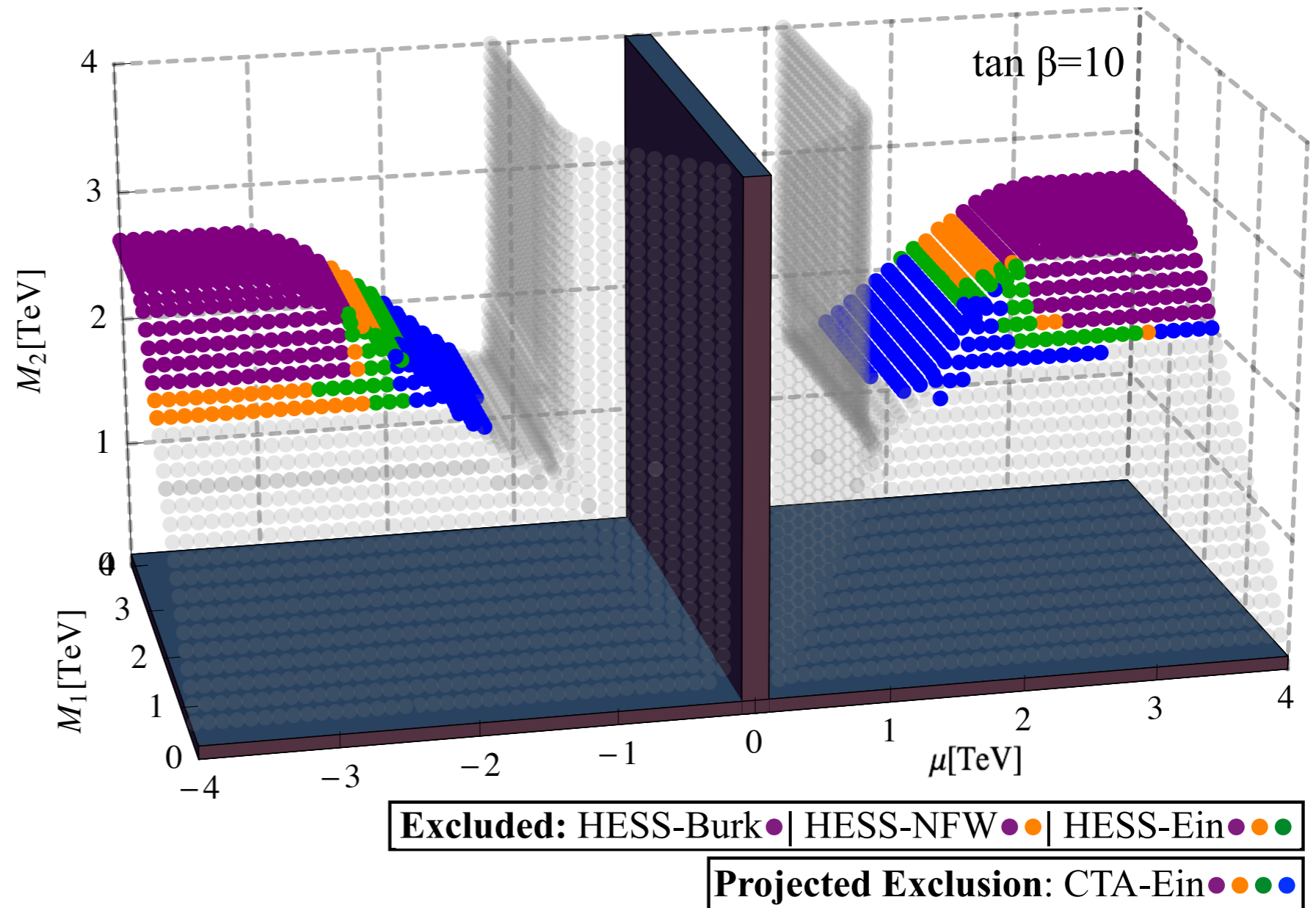
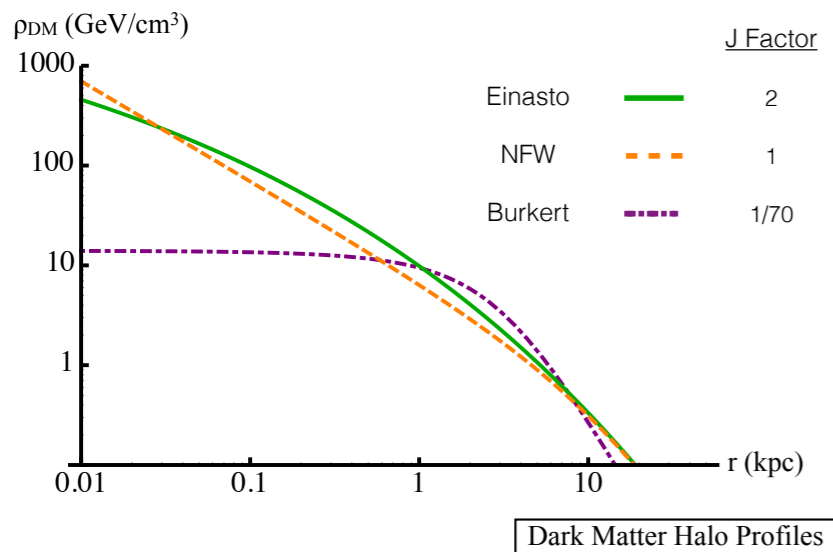
$$\frac{1}{2}\sigma_{\chi\chi\rightarrow\gamma Z} + \sigma_{\chi\chi\rightarrow\gamma\gamma} = \begin{array}{|c|c|c|c|c|c|c|} \hline \bullet & \bullet & \bullet & \bullet & \bullet & \bullet & \bullet \\ \hline >10^{-24} & 10^{-25} & 10^{-26} & 10^{-27} & 10^{-29} & 10^{-31} & <10^{-33} \text{ cm}^3/\text{s} \\ \hline \end{array}$$



Sommerfeld effect enhances annihilation to photons

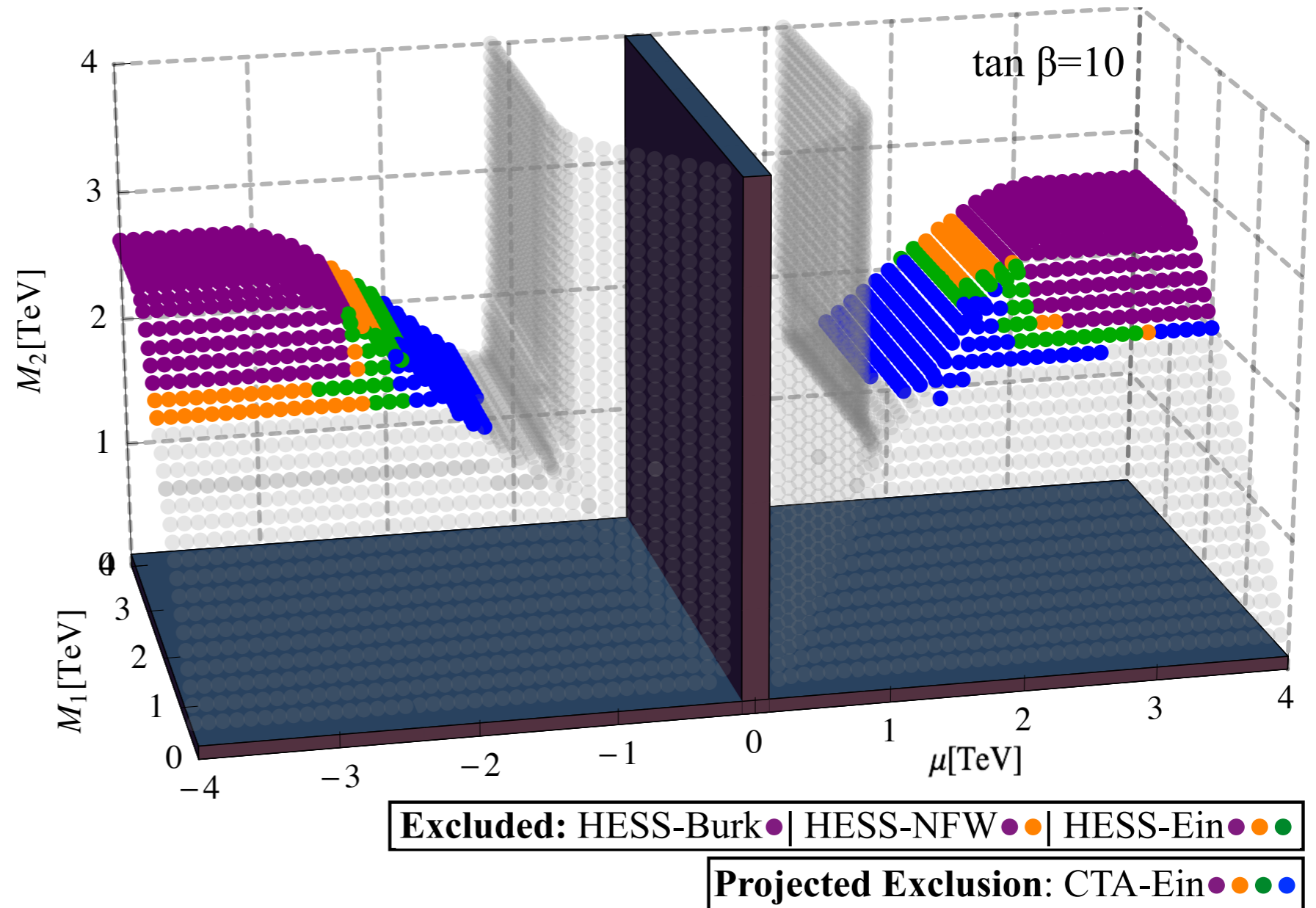
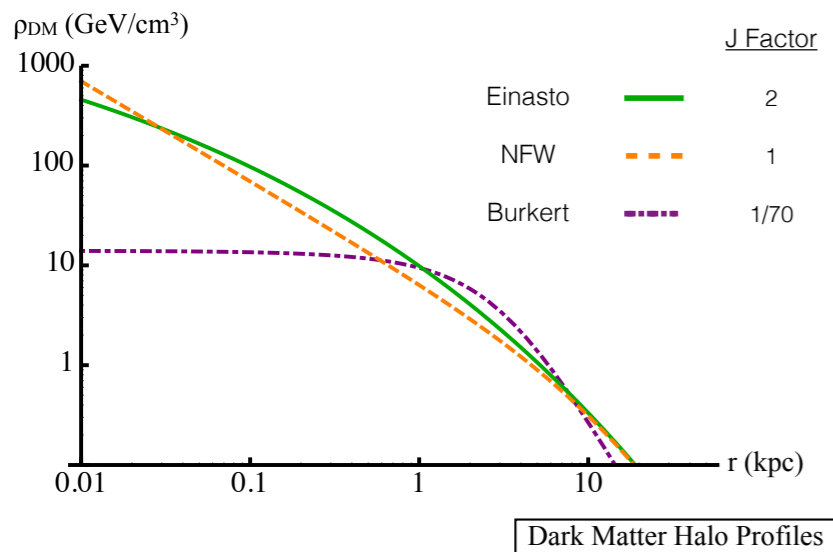
Indirect Detection

- Constraint depends on DM profile



Indirect Detection

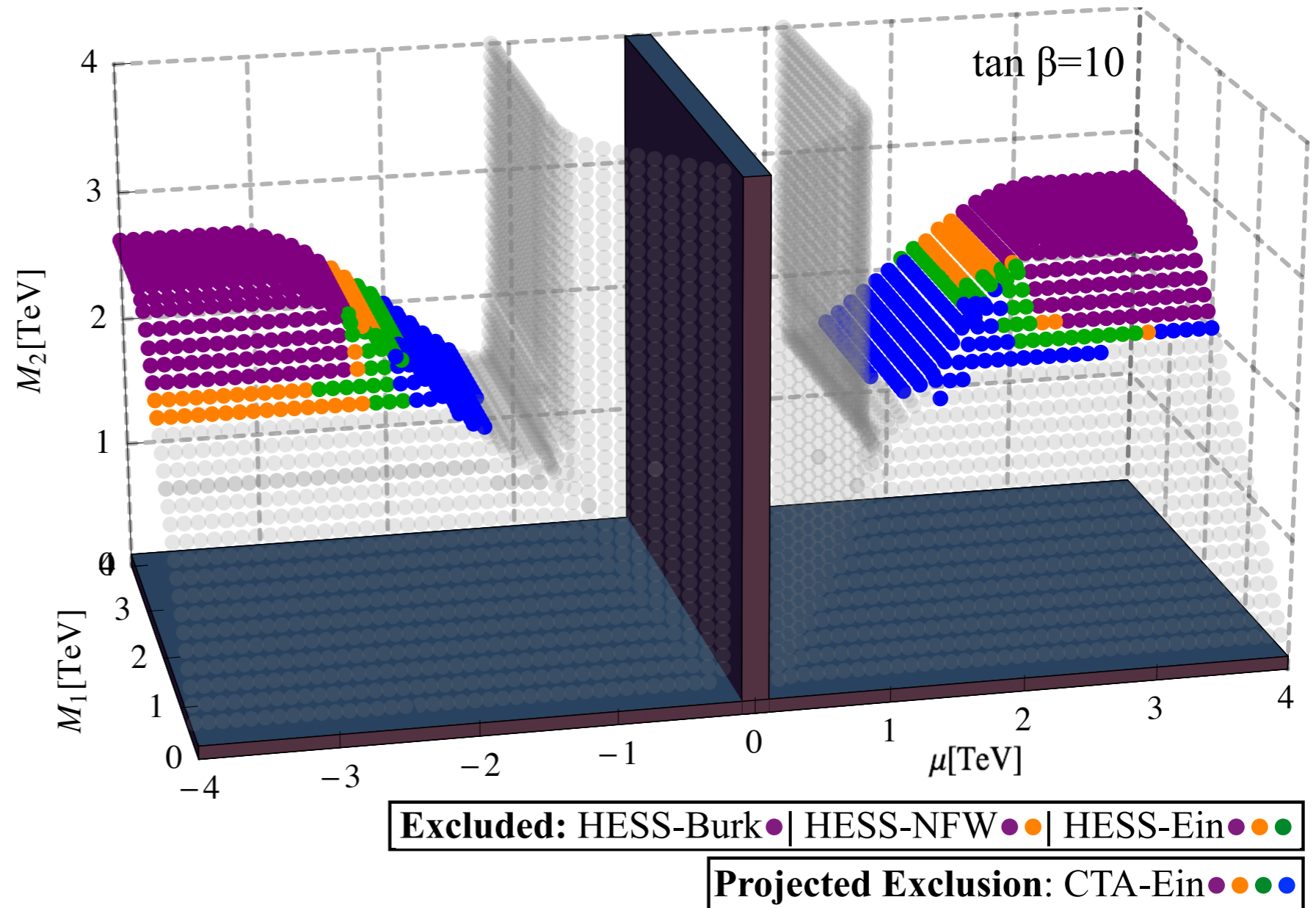
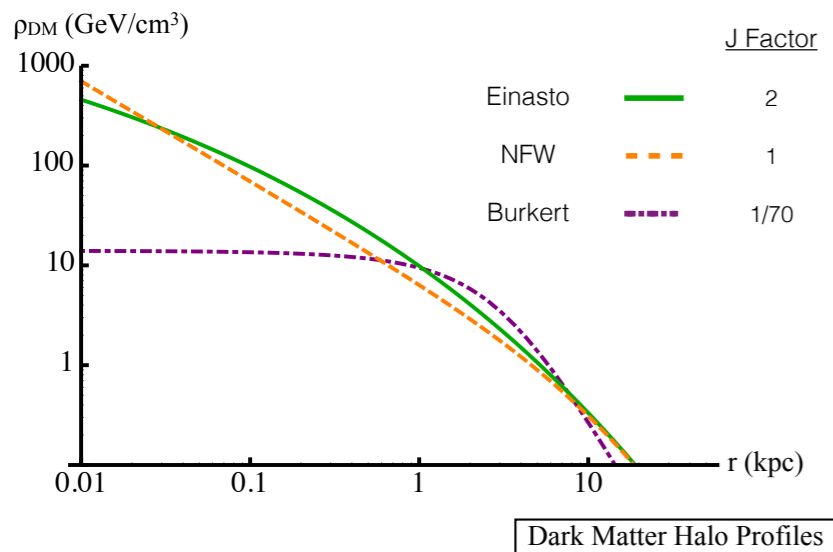
- Constraint depends on DM profile



- ★ Pure Wino plane excluded
- Wino-Higgsino sheet in future

Indirect Detection

- Constraint depends on DM profile



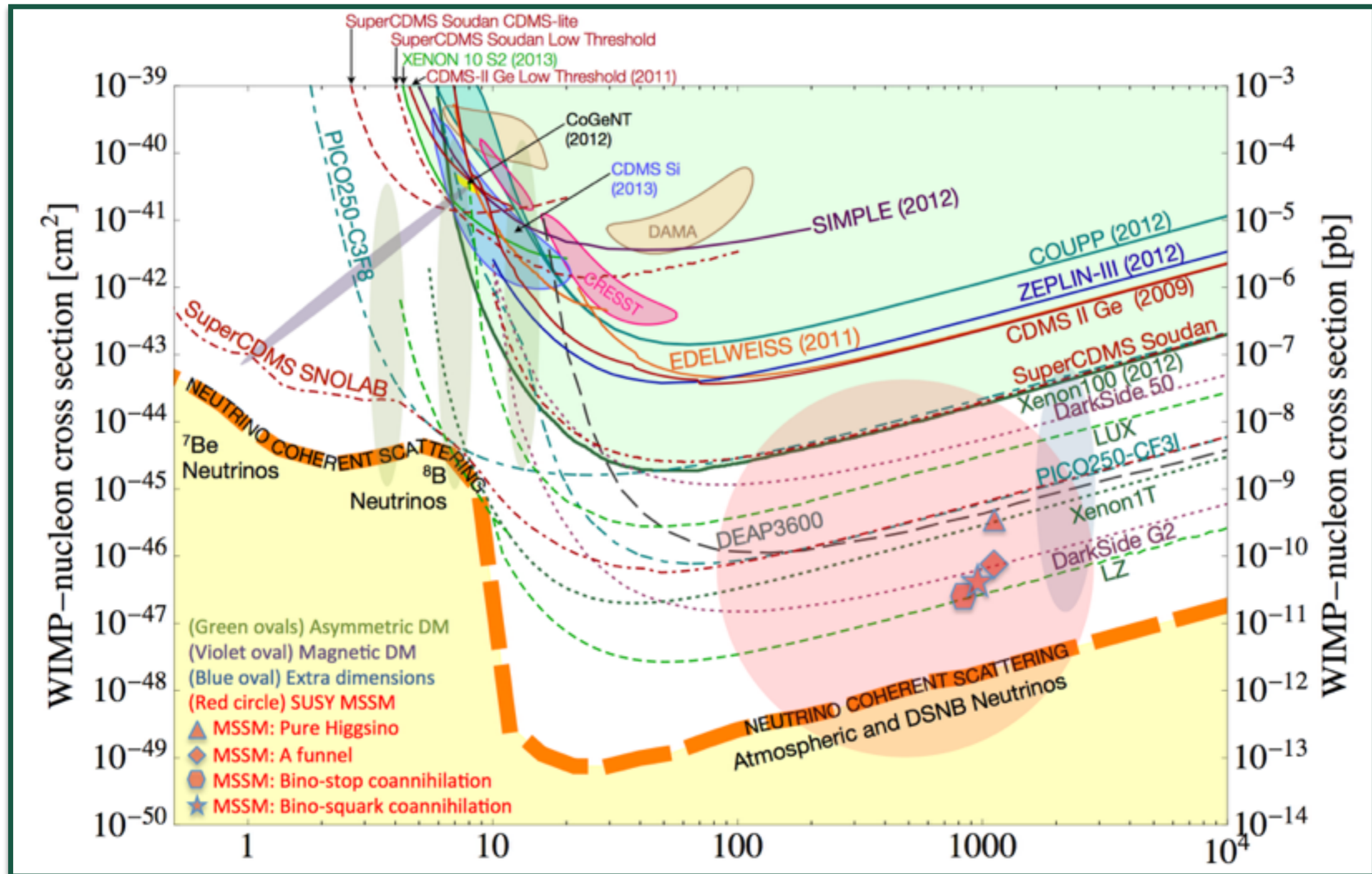
- ★ Pure Wino plane excluded
- Wino-Higgsino sheet in future

Direct Detection

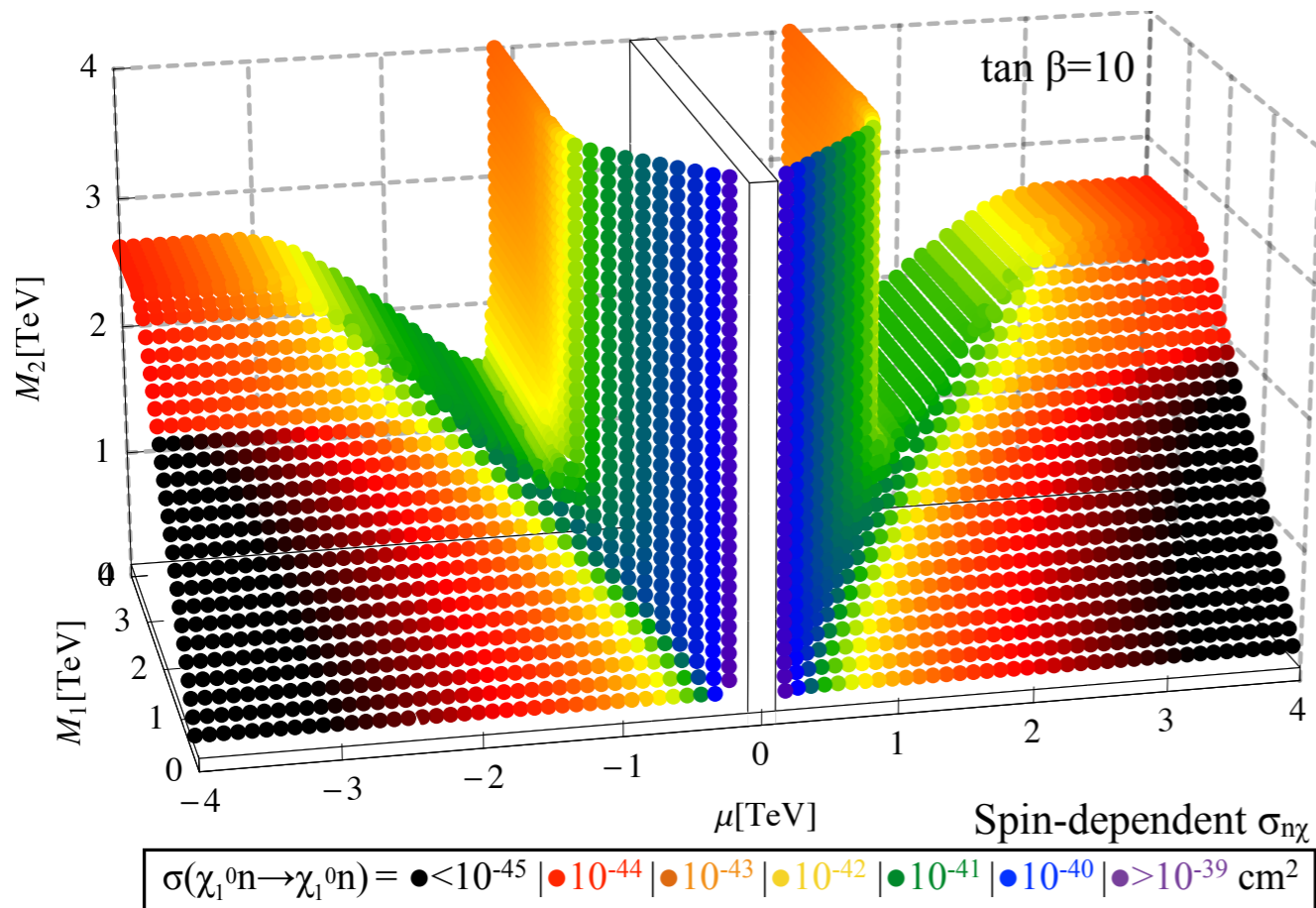
- Look for DM scattering off nucleon in detector material

