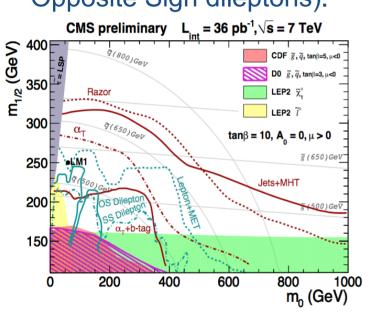
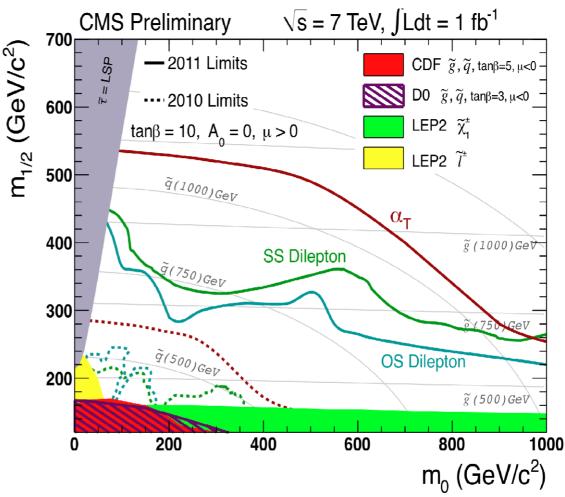
LFV, DM and LHC: how's SUSY health these days?

Antonio Masiero
Univ. of Padova and INFN, Padova

Progress on SUSY

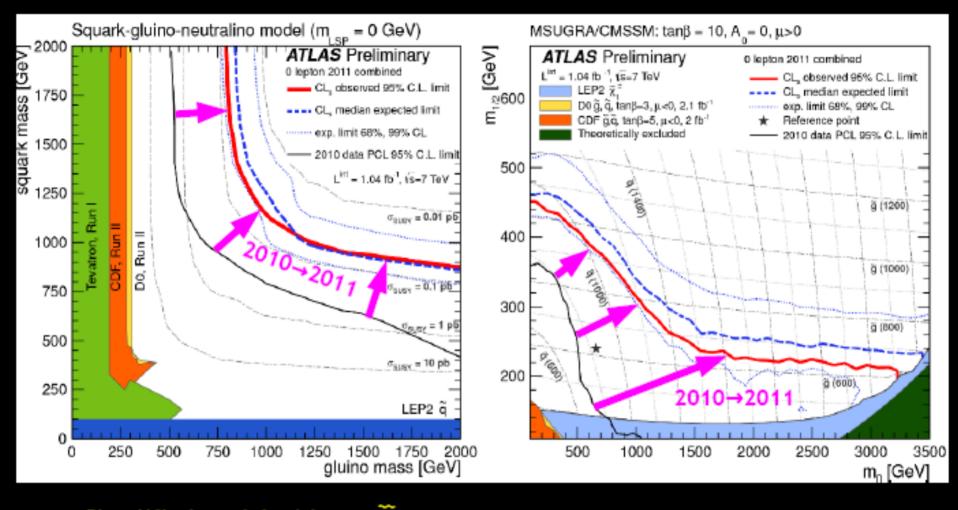
G. Tonelli EPS-HEP 2011 Results of the first three SUSY analyses completed on 2011 data (α_T , Same Sign and Opposite Sign dileptons).





Within the constrained SSM models we are crossing the border of excluding gluinos and squarks up to 1TeV and beyond. The air is getting thin for constrained SUSY. More conclusive results after summer.

SUSY in 0-lepton channel



Simplified model with two q generations, $m(\tilde{\chi}_{i}^{0})\sim 0$ $m_{\tilde{q}} > 800$ GeV $m_{\tilde{q}} > 850$ GeV Equal mass case: $m_{\tilde{g}} = m_{\tilde{g}} > 1.075$ TeV

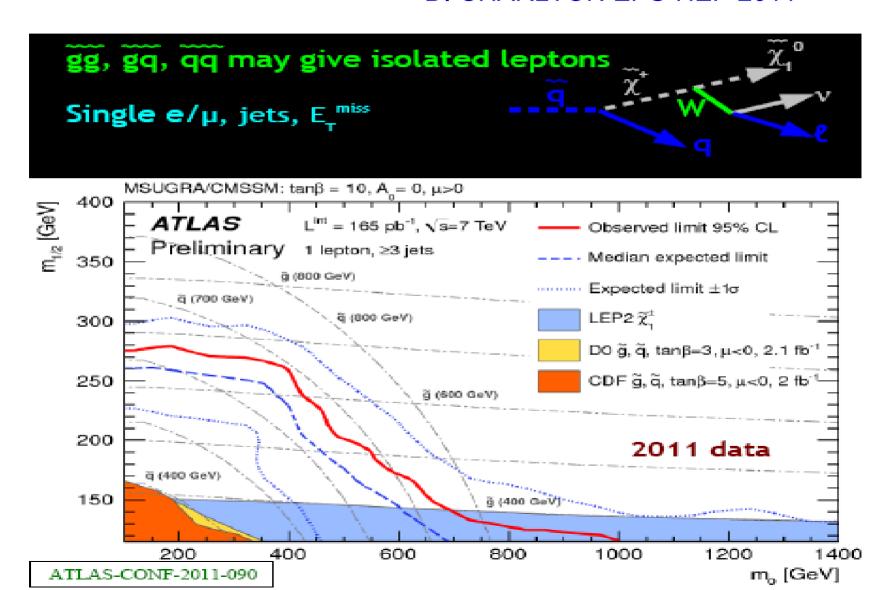
MSUGRA/CMSSM: tanB=10, A₀=0, µ>0

Equal mass case: m==m= > 980 GeV

D. CHARLTON EPS-HEP 2011

SUSY in 1-lepton channel

D. CHARLTON EPS-HEP 2011



ATLAS Searches* - 95% CL Lower Limits (EPS-HEP 2011)

Mass scale [TeV]

Impressive bounds on squarks and gluinos, into TeV range...

What do we learn? \rightarrow Papucci talk

1. Plain vanilla SUSY models (like MSSM with flavor-universal soft masses) are being pushed into a corner

but

Rychkov EPS-HEP 2011

2. Several other, theoretically motivated, scenarios remain very poorly constrained by existing searches

"Flavor-Split" spectra (heavy 1st-2nd gen squarks, gluino below 1-1.5 TeV, light 3rd gen)

"Squashed" spectra (everything below ~500GeV but splittings are small, O(10GeV)) Low MET
scenarios
(not necessarily
RPV)

WHY TO GO BEYOND THE SM

"OBSERVATIONAL" REASONS

•HIGH ENERGY PHYSICS

$$(NO)$$
 (but $A_{FB}...^{Z} \rightarrow bb$)

•FCNC, CP≠

NO (but CPV in Bs, sin2β tension...)

•HIGH PRECISION LOW-EN.

NO (but
$$(g-2)_{\mu}$$
 ...)

NEUTRINO PHYSICS

$$(YES)_y \neq 0, \theta_y \neq 0$$

•COSMO - PARTICLE PHYSICS

YES M