Direct Detection

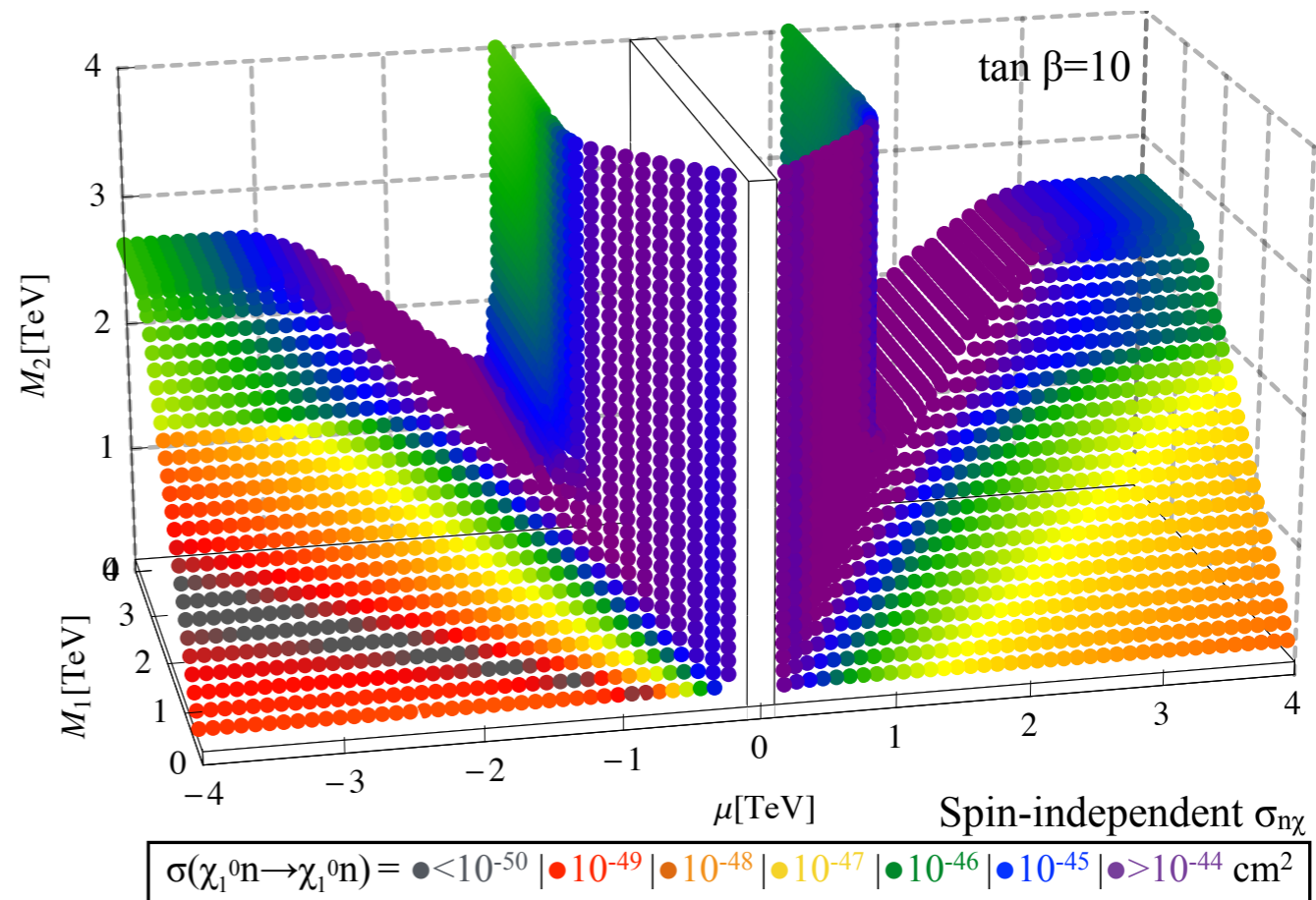
- Look for DM scattering off nucleon in detector material



Direct Detection



$\sigma \propto$ Coupling to Z

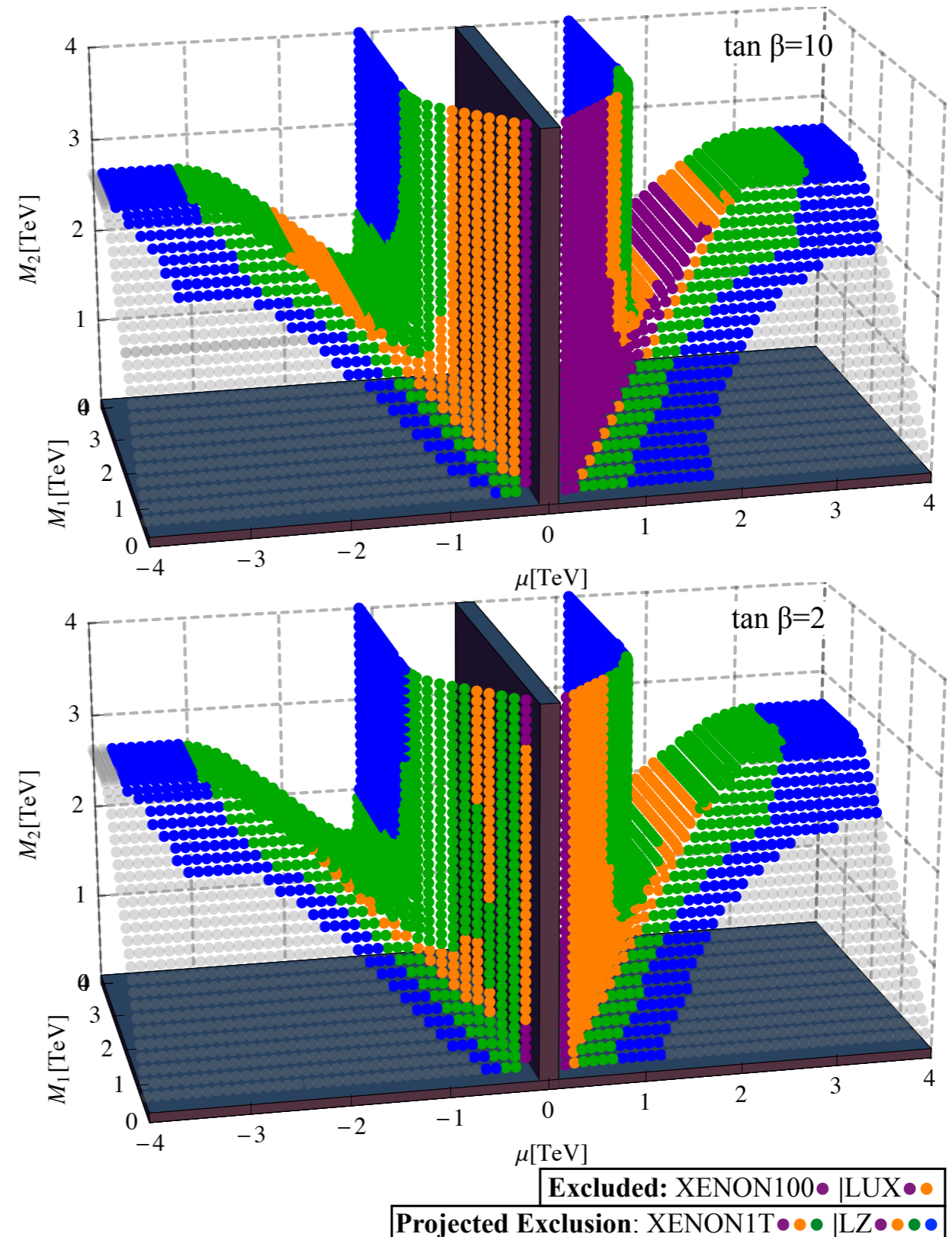


$\sigma \propto$ Coupling to h

Accidental cancelations in coupling of LSP to Z(h) leads to blind spots

Direct Detection

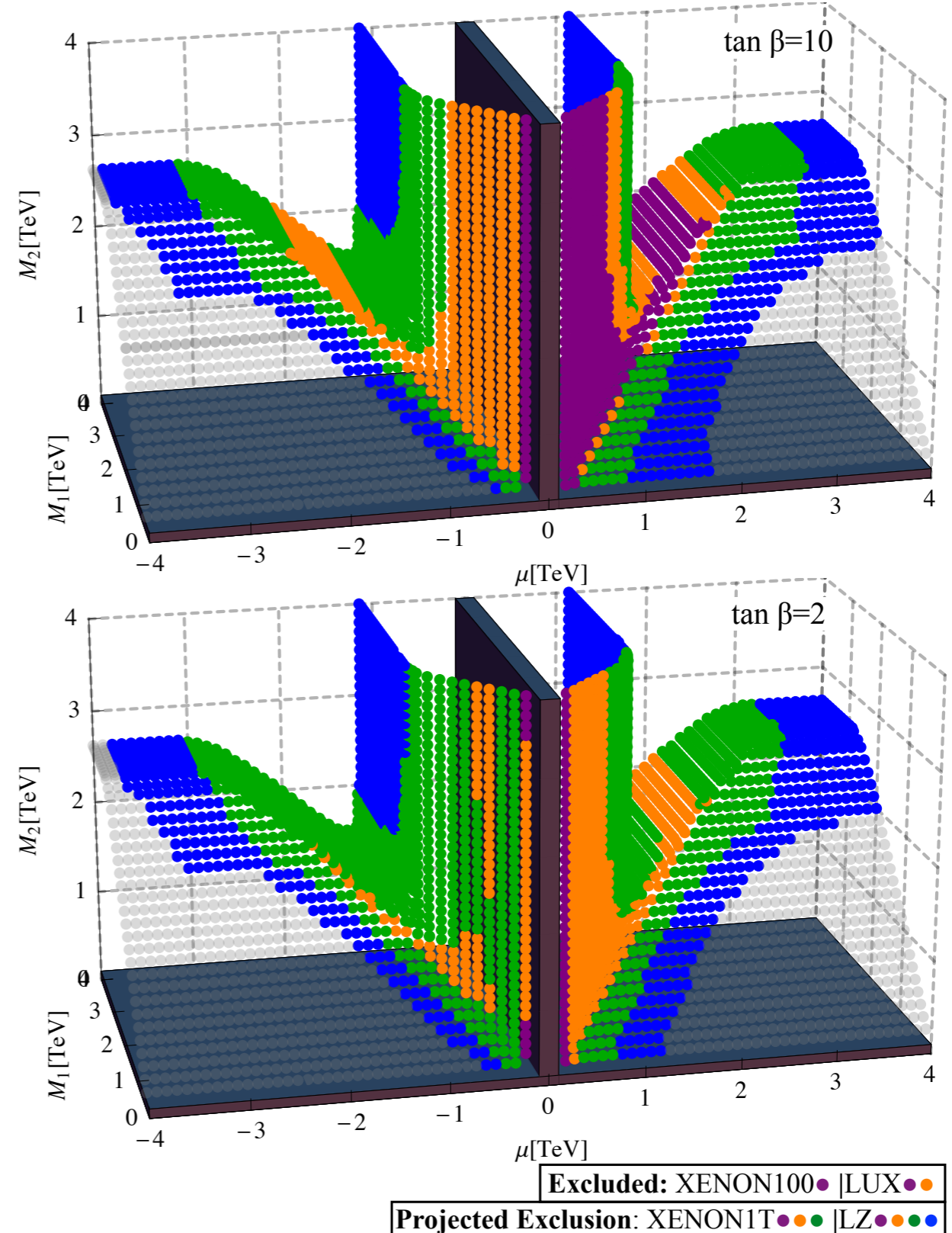
- Spin-independent bounds are much better
- Exclusions depend on $\tan \beta$



Direct Detection

- Spin-independent bounds are much better
- Exclusions depend on $\tan \beta$

- ★ Bino-Higgsino and Wino-Higgsino are or will be excluded
- ★ $M_1, M_2 \leq 4000$ GeV not decoupled enough...
`Pure' Higgsino discoverable in future
- ★ Pure Wino mostly covered



Can entire surface be discovered with current/future experiments?

- Examine Direct/Indirect Detection constraints
 - ★ Pure Wino and Wino-Higgsino covered by indirect detection
 - ★ Bino-Higgsino , Wino-Higgsino, 'Pure Higgsino', and some Wino surface covered by direct detection
 - * No coverage for Bino-Wino
 - * Depend on *astrophysical assumptions*
- Explore search techniques at high energy colliders
- Complementarity between experiments

Can entire surface be discovered with current/future experiments?

- Examine Direct/Indirect Detection constraints
- Explore search techniques at high energy colliders

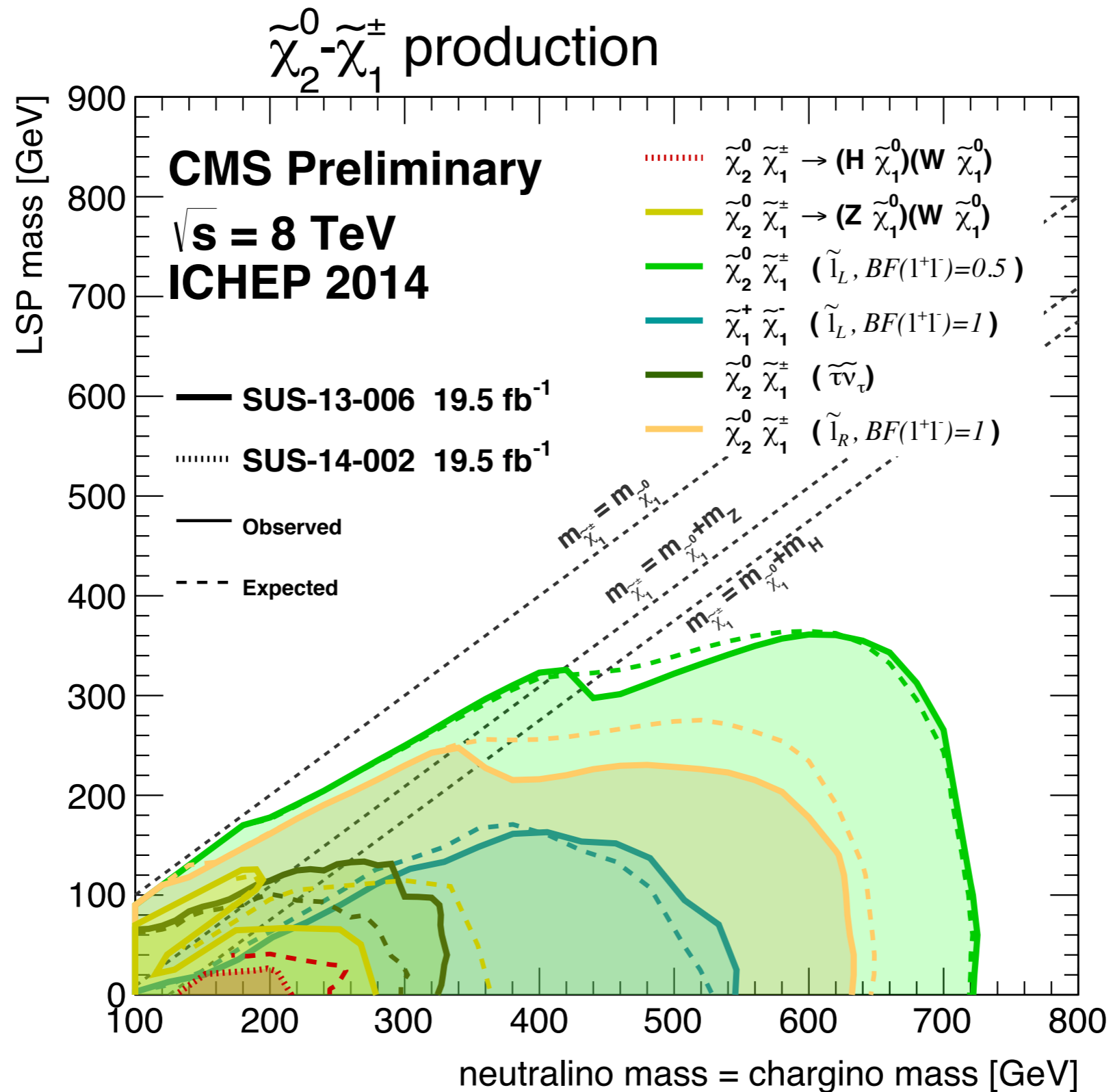
- Cover Bino-Wino surface
- Resolve conflicts with pure wino

LHC Detectors can distinguish

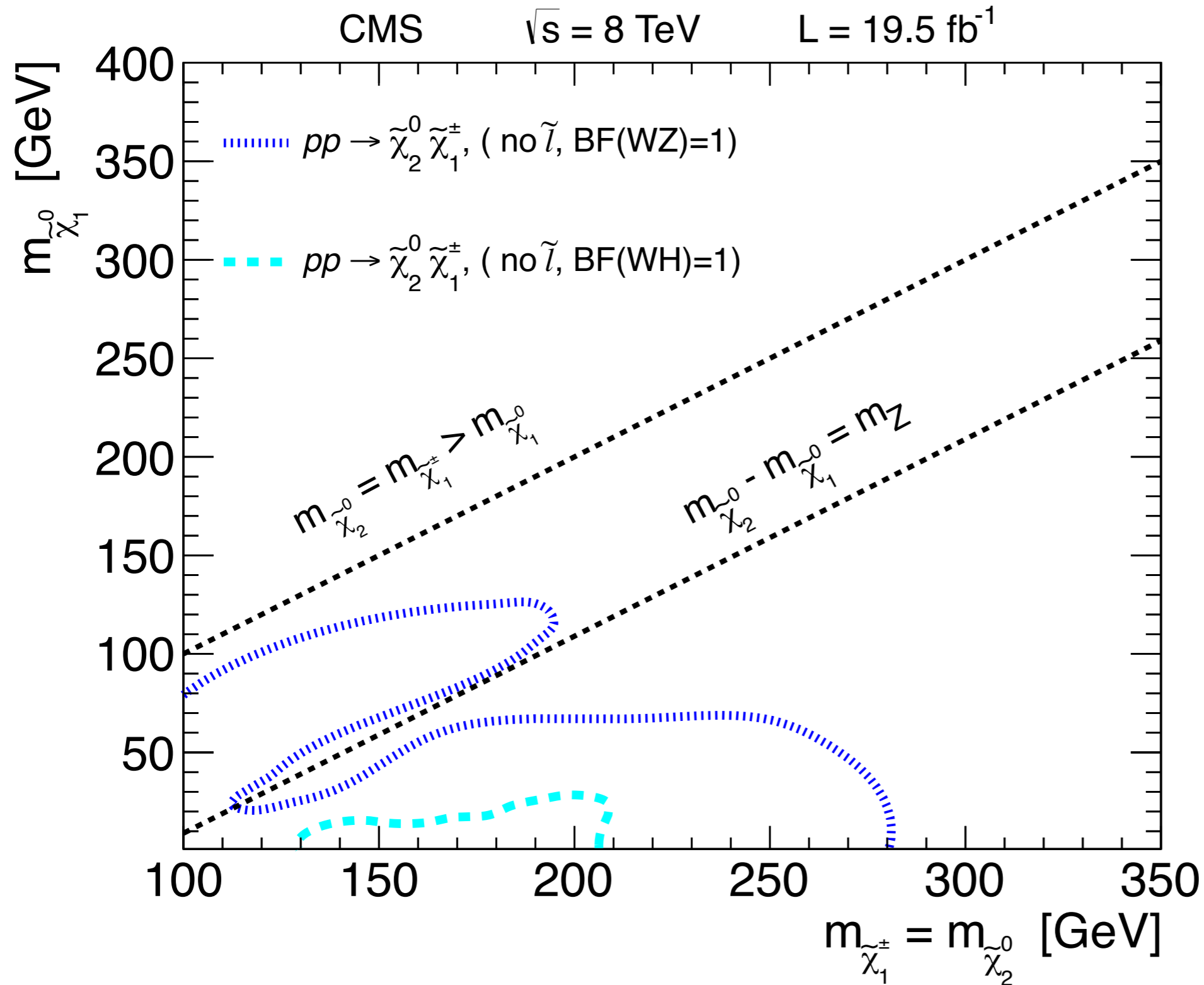
- Leptons (electrons and muons)
- Photons
- Anything with quarks seen as 'jet'
- Neutrinos and dark matter not detected; conservation of transverse momentum

- Complementarity between experiments

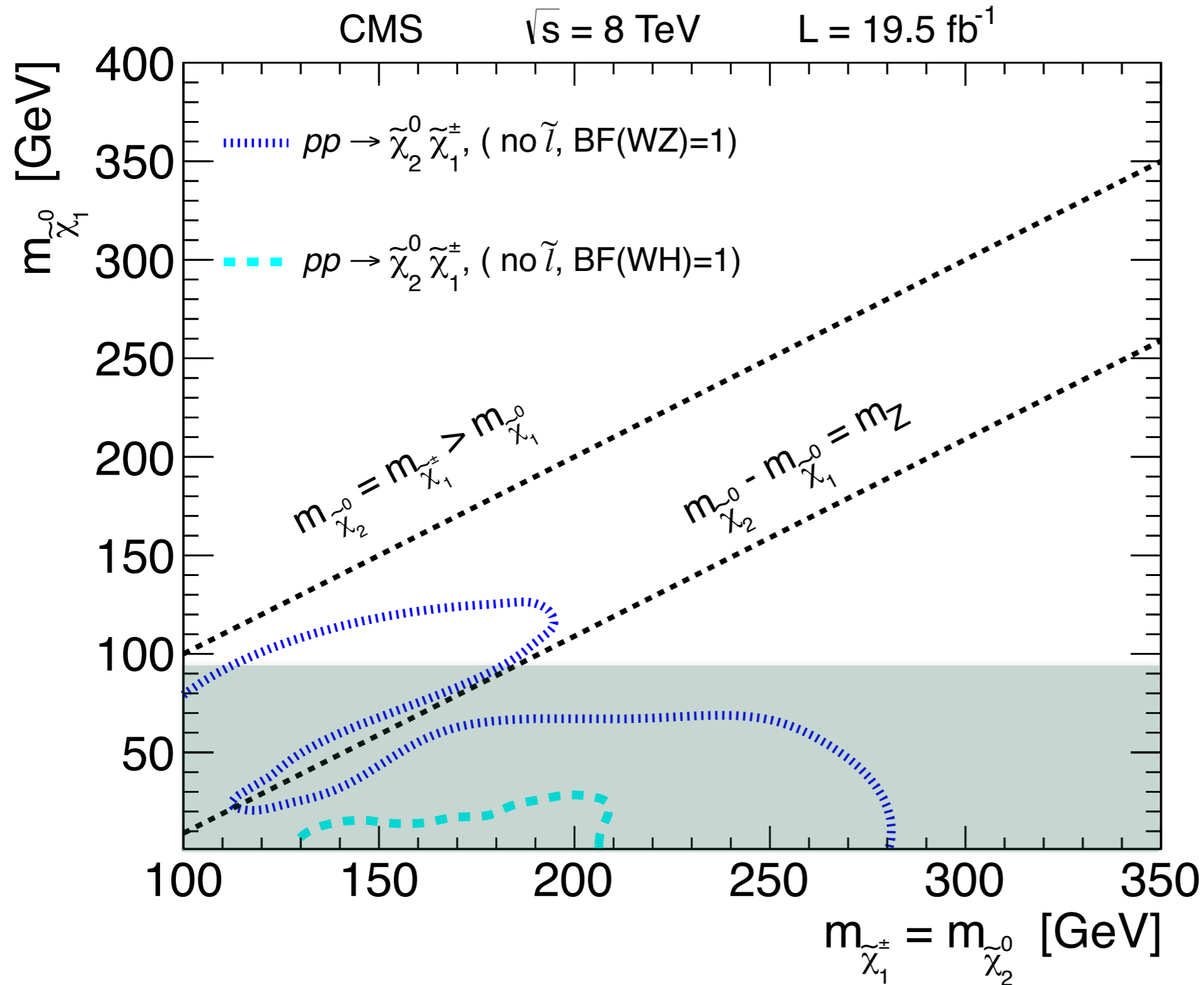
What limits exist already?



What limits exist already?



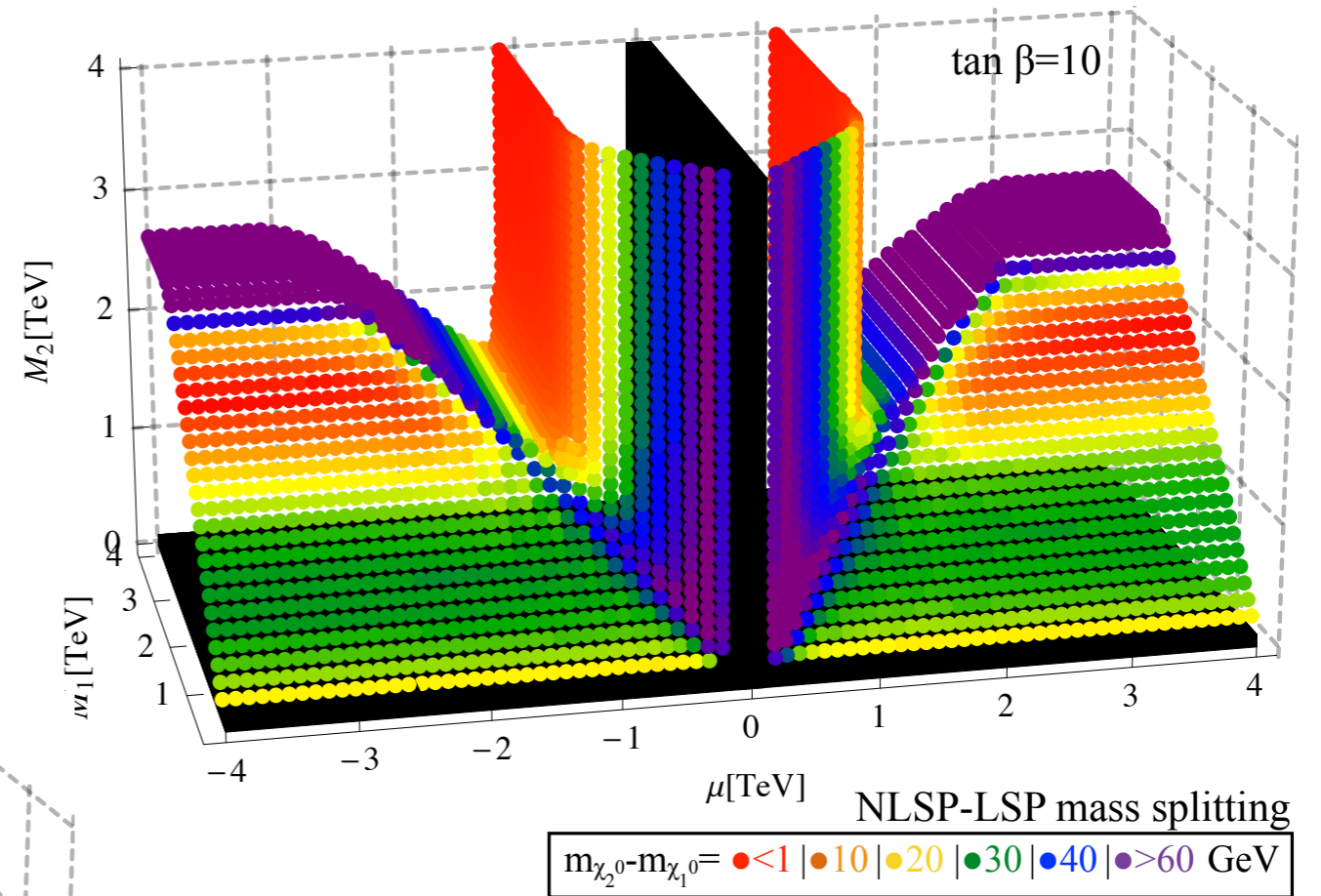
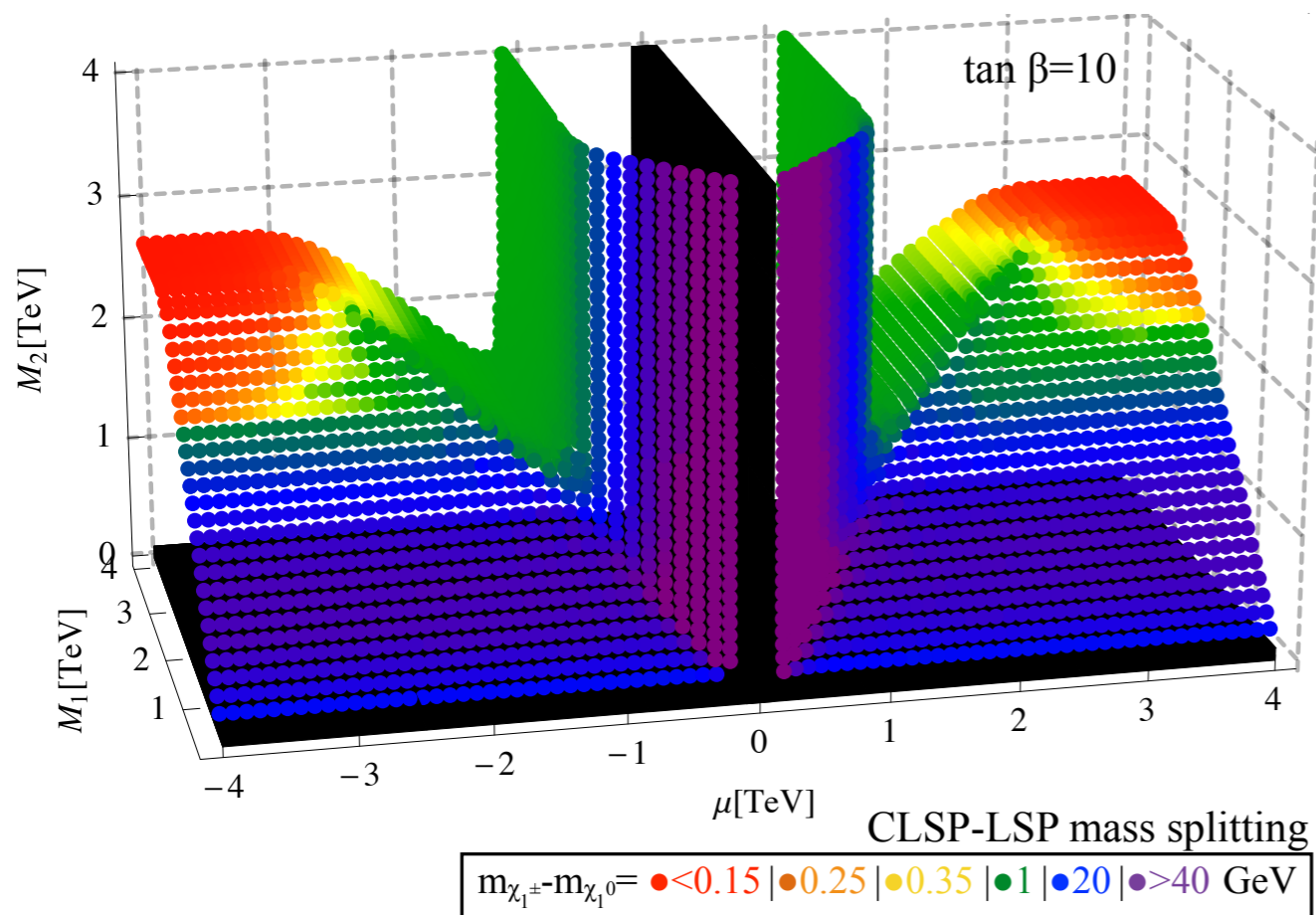
What limits exist already?



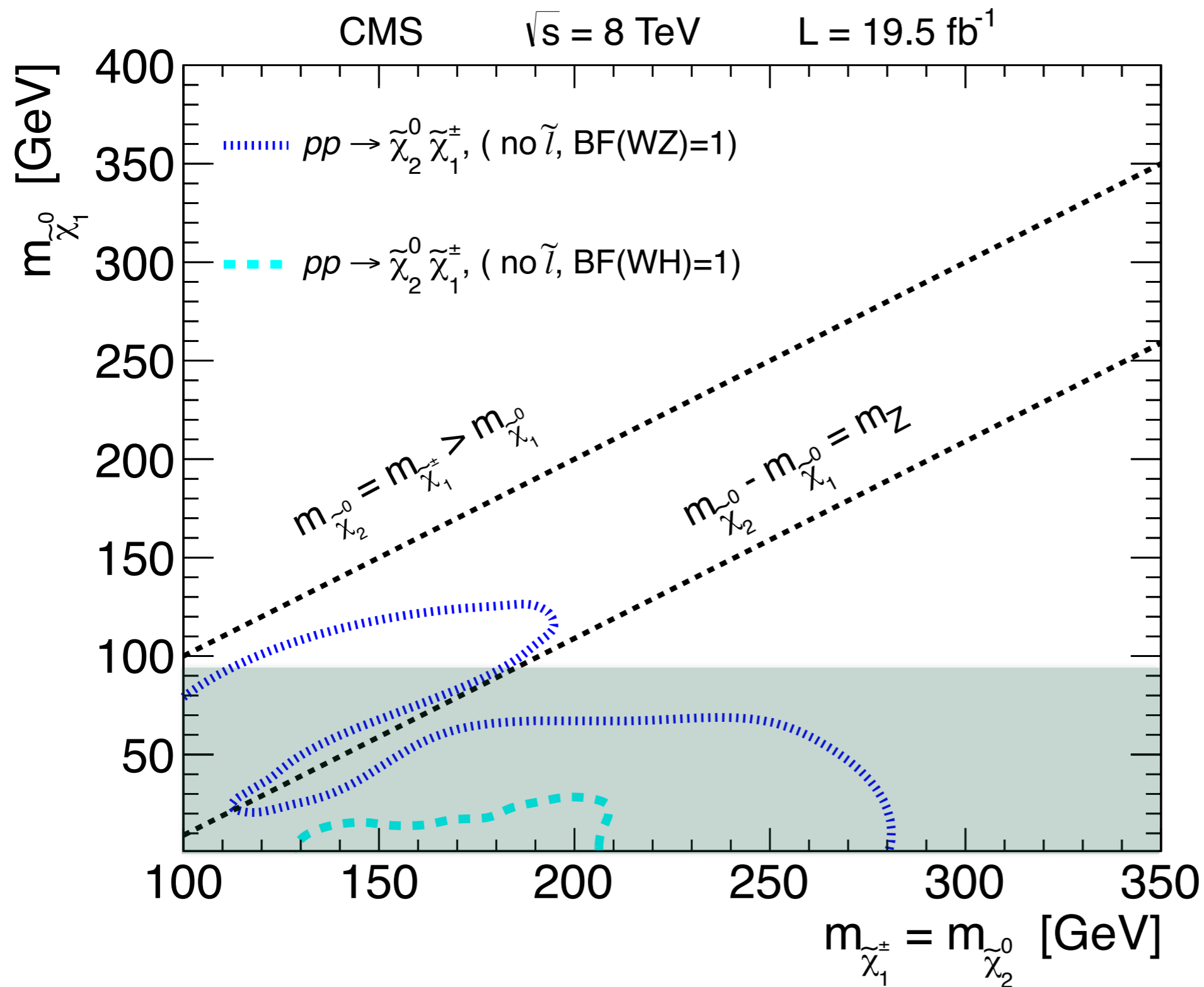
LEP: $m_{\chi^+} > 94 \text{ GeV}$

What limits exist already?

Reminder of mass splittings

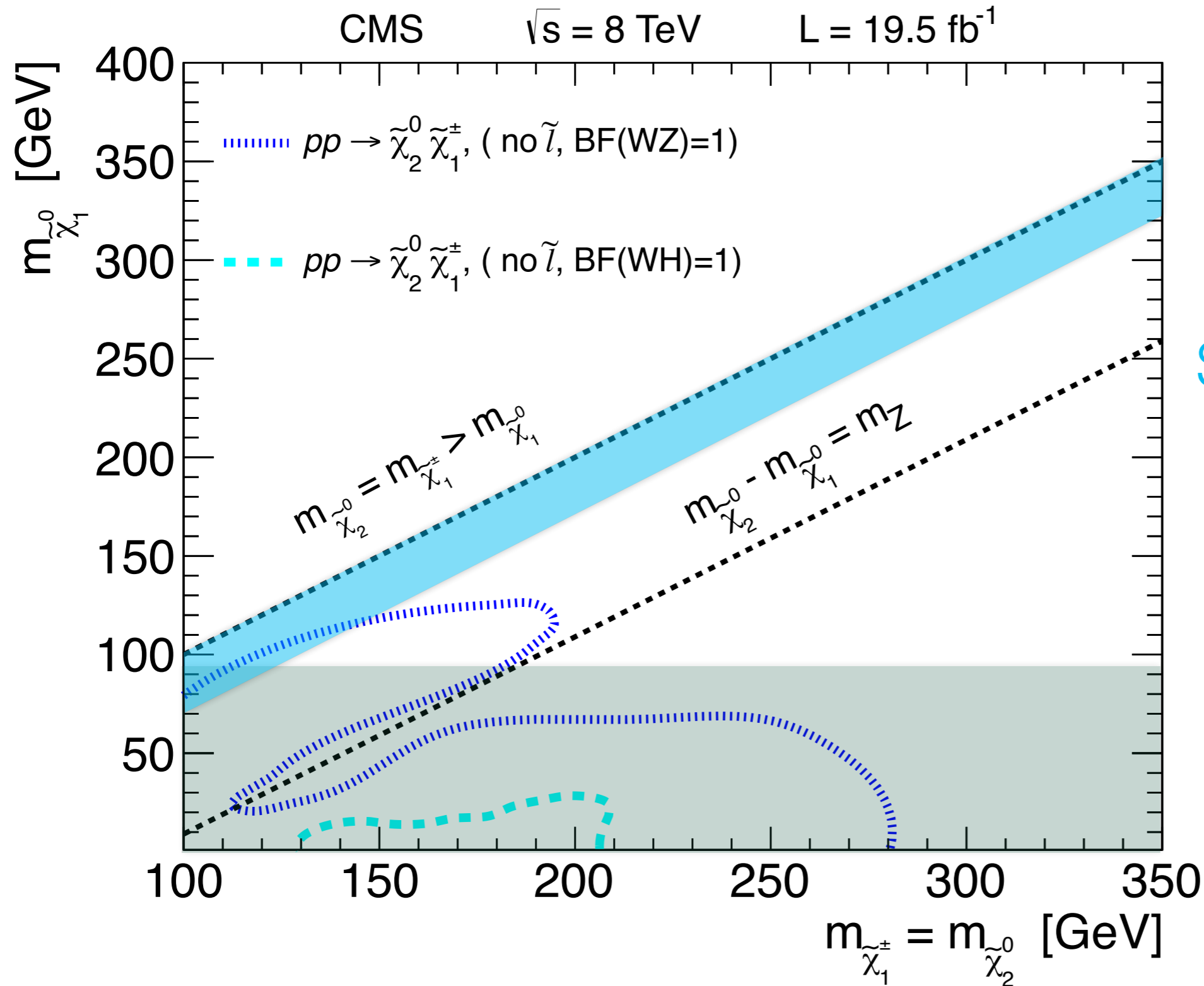


What limits exist already?



LEP: $m_{\chi^+} > 94 \text{ GeV}$

What limits exist already?



LEP: $m_{\chi^+} > 94 \text{ GeV}$

Surface:

$$m_{\chi^+}, m_{\chi_2^0} \lesssim m_{\chi_1^0} + 30$$

- Few SUSY searches so far use photons
- Can light be used to search for dark matter?

$$pp \rightarrow \gamma + \ell^+ \ell^- + \cancel{E}_T$$

Neutral current aimed at production of $\chi_2^0 \chi_3^0$

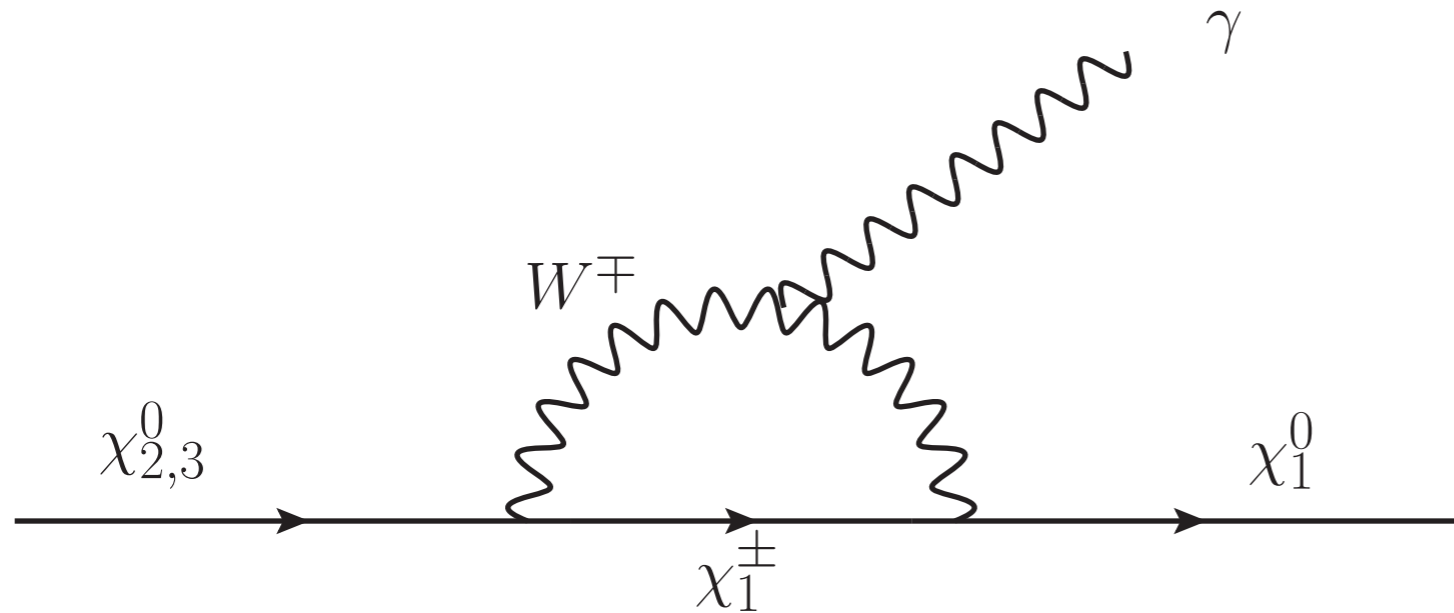
- J. Bramante, A. Delgado, F. Elahi, A. Martin and BO, “Catching sparks from well-forged neutralinos,” Phys. Rev. D 90, no. 9, 095008 (2014) [arXiv:1408.6530 [hep-ph]]

$$pp \rightarrow \gamma + \ell^\pm + \cancel{E}_T$$

Charged current aimed at production of $\chi_2^0 \chi_1^\pm$

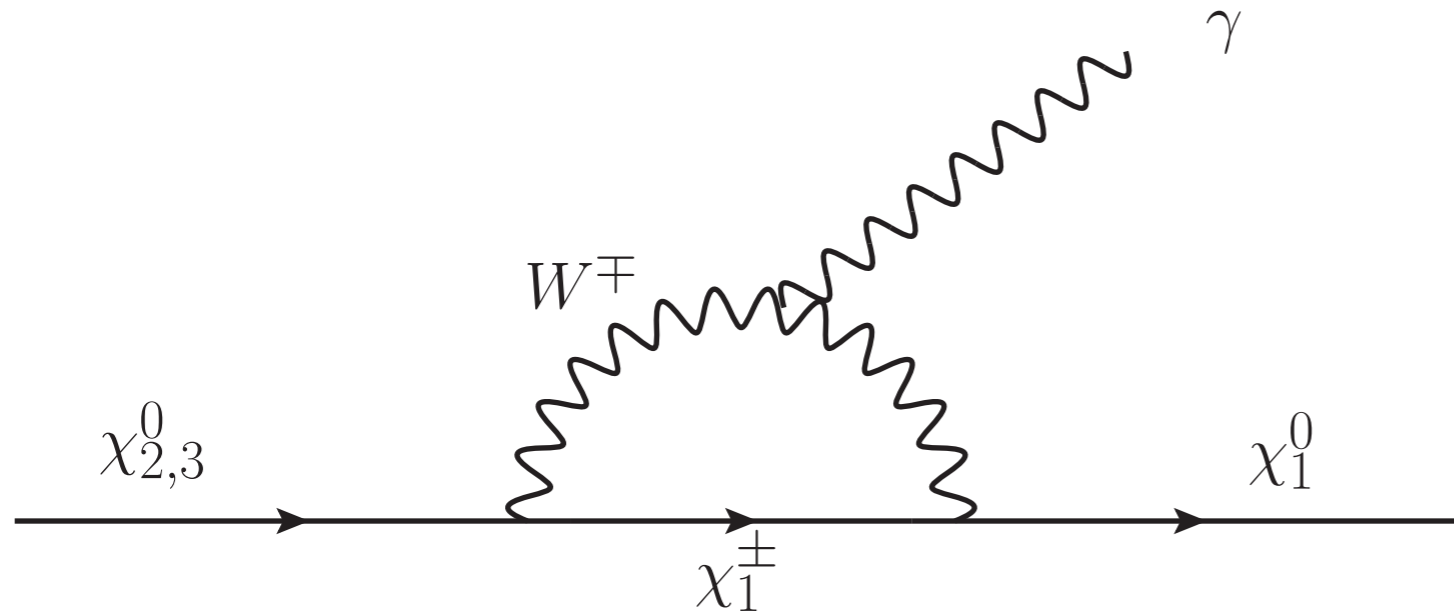
- J. Bramante, P. J. Fox, A. Martin, BO, T. Plehn, T. Schell and M. Takeuchi, “Relic neutralino surface at a 100 TeV collider,” Phys. Rev. D 91, no. 5, 054015 (2015) [arXiv:1412.4789 [hep-ph]]
- J. Bramante, N. Desai, P. Fox, A. Martin, BO, and T. Plehn, “Towards the Final Word on Neutralino Dark Matter,” arXiv:1510.03460 [hep-ph]

How to get a photon

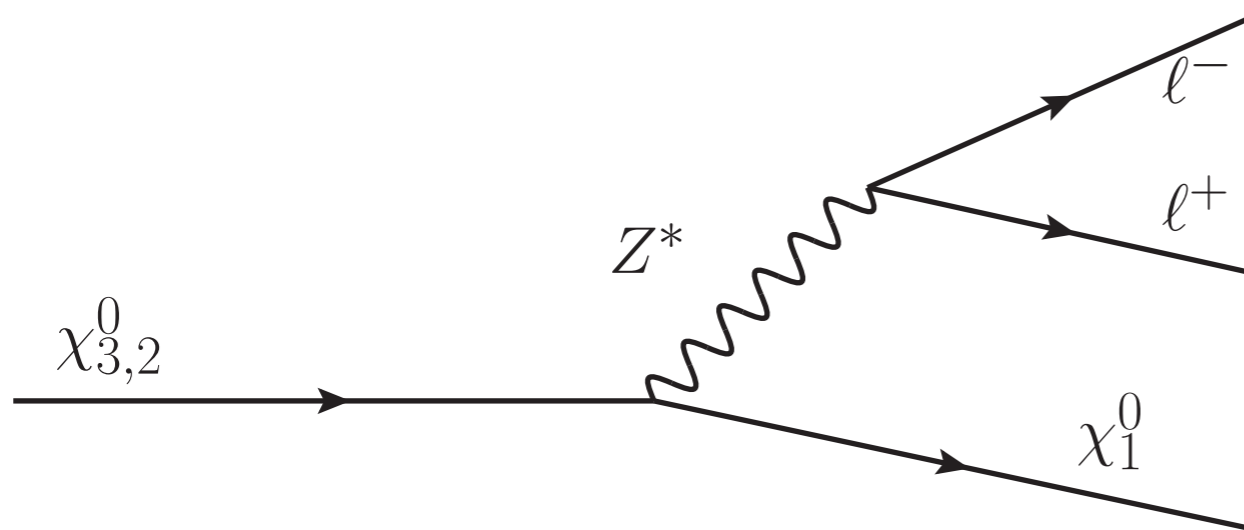


- 2-body phase space
- Loop factor $\sim 1/16\pi^2$

How to get a photon

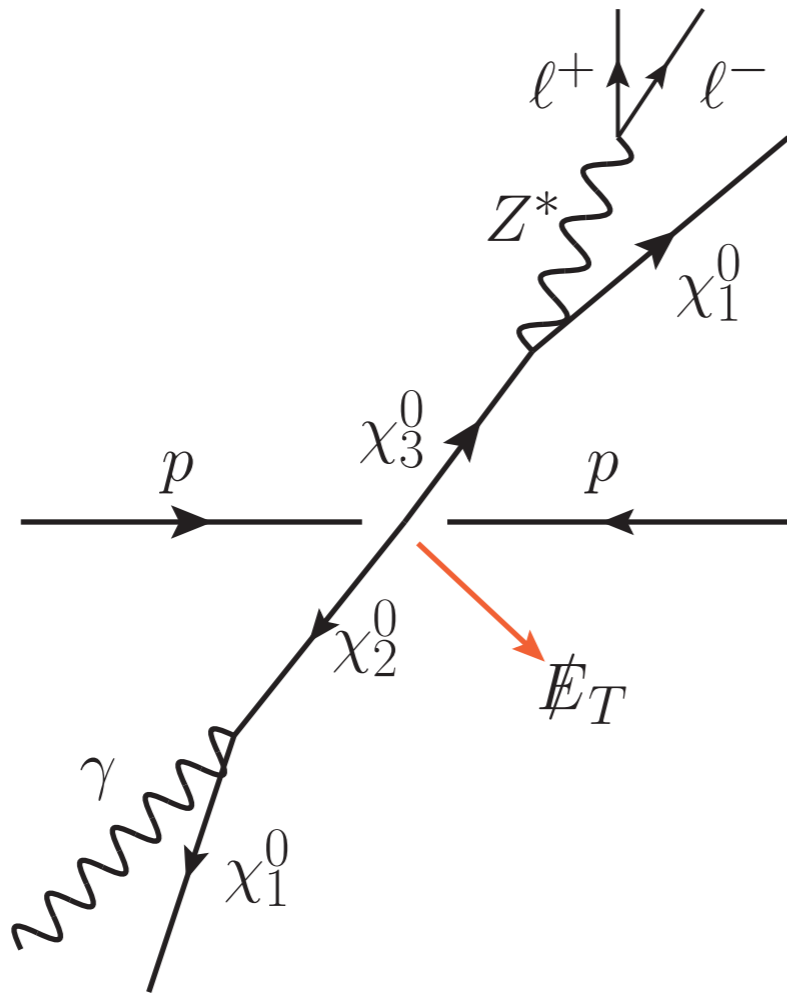


- 2-body phase space
- Loop factor $\sim 1/16\pi^2$



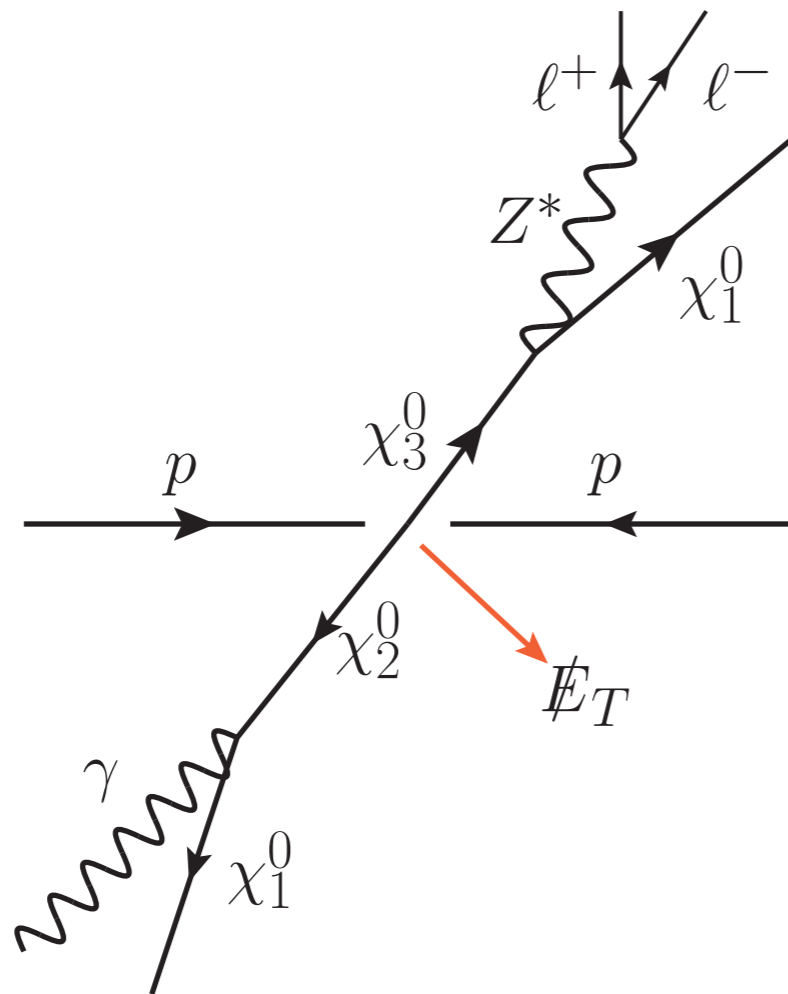
- 3-body phase space
 $\sim 2\text{BPS} * \frac{1}{(2\pi)^3} \frac{1}{2} \frac{d^3p}{E_p}$
- Z propagator $\sim \frac{1}{p_Z^2 - m_Z^2}$
- $\text{BR}(Z \rightarrow l^+l^-) \sim 7\%$

$$pp \rightarrow \chi_2^0 \chi_3^0 \rightarrow \gamma + \ell^+ \ell^- + \cancel{E}_T$$



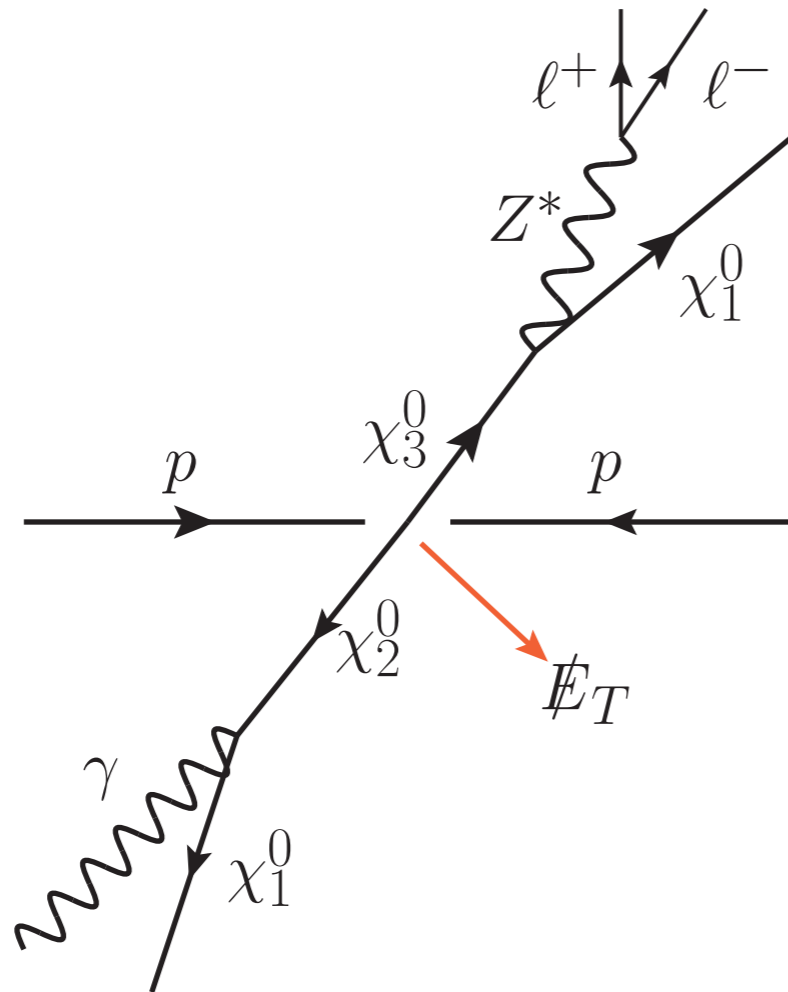
$$pp \rightarrow \chi_2^0 \chi_3^0 \rightarrow \gamma + \ell^+ \ell^- + \cancel{E}_T$$

Only works in the Bino-Higgsino part of the surface



$$pp \rightarrow \chi_2^0 \chi_3^0 \rightarrow \gamma + \ell^+ \ell^- + \cancel{E}_T$$

Only works in the Bino-Higgsino part of the surface

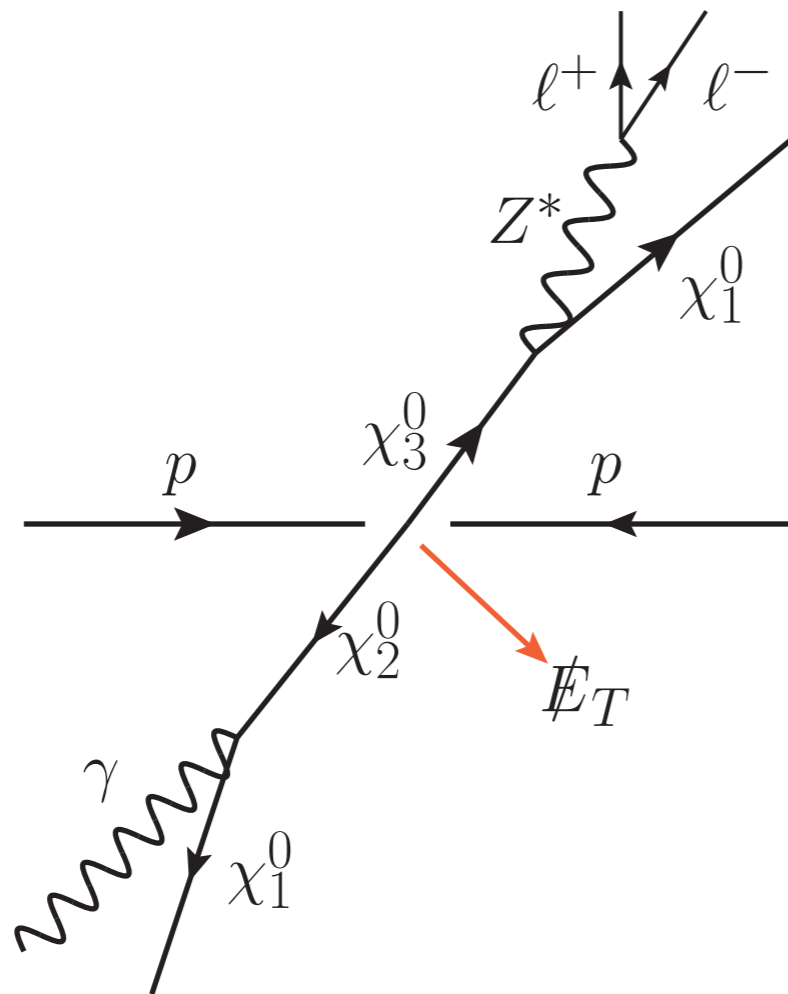


SM Backgrounds

- $\gamma t\bar{t}$ | dilepton
- $\gamma \gamma^* / Z (\tau^+ \tau^-)$ | dilepton
- γVV | dilepton

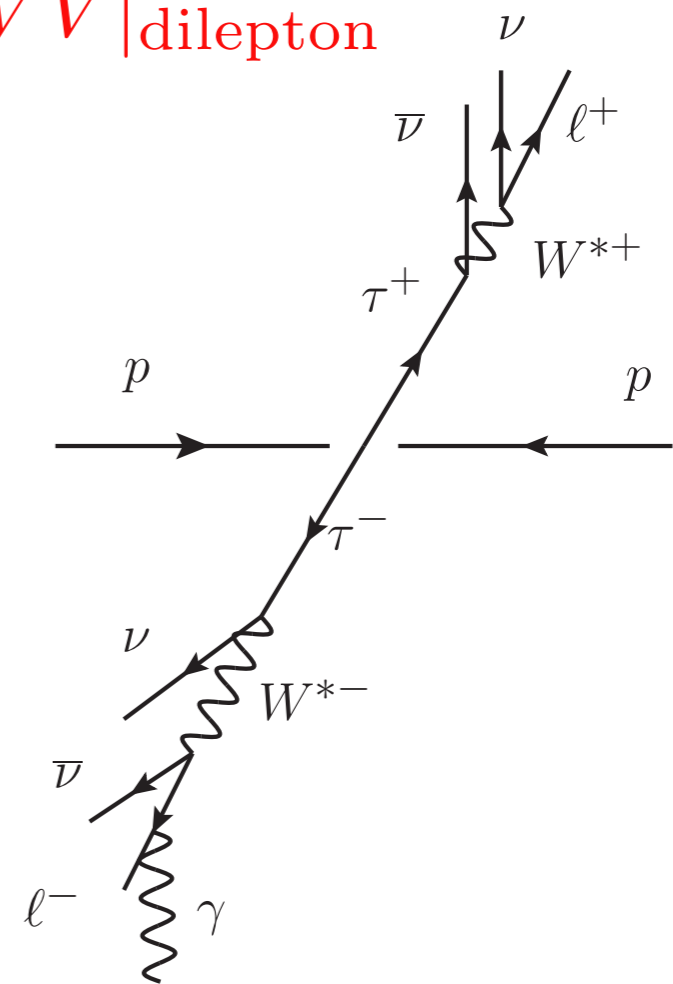
$$pp \rightarrow \chi_2^0 \chi_3^0 \rightarrow \gamma + \ell^+ \ell^- + \cancel{E}_T$$

Only works in the Bino-Higgsino part of the surface



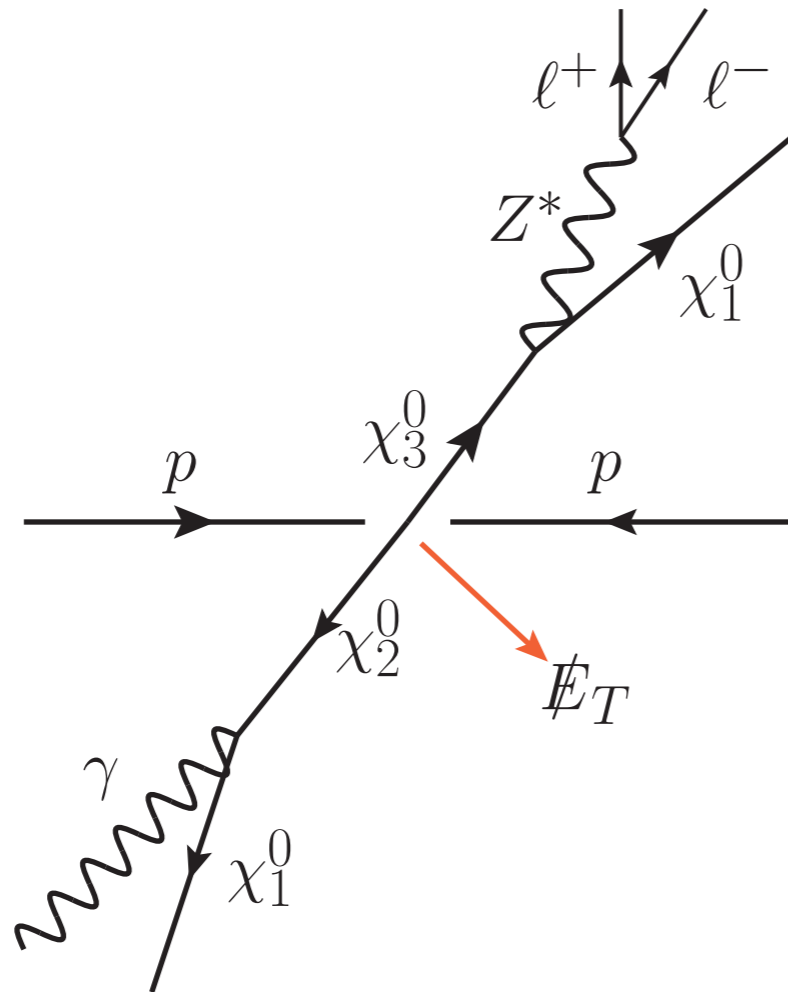
SM Backgrounds

- $\gamma t\bar{t}$ | dilepton
- $\gamma \gamma^* / Z (\tau^+ \tau^-)$ | dilepton
- γVV | dilepton



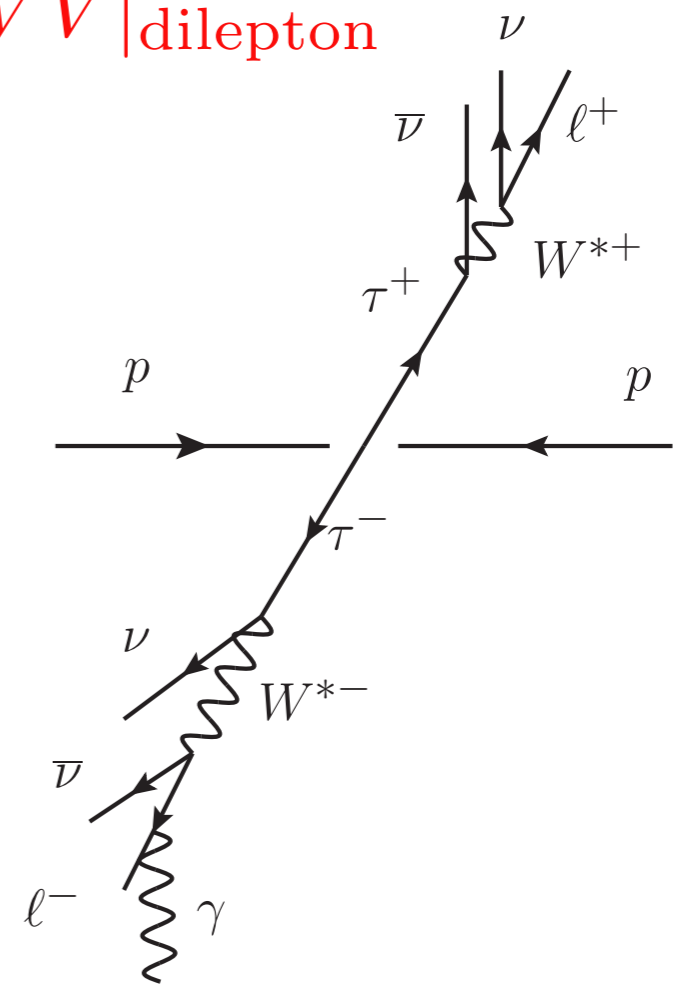
$$pp \rightarrow \chi_2^0 \chi_3^0 \rightarrow \gamma + \ell^+ \ell^- + \cancel{E}_T$$

Only works in the Bino-Higgsino part of the surface



SM Backgrounds

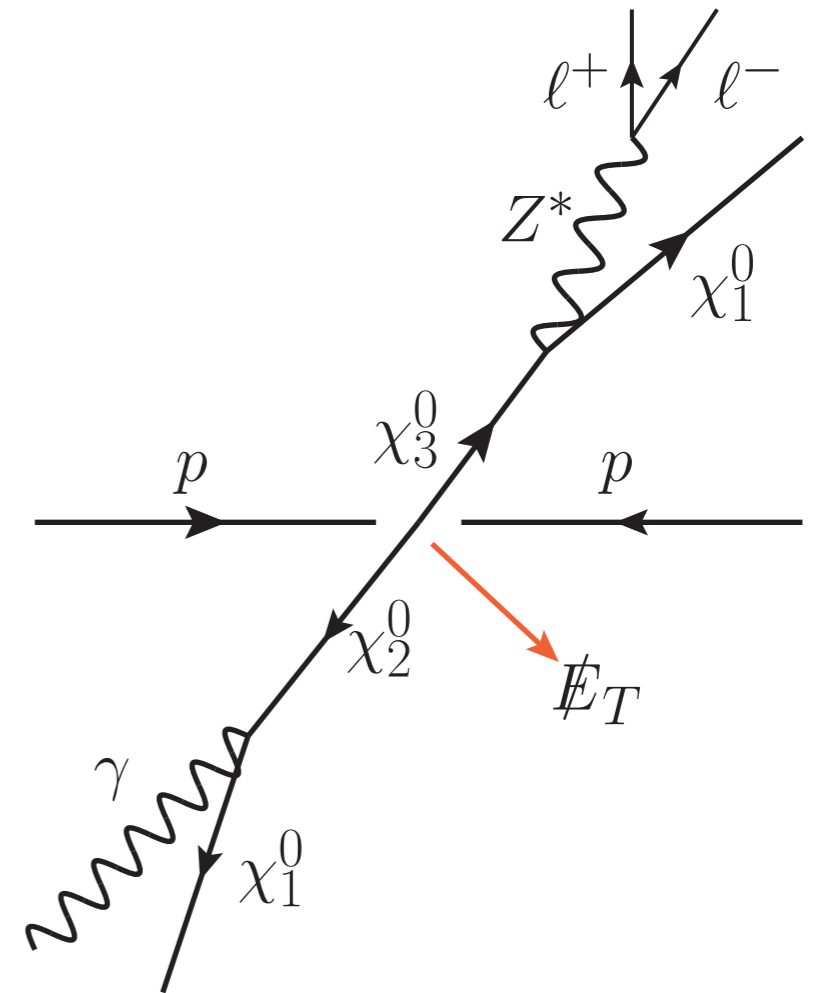
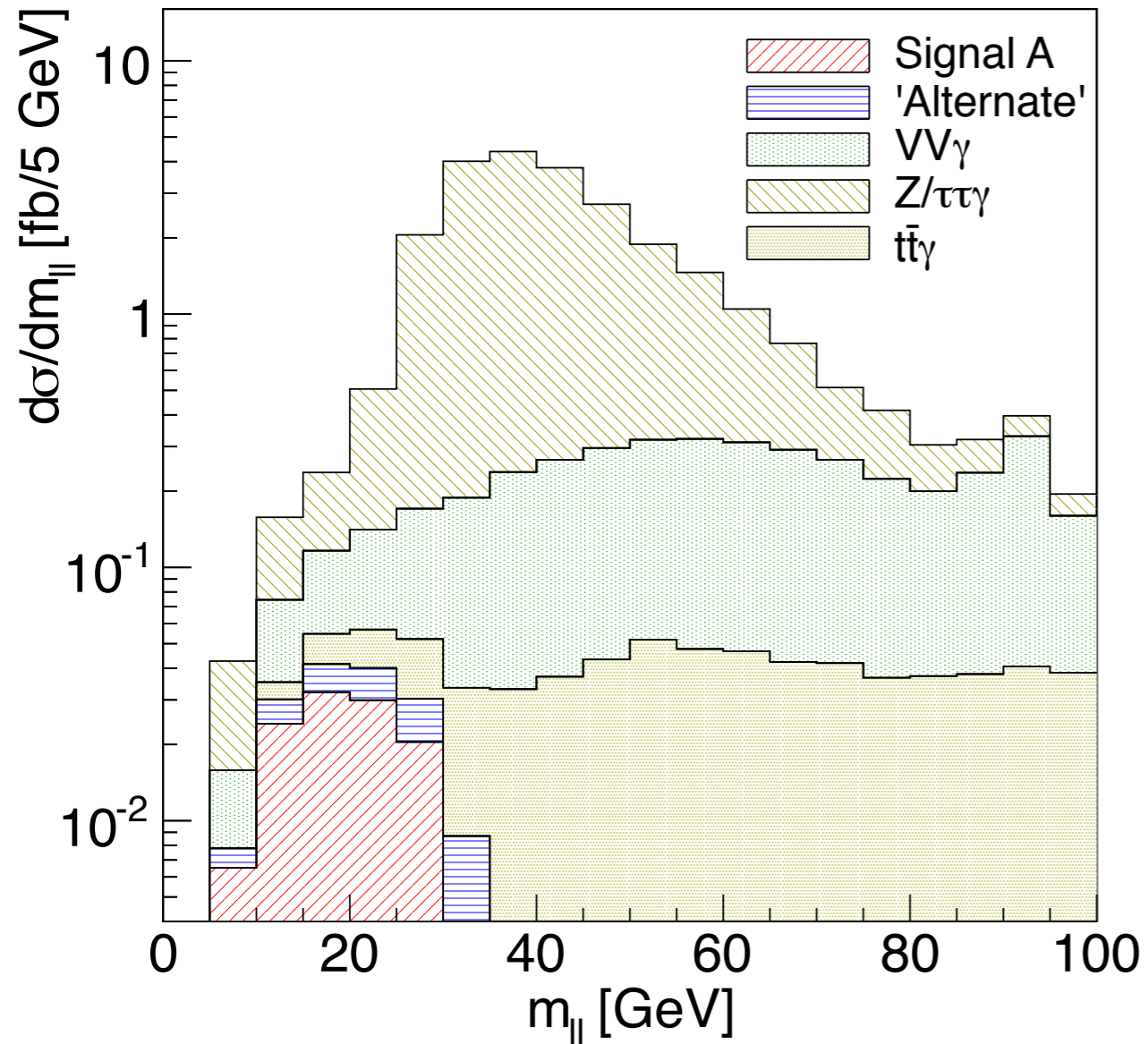
- $\gamma t\bar{t}$ | dilepton
- $\gamma \gamma^* / Z (\tau^+ \tau^-)$ | dilepton
- γVV | dilepton



- Angle between leptons
- Transverse mass
- Angle between leptons and photon
- $p_T(\gamma)$, Total MET

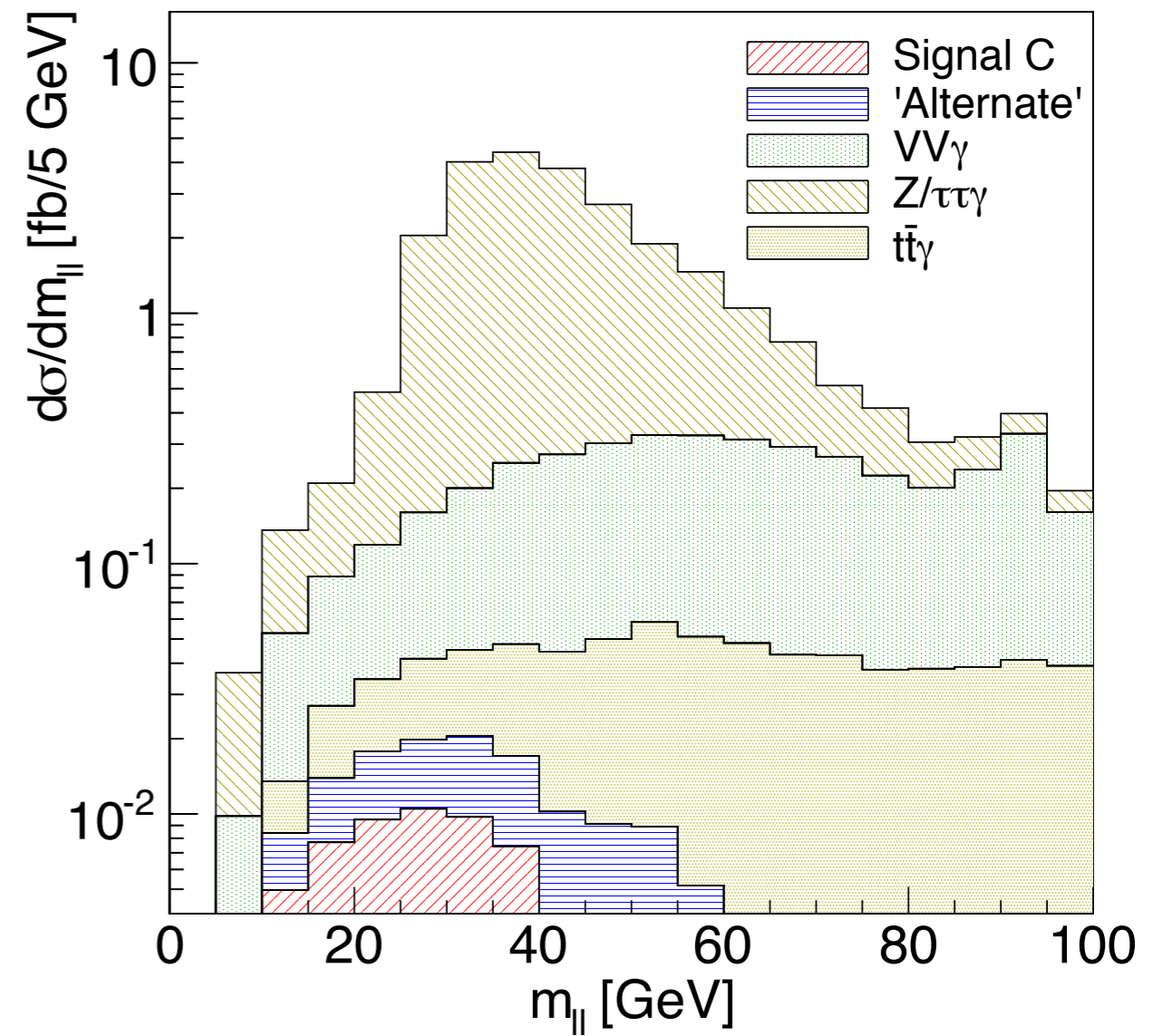
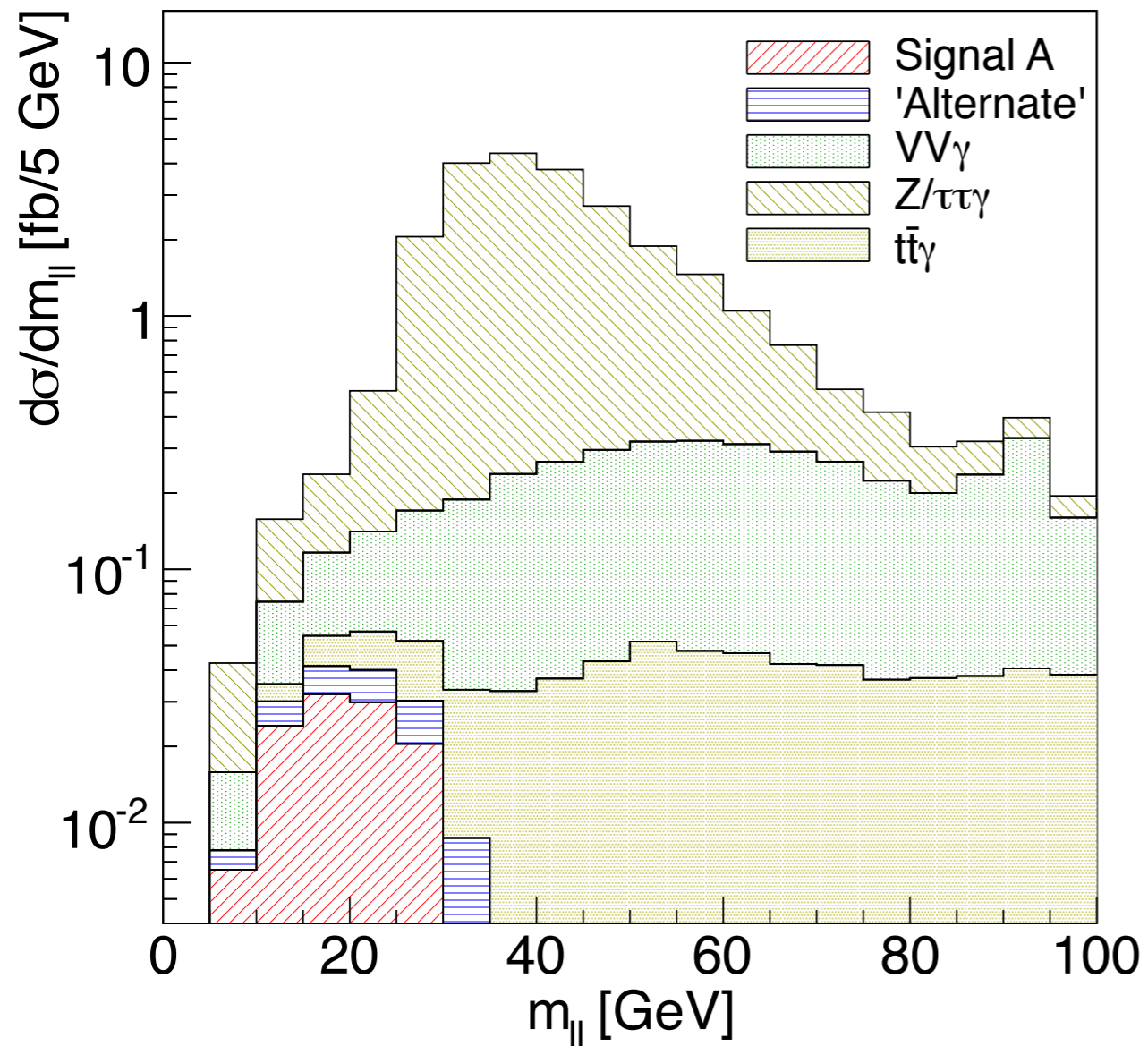
$$pp \rightarrow \chi_2^0 \chi_3^0 \rightarrow \gamma + \ell^+ \ell^- + \cancel{E}_T$$

Most discriminating cut is m_{ll}



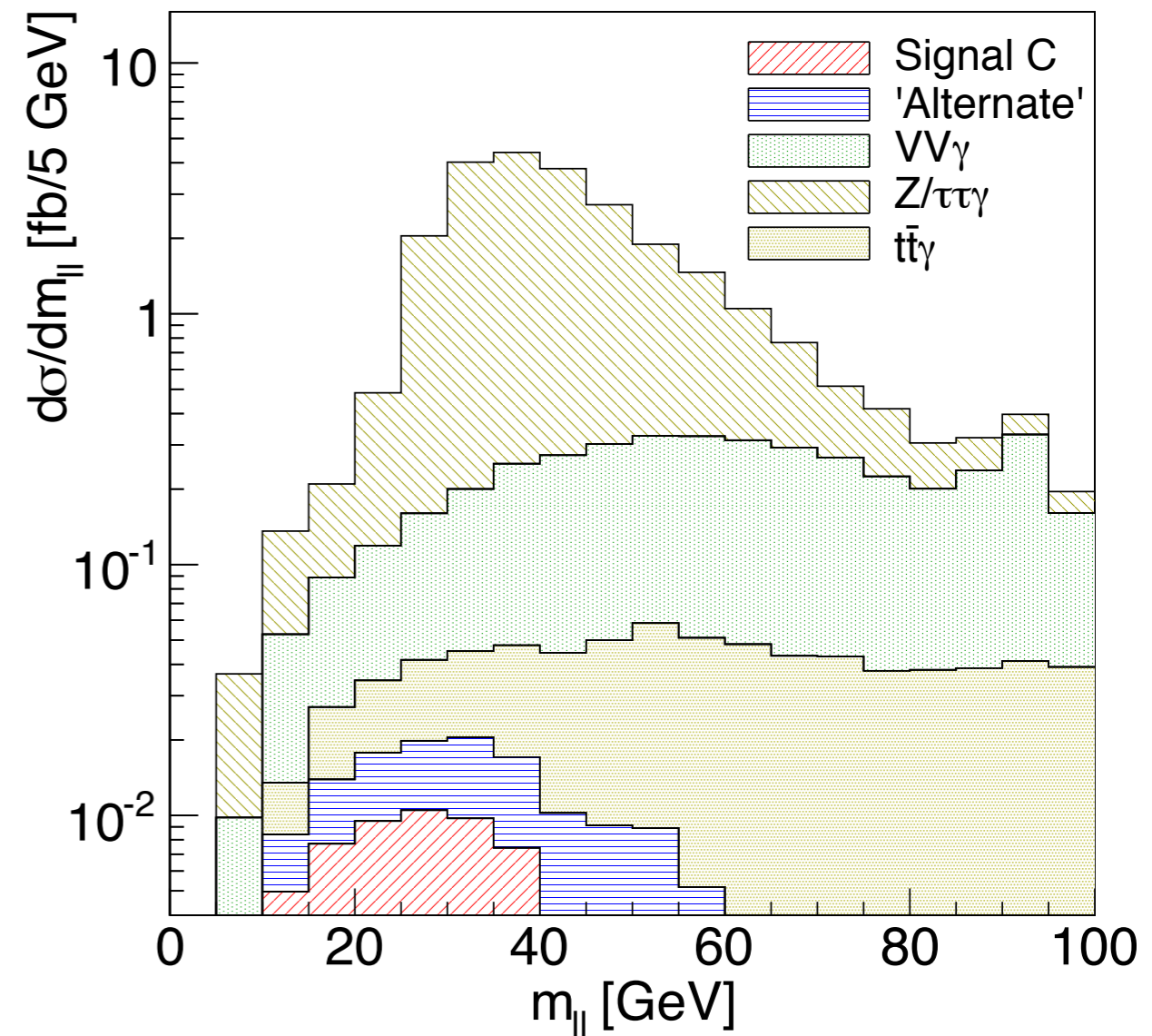
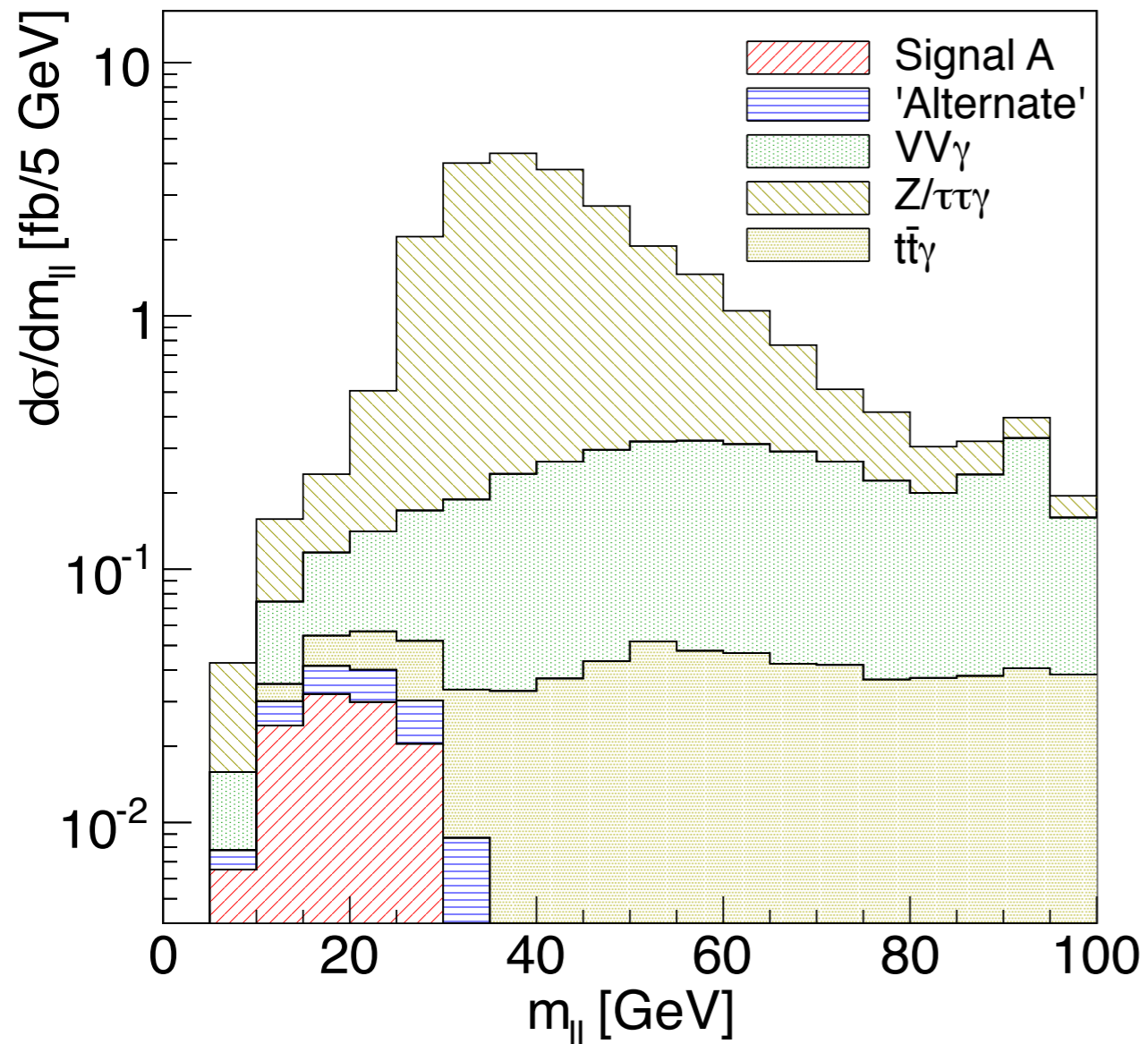
$$pp \rightarrow \chi_2^0 \chi_3^0 \rightarrow \gamma + \ell^+ \ell^- + \cancel{E}_T$$

Most discriminating cut is m_{ll}



$$pp \rightarrow \chi_2^0 \chi_3^0 \rightarrow \gamma + \ell^+ \ell^- + \cancel{E}_T$$

Most discriminating cut is m_{ll}



*Search cuts are optimal for smaller mass splittings**

$$pp \rightarrow \chi_2^0 \chi_3^0 \rightarrow \gamma + \ell^+ \ell^- + \cancel{E}_T$$

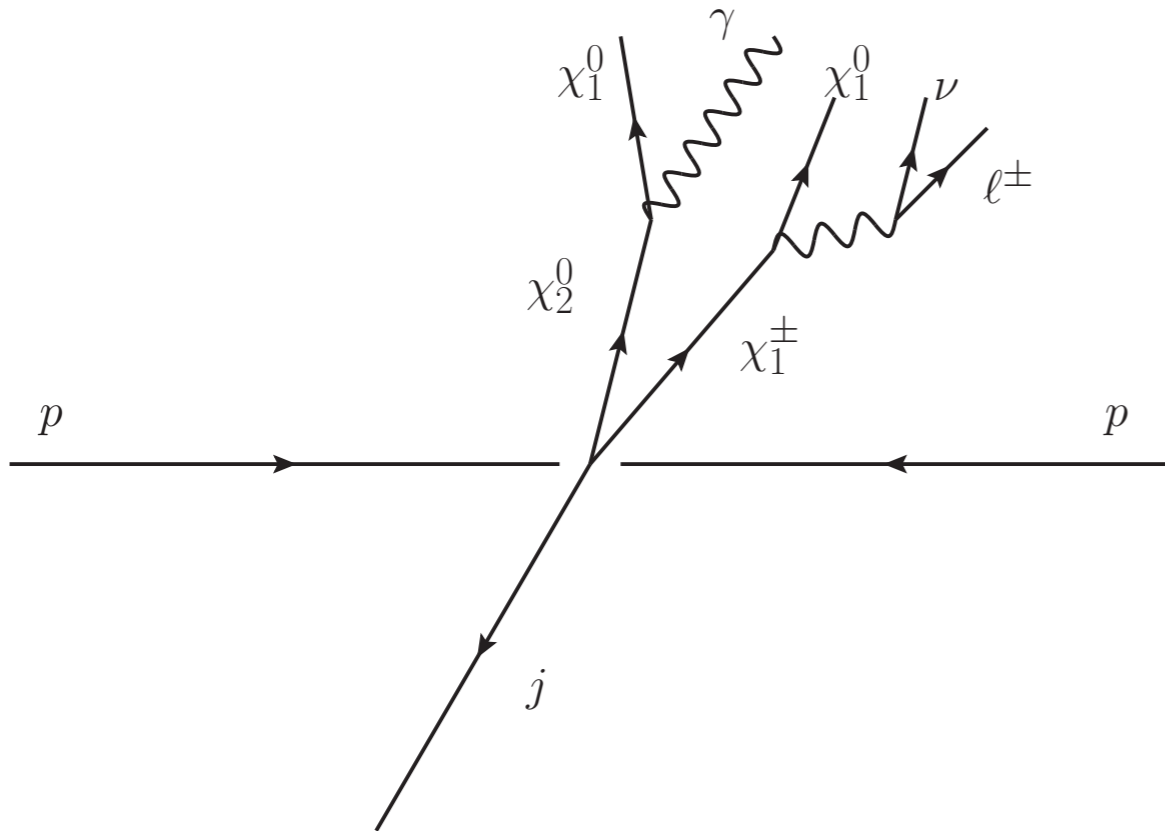
Benchmark points	Point A	Point C
μ	-150 GeV	-145 GeV
M_1	125 GeV	120 GeV
$\tan \beta$	2	10
$m_{\tilde{\chi}_1^0}$	124.0 GeV	105 GeV
$m_{\tilde{\chi}_2^0}$	156.9 GeV	150 GeV
$m_{\tilde{\chi}_3^0}$	157.4 GeV	163 GeV
$(\sqrt{s} = 14 \text{ TeV}) \int \mathcal{L} \text{ needed } [\text{fb}^{-1}]$	430	4300

Small mass splitting is good for cuts, bad for triggering
we were using 8-TeV di-lepton trigger, small efficiency

What if the system is boosted off a hard ISR jet to trigger on?

What if the system is boosted off a hard ISR jet to trigger on?

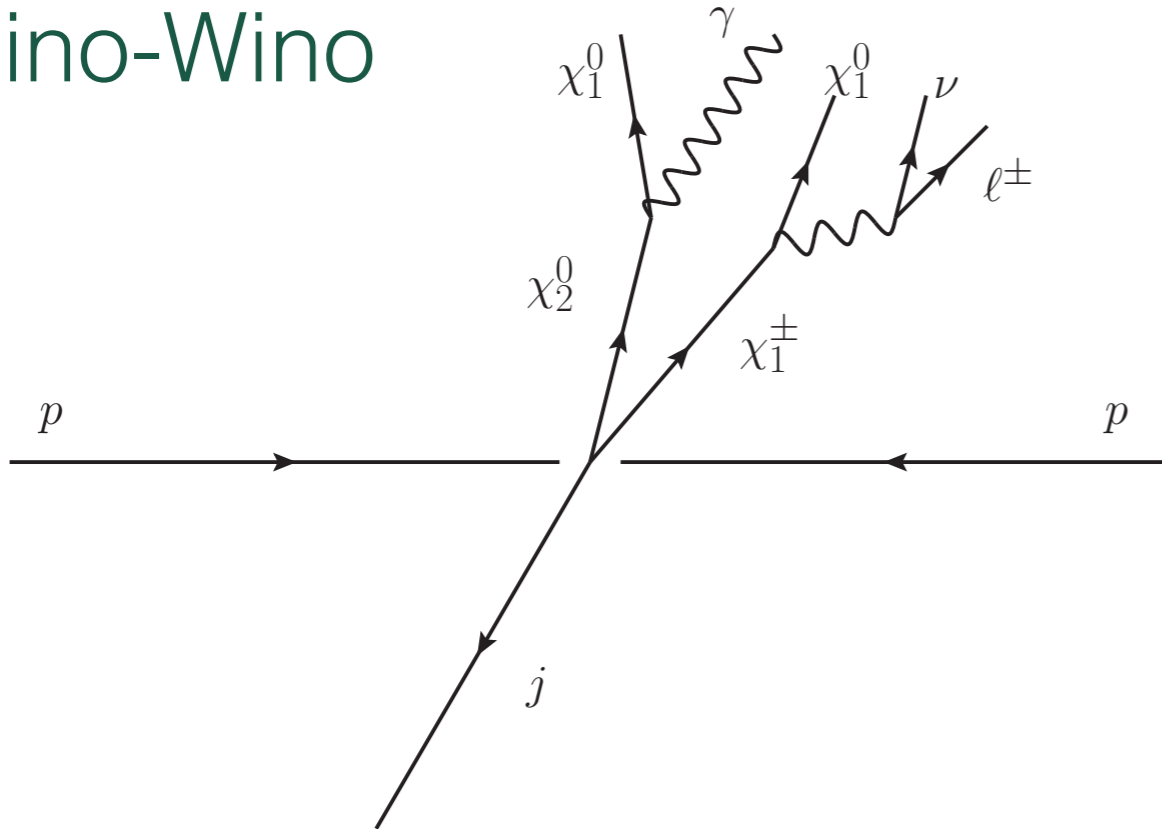
$$pp \rightarrow \chi_2^0 \chi_1^+ \rightarrow \gamma + \ell^+ + \cancel{E}_T$$



What if the system is boosted off a hard ISR jet to trigger on?

$$pp \rightarrow \chi_2^0 \chi_1^+ \rightarrow \gamma + \ell^+ + \cancel{E}_T$$

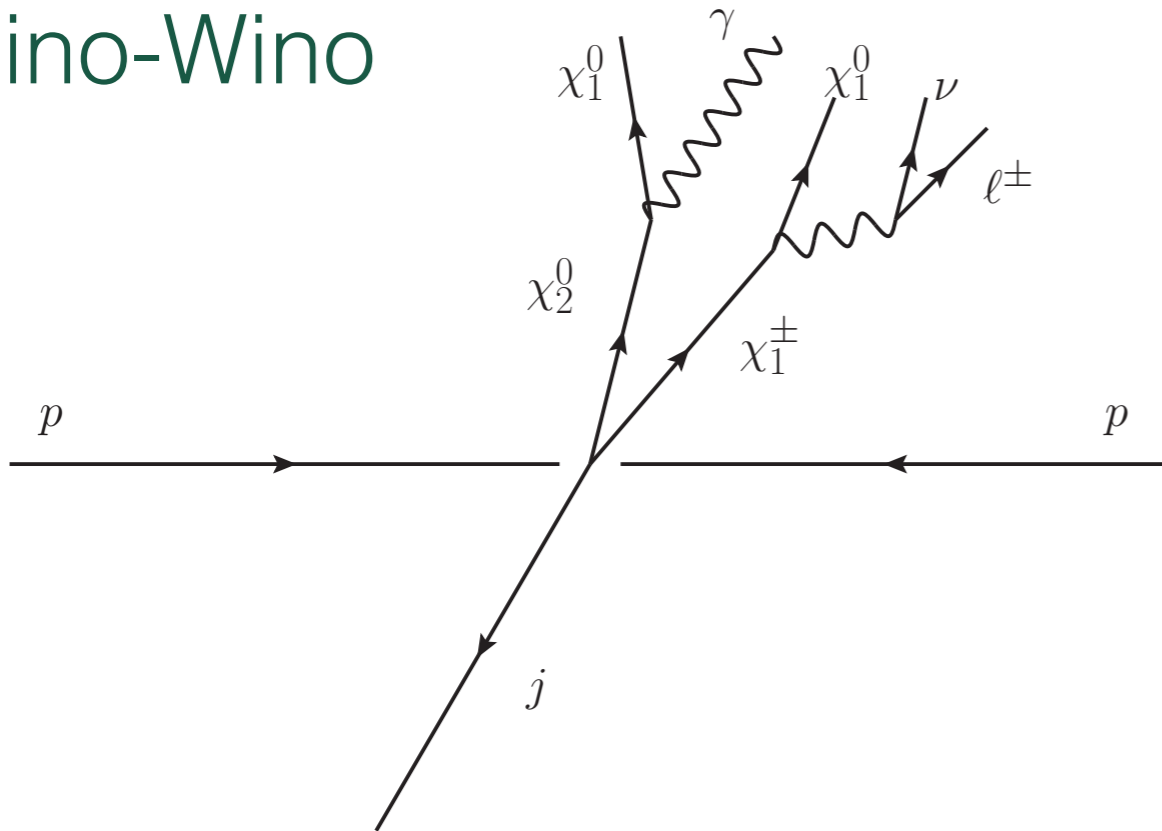
Bino-Wino



What if the system is boosted off a hard ISR jet to trigger on?

$$pp \rightarrow \chi_2^0 \chi_1^+ \rightarrow \gamma + \ell^+ + \cancel{E}_T$$

Bino-Wino

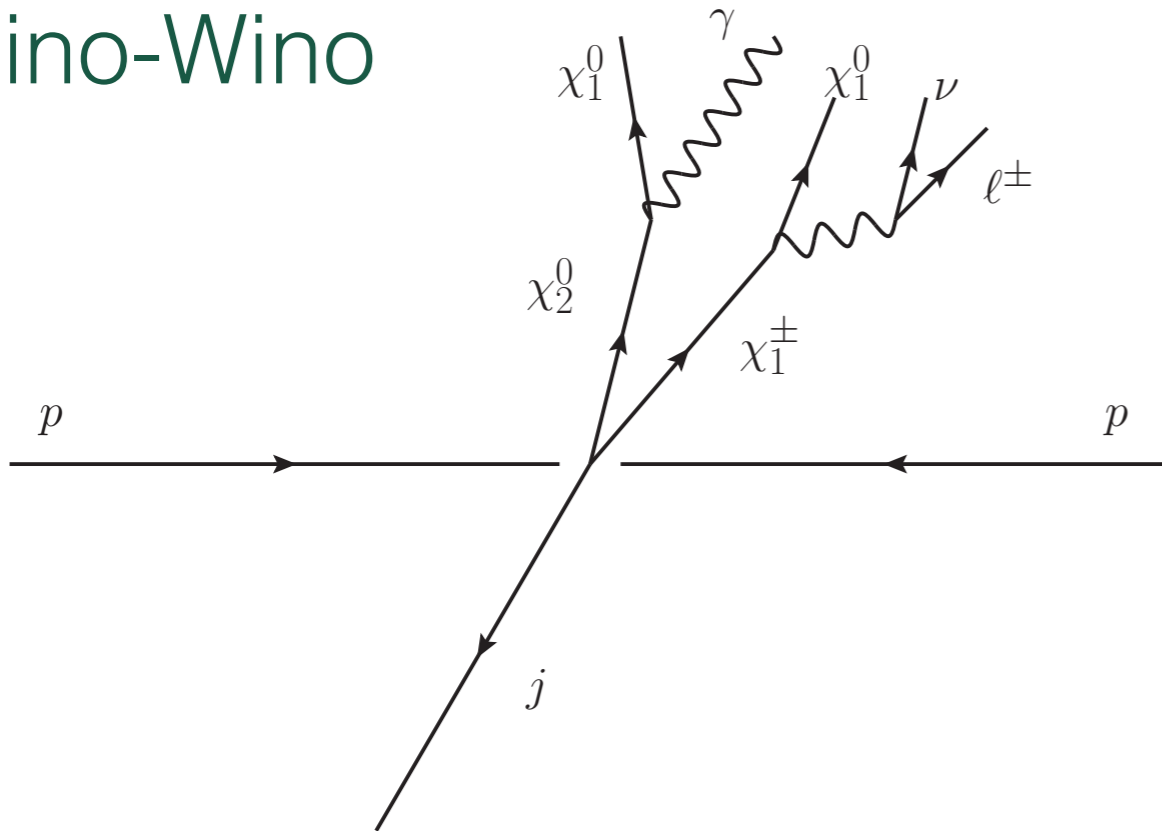


- $\cancel{E}_T \propto p_T(j)$
- $p_T(\gamma) \propto \Delta(m_{\chi_2^0}, m_{\chi_1^0})$
- $p_T(\ell) \propto \Delta(m_{\chi_1^\pm}, m_{\chi_1^0})$

What if the system is boosted off a hard ISR jet to trigger on?

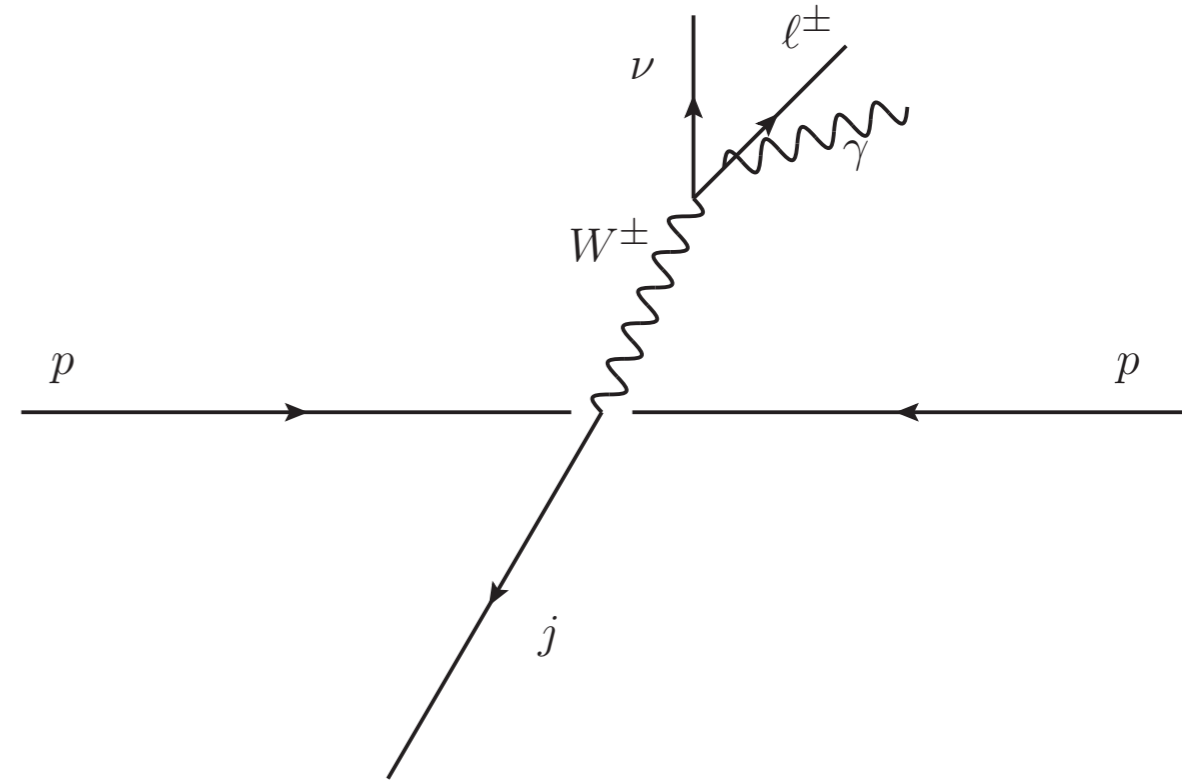
$$pp \rightarrow \chi_2^0 \chi_1^+ \rightarrow \gamma + \ell^+ + \cancel{E}_T$$

Bino-Wino



- $\cancel{E}_T \propto p_T(j)$
- $p_T(\gamma) \propto \Delta(m_{\chi_2^0}, m_{\chi_1^0})$
- $p_T(\ell) \propto \Delta(m_{\chi_1^\pm}, m_{\chi_1^0})$

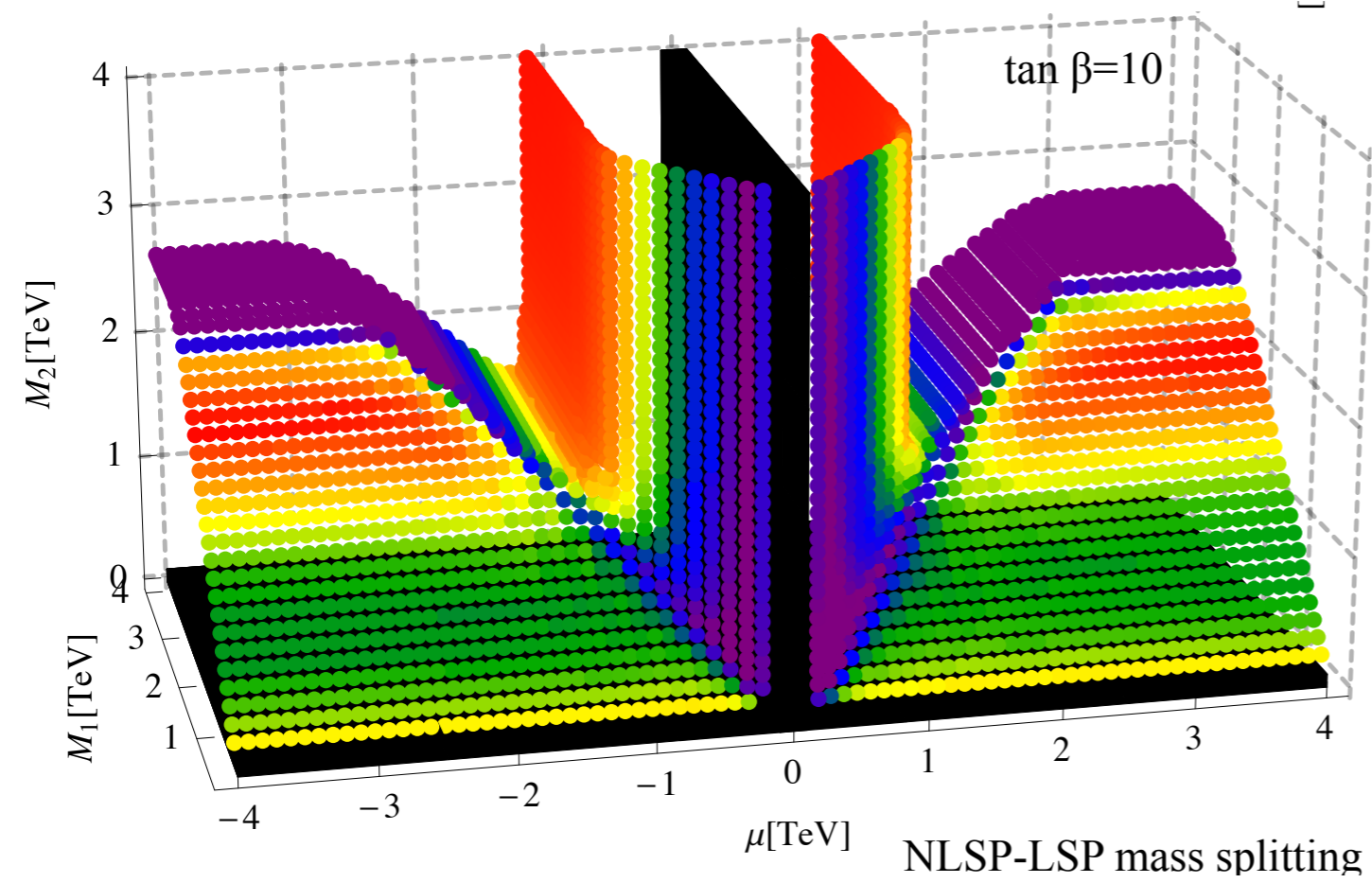
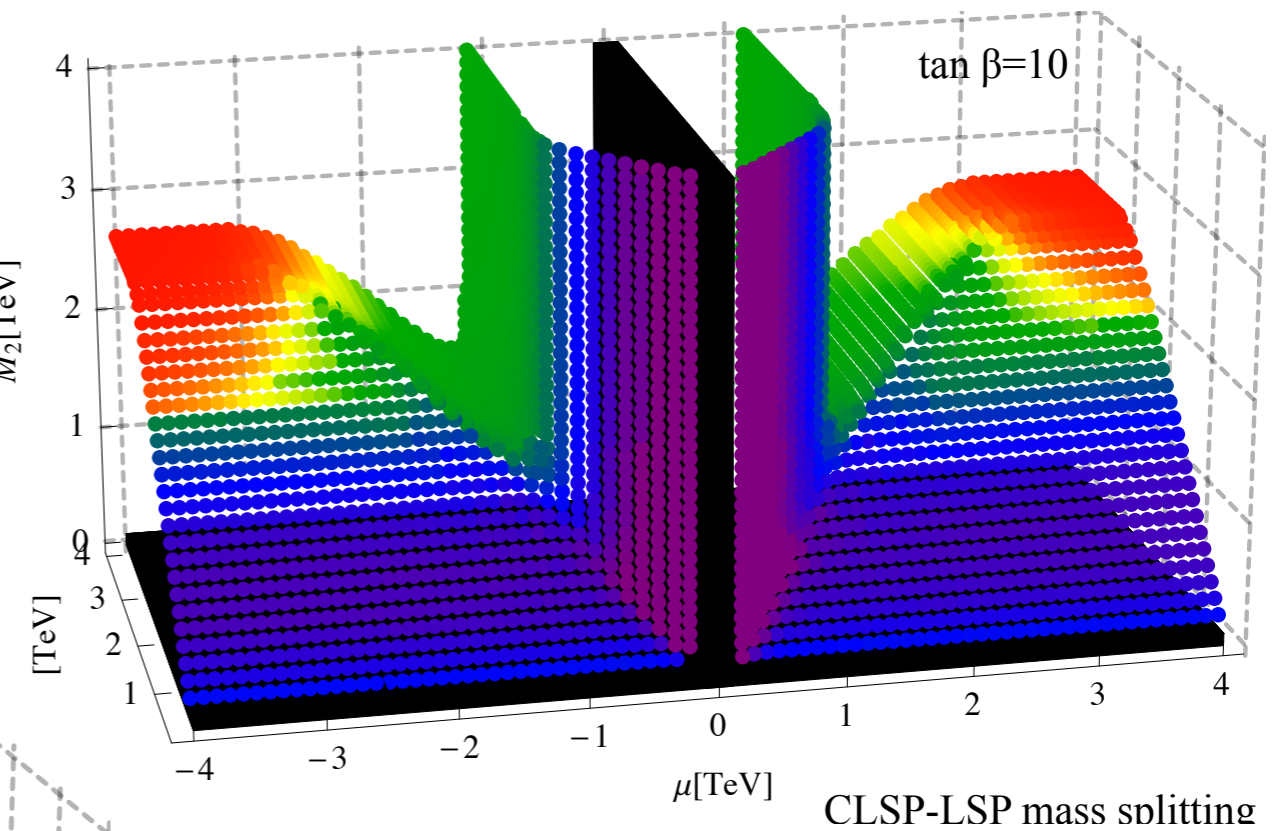
Dominant Background



- $\cancel{E}_T \propto p_T(j)$
- $p_T(\gamma) \propto p_T(j)$
- $p_T(\ell) \propto p_T(j)$

Parameter space for $pp \rightarrow \chi_2^0 \chi_1^\pm \rightarrow \gamma + \ell^+ + \cancel{E}_T$

CLSP-LSP mass splitting
 $m_{\chi_{1^\pm}} - m_{\chi_1^0} =$ ● <0.15 | ● 0.25 | ● 0.35 | ● 1 | ● 20 | ● >40 GeV



NLSP-LSP mass splitting
 $m_{\chi_2^0} - m_{\chi_1^0} =$ ● <1 | ● 10 | ● 20 | ● 30 | ● 40 | ● >60 GeV

$$pp \rightarrow \chi_2^0 \chi_1^+ \rightarrow \gamma + \ell^+ + \cancel{E}_T$$

Larger $p_{\tau}(j)$ yields more separation from background

$$pp \rightarrow \chi_2^0 \chi_1^+ \rightarrow \gamma + \ell^+ + \cancel{E}_T$$

Larger $p_{T(j)}$ yields more separation from background

Higher Energy Collider

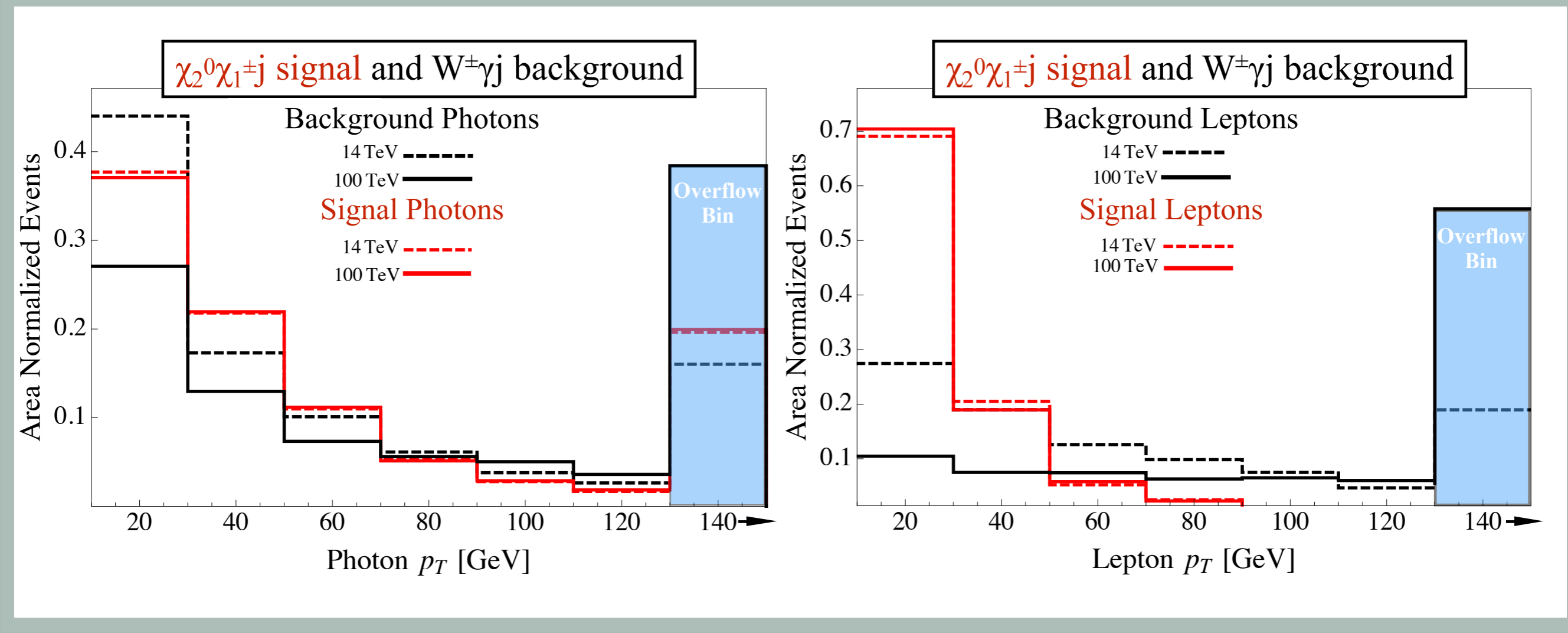
$$pp \rightarrow \chi_2^0 \chi_1^\pm \rightarrow \gamma + \ell^\pm + \cancel{E}_T$$

Larger $p_T(j)$ yields more separation from background

Higher Energy Collider

$$m_{\chi_2^0} = m_{\chi_1^\pm} = 200 \text{ GeV}; m_{\chi_1^0} = 190 \text{ GeV}$$

$$p_T(j) > 100(600) \text{ GeV}$$



$pp \rightarrow \chi_2^0 \chi_1^+ \rightarrow \gamma + \ell^+ + \cancel{E}_T$ Results

$$p_{T,j} > 0.8 \text{ TeV} \quad |\eta_j| < 2.5$$

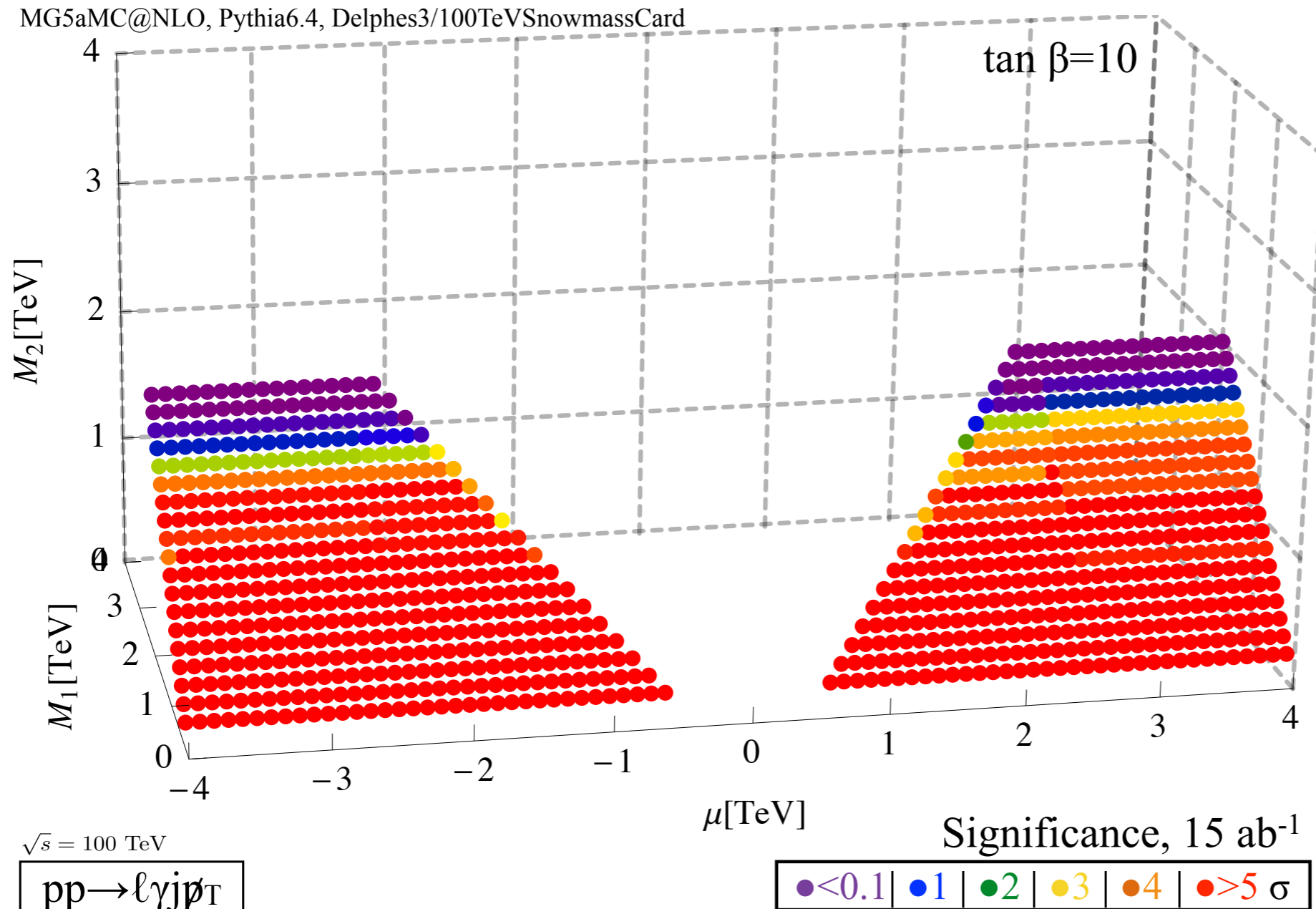
$$\cancel{E}_T > 1.2 \text{ TeV}$$

$$p_{T,\ell} = [10 - 60] \text{ GeV} \quad |\eta_\ell| < 2.5$$

$$p_{T,\gamma} = [10 - 60] \text{ GeV} \quad |\eta_\gamma| < 2.5$$

$$M_{T2}^{\gamma,\ell} < 10 \text{ GeV} \quad \Delta R_{\ell\gamma} > 0.5$$

MG5aMC@NLO, Pythia6.4, Delphes3/100TeVSnowmassCard



$pp \rightarrow \chi_2^0 \chi_1^+ \rightarrow \gamma + \ell^+ + \cancel{E}_T$ Results

$$p_{T,j} > 0.8 \text{ TeV} \quad |\eta_j| < 2.5$$

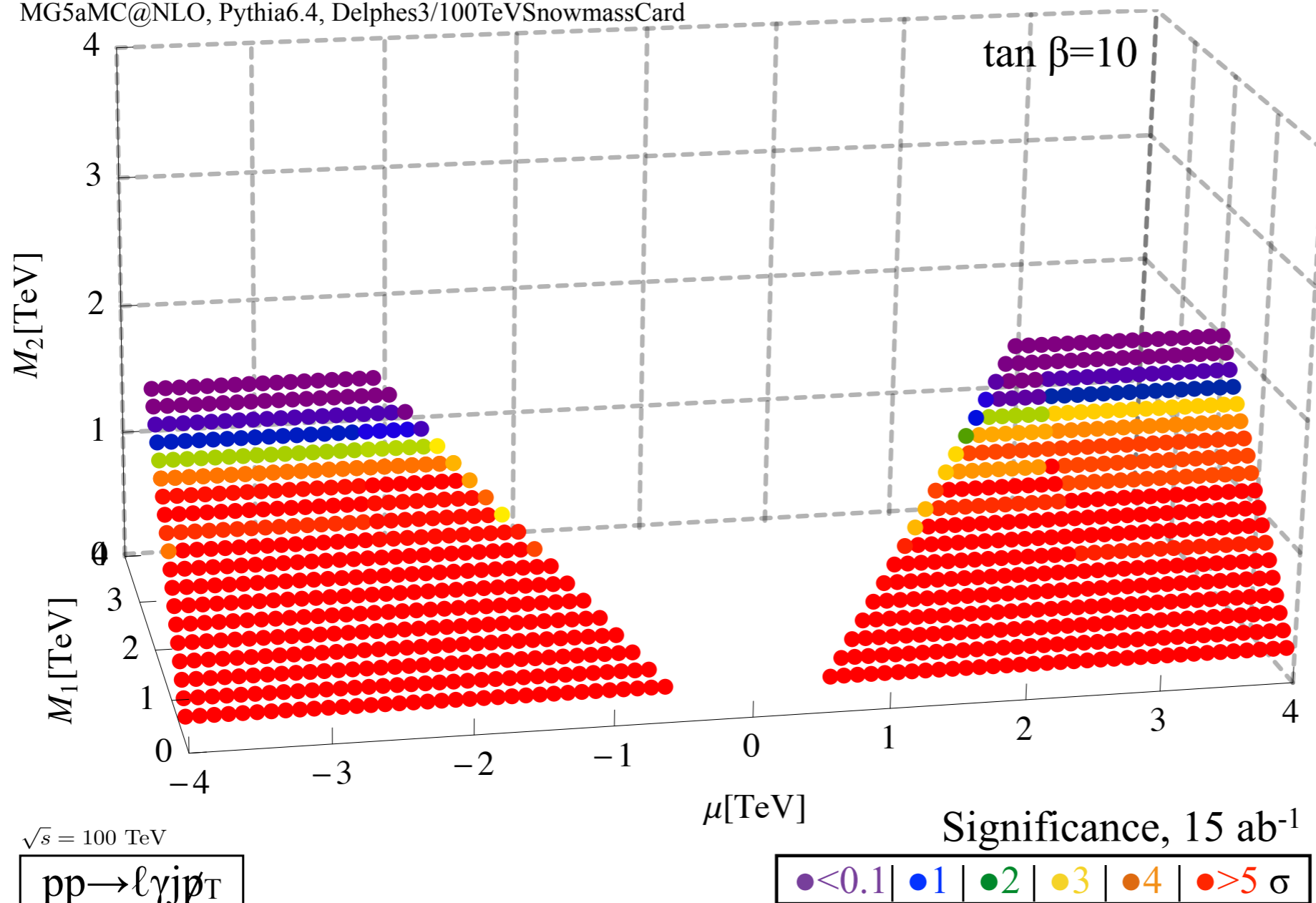
$$\cancel{E}_T > 1.2 \text{ TeV}$$

$$p_{T,\ell} = [10 - 60] \text{ GeV} \quad |\eta_\ell| < 2.5$$

$$p_{T,\gamma} = [10 - 60] \text{ GeV} \quad |\eta_\gamma| < 2.5$$

$$M_{T2}^{\gamma,\ell} < 10 \text{ GeV} \quad \Delta R_{\ell\gamma} > 0.5$$

MG5aMC@NLO, Pythia6.4, Delphes3/100TeVSnowmassCard



Cover most of the Bino-Wino plane!

Can entire surface be discovered with current/future experiments?

- Examine Direct/Indirect Detection constraints
- Explore search techniques at high energy colliders
 - ❑ Cover Bino-Wino surface
 - ❑ Resolve conflicts with pure wino
- Complementarity between experiments

Can entire surface be discovered with current/future experiments?

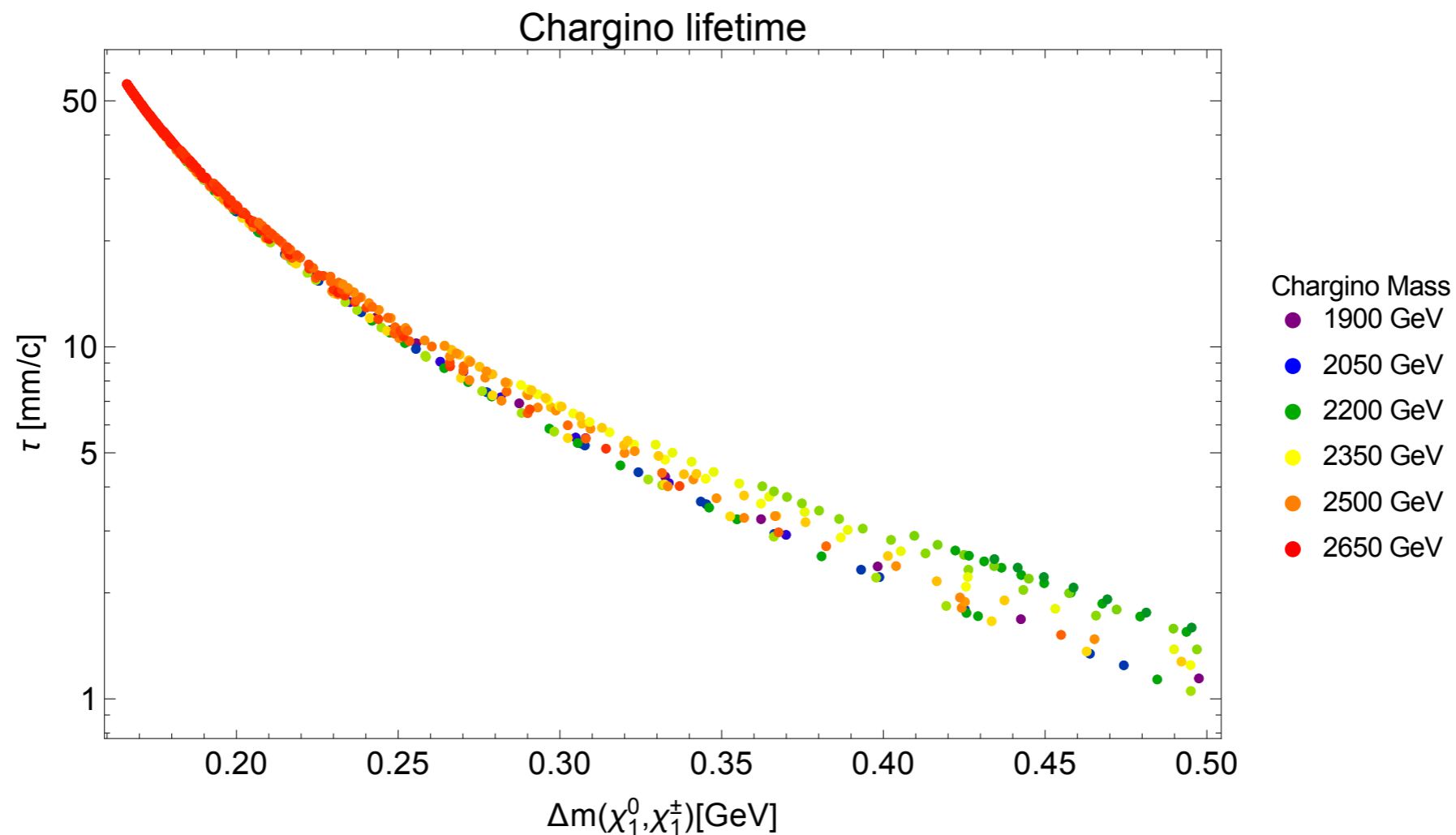
- Examine Direct/Indirect Detection constraints
- Explore search techniques at high energy colliders

- Cover Bino-Wino surface
- Resolve conflicts with pure wino

- Complementarity between experiments

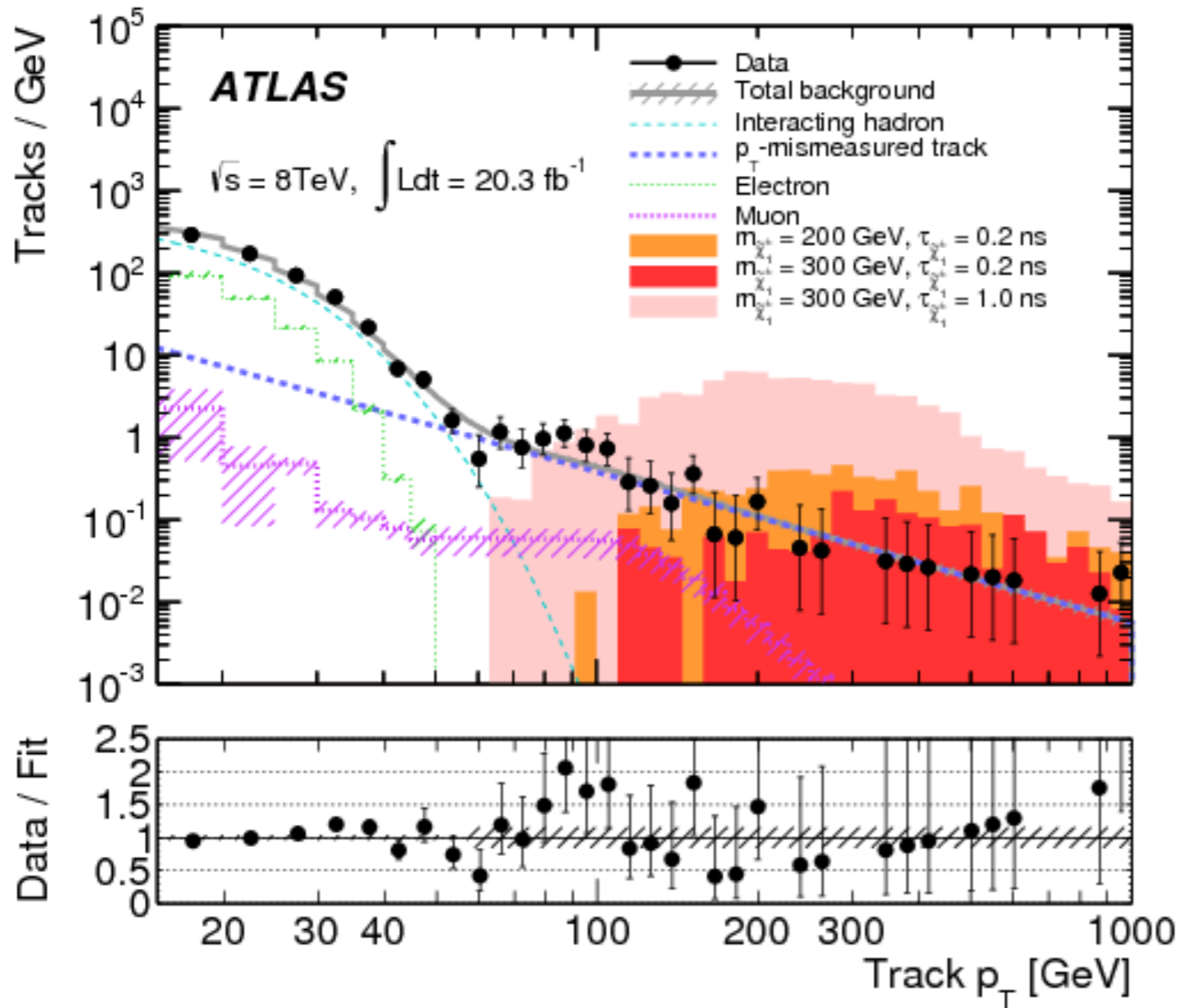
Resolve conflicts with pure wino

- Pure Wino mass splittings come at loop level
- Small mass splitting \rightarrow little phase space \rightarrow large lifetimes
- Look for chargino traveling a macroscopic distance then decaying to pions (soft and undetected). i.e. **disappearing track**



Disappearing Track

What process gives a background for a disappearing track?



- At large p_T dominated by p_T mismeasured tracks
- Fit by p_T^{-a} with a = 1.78 ± 0.05

Extrapolation

- Assume same shape
- Scale total background @ 8 TeV to ratio of $pp \rightarrow \nu\bar{\nu} + j$ between 100 TeV/8 TeV
- Same detector size

$$p_{T,j} > 90\text{ GeV}, \cancel{E}_T > 90\text{ GeV}, \Delta\phi_{min}^{\text{jet}-E_T^{\text{miss}}} > 1.5$$

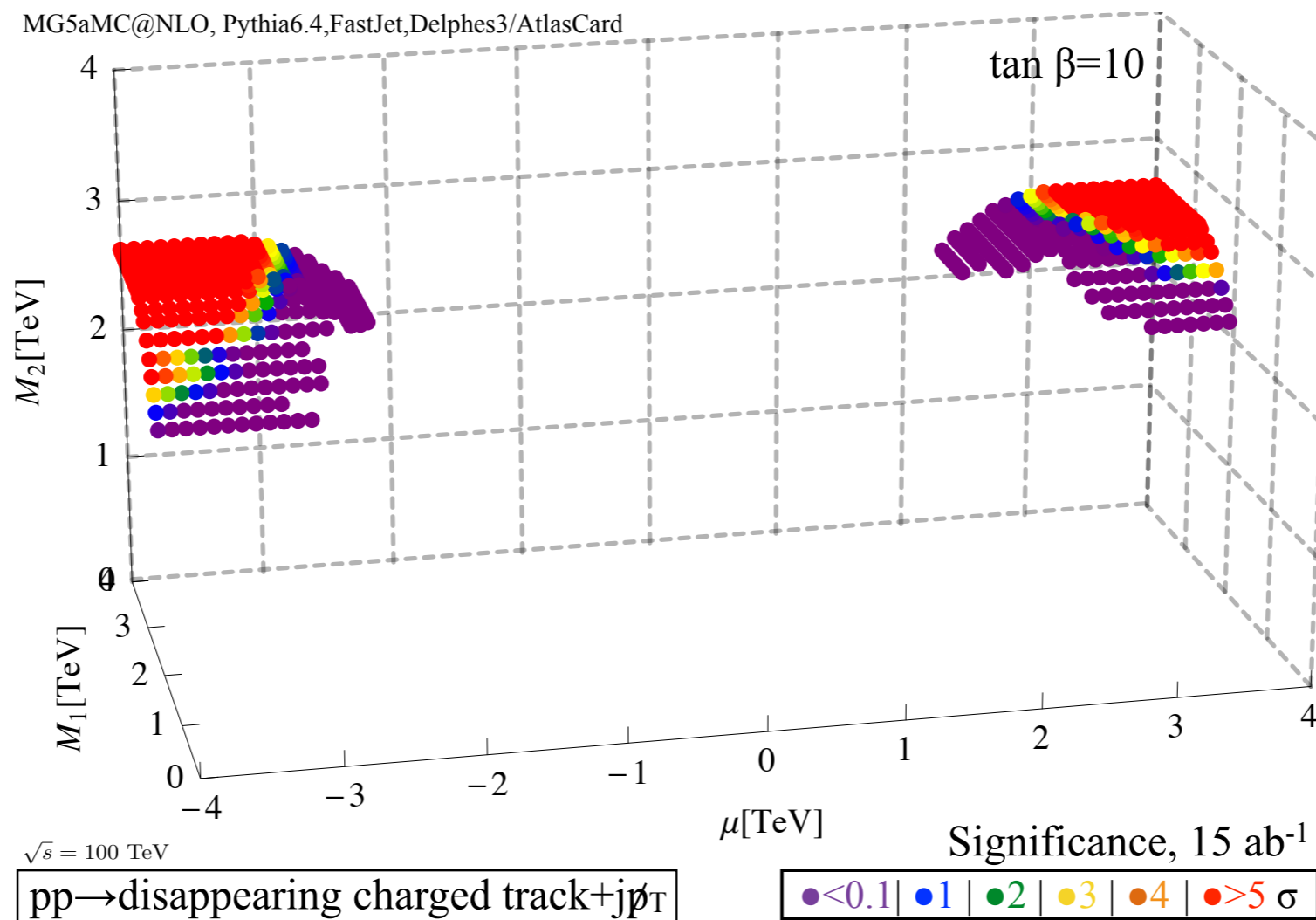
Isolated track (with largest p_T) with
 30 cm < L_T < 80 cm

Disappearing Track

$$p_{T,j} > 1 \text{ TeV}, \quad \cancel{E}_T > 1.4 \text{ TeV}, \quad \Delta\phi_{min}^{\text{jet}-E_T^{\text{miss}}} > 1.5$$

$$p_{T,j_2} > 500 \text{ GeV}, \quad p_{T,\text{track}} > 2.1 \text{ TeV}$$

Isolated track (with largest p_T) with $30 \text{ cm} < L_T < 80 \text{ cm}$



M. Low and L. T. Wang, JHEP **1408**, 161 (2014) [arXiv:1404.0682 [hep-ph]].

M. Cirelli, F. Sala and M. Taoso, JHEP **1410**, 033 (2014) [JHEP **1501**, 041 (2015)] [arXiv:1407.7058 [hep-ph]]

Can entire surface be discovered with current/future experiments?

- Examine Direct/Indirect Detection constraints
- Explore search techniques at high energy colliders

- Cover Bino-Wino surface
- Resolve conflicts with pure wino

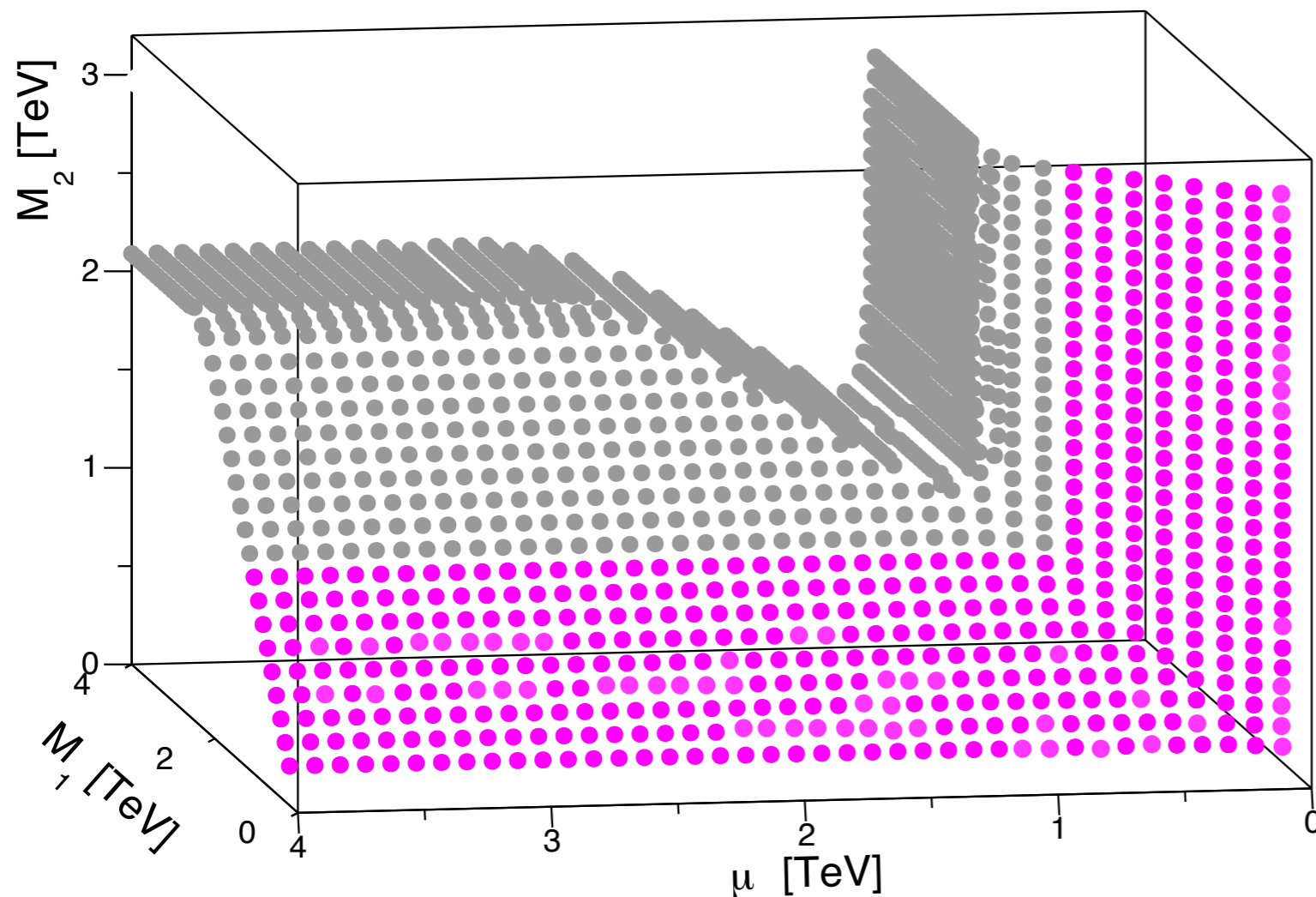
- Complementarity between experiments

More soft objects in energetic events

Extend methodology of the $j+\gamma+l^\pm$ search to soft dileptons

$$p_{T,\ell} = [10 - 50] \text{ GeV}, \quad p_{T,j} > 100 \text{ GeV},$$
$$m_{\ell\ell} < m_{\ell\ell}^{\text{max}}, \quad \cancel{E}_T > 500 \text{ GeV}$$

Relic neutralino 5σ discovery with soft dileptons (3 ab^{-1})

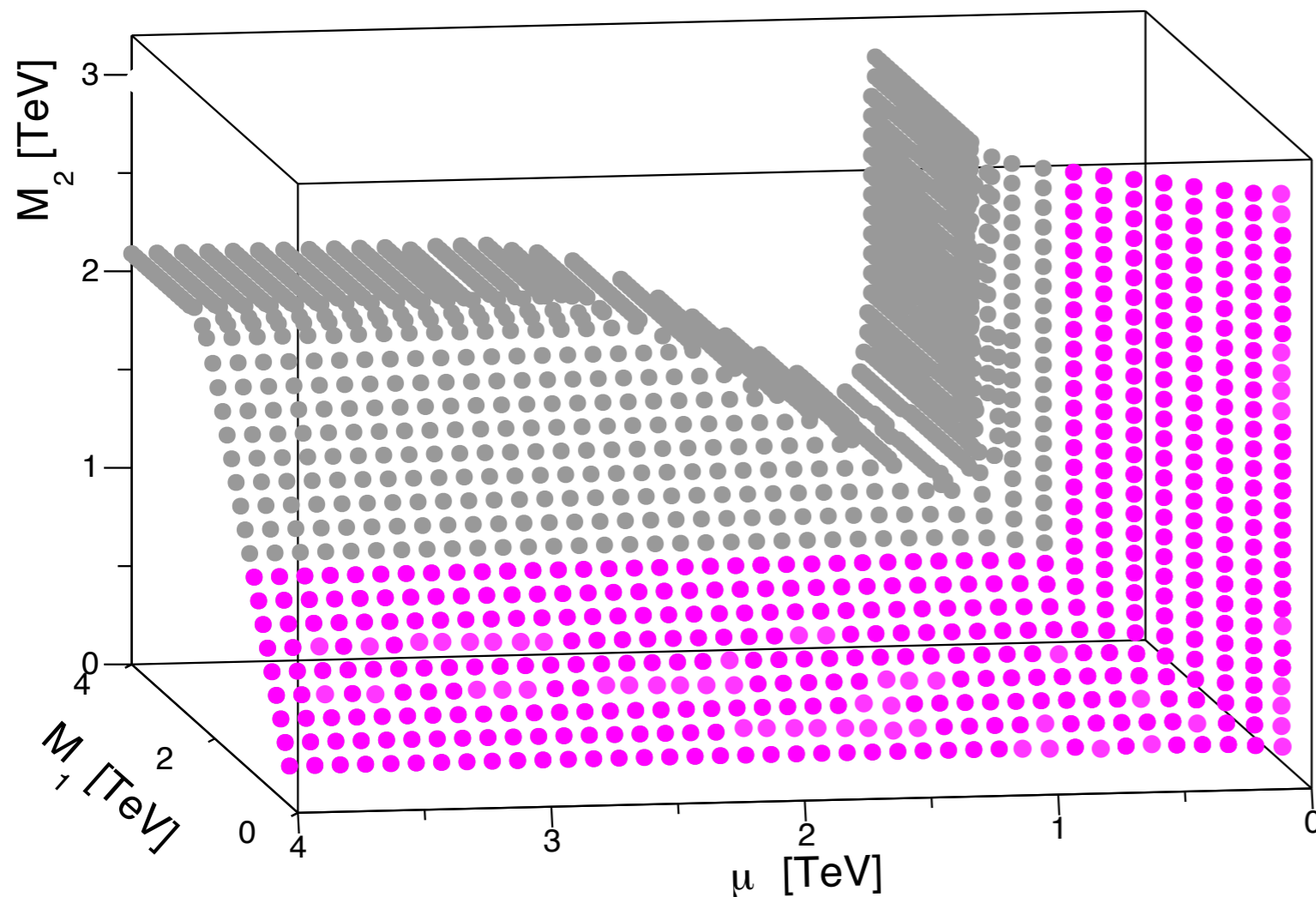


More soft objects in energetic events

Extend methodology of the $j+\gamma+l^\pm$ search to soft dileptons

$$p_{T,\ell} = [10 - 50] \text{ GeV}, \quad p_{T,j} > 100 \text{ GeV},$$
$$m_{\ell\ell} < m_{\ell\ell}^{\text{max}}, \quad \cancel{E}_T > 500 \text{ GeV}$$

Relic neutralino 5σ discovery with soft dileptons (3 ab^{-1})

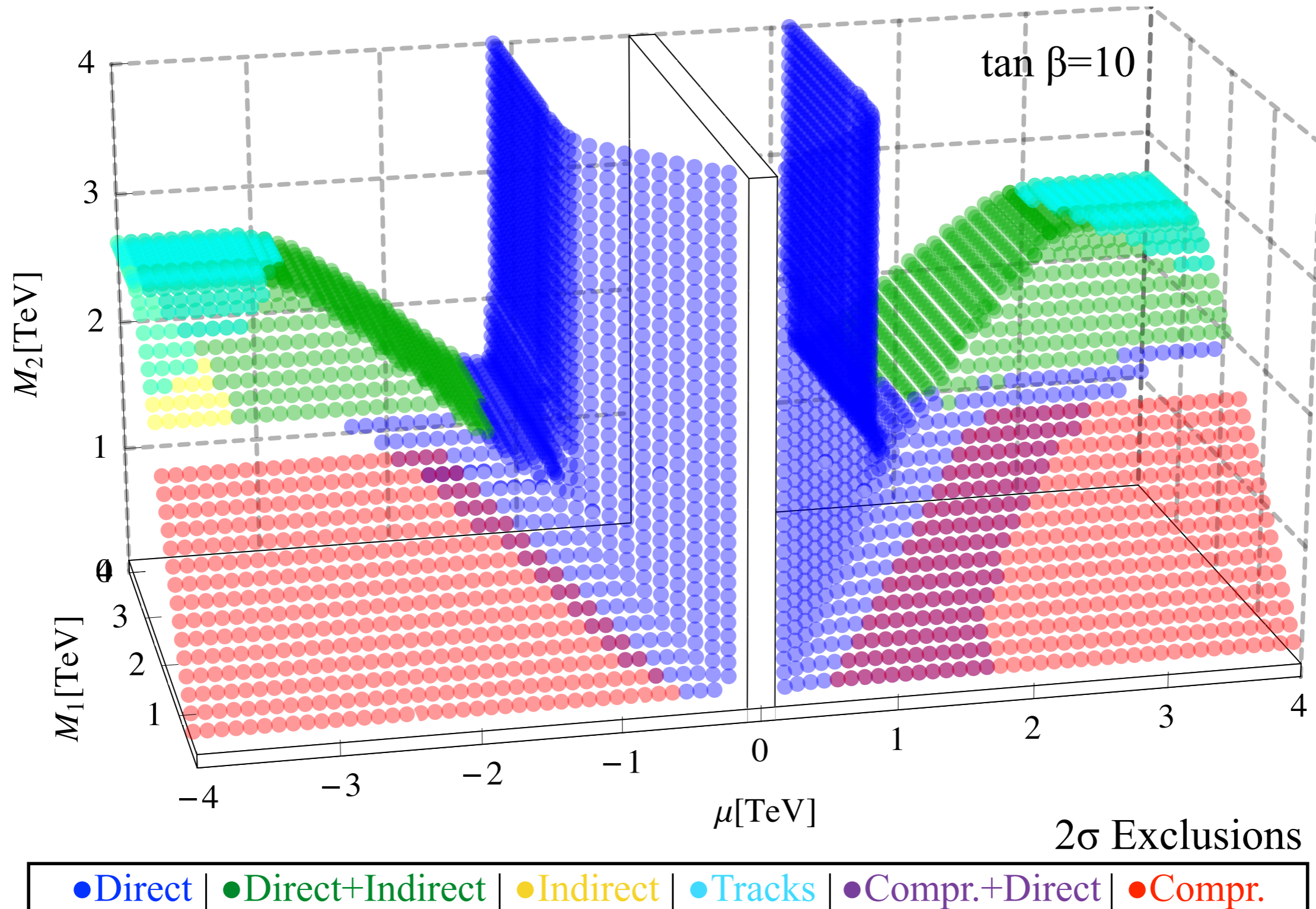


Soft dileptonic decay of χ_2^0 allows for discovery of much of Bino-Higgsino and some Bino-Wino

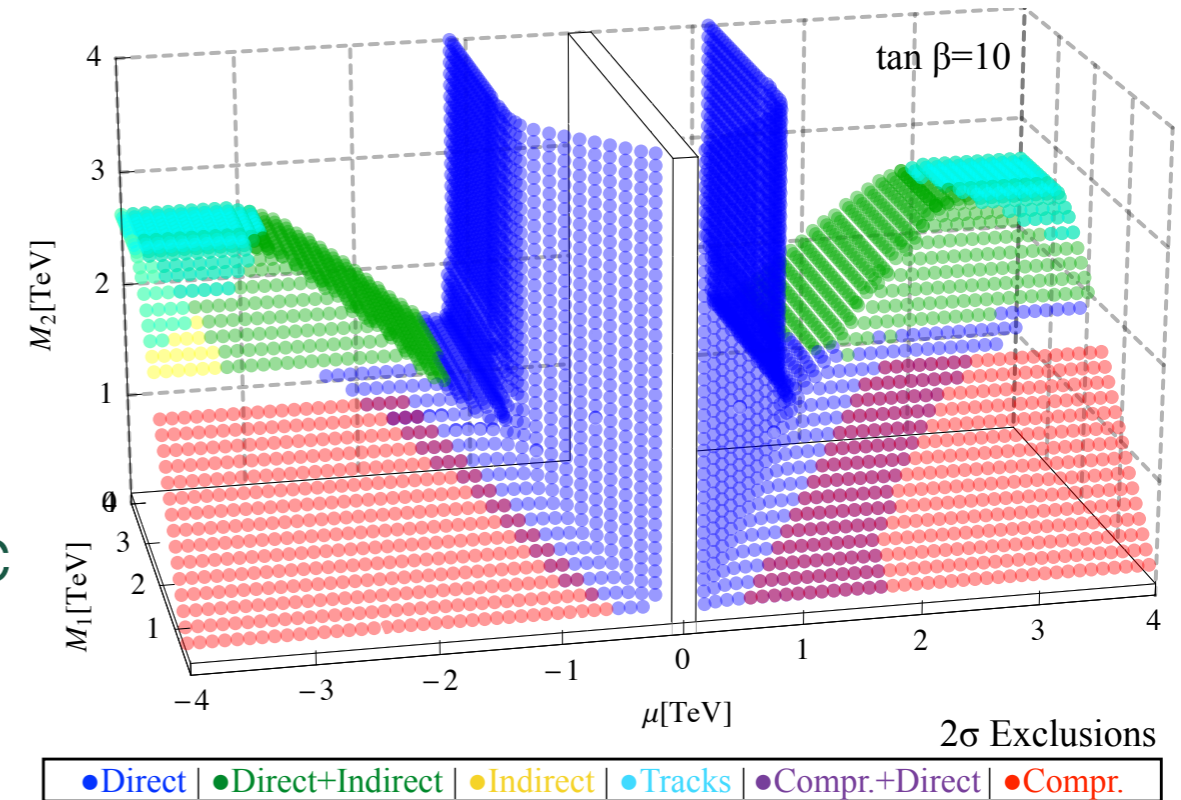
Can entire surface be discovered with current/future experiments?

- Examine Direct/Indirect Detection constraints
- Explore search techniques at high energy colliders
- Complementarity between experiments

Putting the searches together



- Well tempering uses co-annihilation of electroweakinos to set the observed relic abundance
 - Does not have to be supersymmetric
- Small mass splittings result
- Weak LHC limits (hard to trigger)
- Searching for soft objects recoiling off hard jet allow for collider coverage of most of the surface



- A few holes remain
- 'Pure Higgsino' needs much more decoupled Wino and Bino

- J. Bramante, N. Desai, P. Fox, A. Martin, BO and T. Plehn, “Towards the Final Word on Neutralino Dark Matter,” arXiv:1510.03460 [hep-ph].
- J. Bramante, P. J. Fox, A. Martin, BO, T. Plehn, T. Schell and M. Takeuchi, “Relic neutralino surface at a 100 TeV collider,” Phys. Rev. D **91**, 054015 (2015) doi:10.1103/PhysRevD.91.054015 [arXiv:1412.4789 [hep-ph]].
- J. Bramante, A. Delgado, F. Elahi, A. Martin and BO, “Catching sparks from well-forged neutrinos,” Phys. Rev. D **90**, no. 9, 095008 (2014) doi:10.1103/PhysRevD.90.095008 [arXiv:1408.6530 [hep-ph]].

Back Up

Benchmark points	Point A	Point B	Point C	Point D
μ	-150 GeV	-180 GeV	-145 GeV	150 GeV
M_1	125 GeV	160 GeV	120 GeV	125 GeV
$\tan \beta$	2	2	10	10
$m_{\tilde{\chi}_1^0}$	124.0 GeV	157 GeV	105 GeV	103 GeV
$m_{\tilde{\chi}_2^0}$	156.9 GeV	186 GeV	150 GeV	153 GeV
$m_{\tilde{\chi}_3^0}$	157.4 GeV	188 GeV	163 GeV	173 GeV
$\sigma(pp \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_3^0)$	394 fb	200 fb	345 fb	287 fb
$BR(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \gamma)$	0.0441	0.0028	0.0017	0.0014
$BR(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 l^+ l^-)$	0.0671	0.0712	0.0702	0.0700
$BR(\tilde{\chi}_3^0 \rightarrow \tilde{\chi}_1^0 \gamma)$	0.0024	0.0767	0.0115	0.0102
$BR(\tilde{\chi}_3^0 \rightarrow \tilde{\chi}_1^0 l^+ l^-)$	0.0714	0.0613	0.0447	0.0304
$\sigma(pp \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_3^0 \rightarrow \gamma l^+ l^- \tilde{\chi}_1^0 \tilde{\chi}_1^0)$	1.297 fb	1.125 fb	0.279 fb	0.205 fb

Back Up

- Scan over all possibilities for each cut (not $m_{\ell\ell}$)
- Pick cut which maximizes S/\sqrt{B}
- Repeat until no gain in significance or each cut has been used

'small mass splitting' cuts	Cross section [ab]					Significance S/B
	Signal A	Signal B	$VV\gamma$	$t\bar{t}\gamma$	$Z/\tau\tau\gamma$	
0) Basic Selection	281	169	5830	18900	24500	5.7×10^{-3} (3.4×10^{-3})
1) $N_{jets} = 0$	181	108	4820	1220	21400	6.6×10^{-3} (3.9×10^{-3})
2) $ \Delta\phi_{\ell_1, \ell_2} < 1.0$	118	79.5	580	201	567	8.8×10^{-2} (5.9×10^{-2})
3) $\left. \begin{array}{l} 15 \text{ GeV} < m_T(\ell_2) < 50 \text{ GeV} \\ m_T(\ell_1) < 60 \text{ GeV} \end{array} \right\}$	52.4	38.2	93.3	32.8	92.2	0.24 (0.17)
4) $ \Delta\phi_{\ell\ell-\gamma} > 1.45$	49.9	37.0	65.2	25.0	67.8	0.32 (0.23)
5) $30 \text{ GeV} < p_{T,\gamma} < 100 \text{ GeV}$	36.9	28.2	36.6	17.2	19.0	0.51 (0.39)
6) \cancel{E}_T cuts	26.8	20.2	24.6	3.90	0.00	0.94 (0.71)
7) $m_{\ell\ell} < 24 \text{ GeV}$	23.3	19.3	9.29	0.00	0.00	2.5 (2.1)

- Discover 'A' with 430 fb^{-1} (125 GeV DM particle)
- Discover 'B' with 620 fb^{-1} (157 GeV DM particle)

Back Up

- Points 'C' and 'D' have larger mass splittings
- Cuts not as effective
- More possibility of 'Alternative Signal'

'large mass splitting' cuts Cut	Cross section [ab]					Significance
	Signal C	Signal D	$VV\gamma$	$t\bar{t}\gamma$	$Z/\tau\tau\gamma$	S/B
0) Basic Selection	256	411	5830	18900	24500	5.2×10^{-3} (8.3×10^{-3})
1) $N_{jets} = 0$	157	227	4820	1220	21400	5.7×10^{-3} (8.3×10^{-3})
2) $ \Delta\phi_{\ell_1, \ell_2} < 1.05$	68.3	109	618	208	608	4.8×10^{-2} (7.6×10^{-2})
3) $\left. \begin{array}{l} 10 \text{ GeV} < m_T(\ell_1) < 100 \text{ GeV} \\ 10 \text{ GeV} < m_T(\ell_2) < 95 \text{ GeV} \end{array} \right\}$	47.9	72.2	389	127	117	7.5×10^{-2} (0.11)
4) $8 \text{ GeV} < \cancel{E}_T < 95 \text{ GeV}$	45.8	69.4	375	116	84.1	7.9×10^{-2} (0.12)
5) $m_{\ell\ell} < 39 \text{ GeV}$	42.8	64.0	228	35.9	51.5	0.14 (0.20)

- Discover 'C' with 4300 fb^{-1} (105 GeV DM particle)
- Discover 'D' with 1900 fb^{-1} (103 GeV DM particle)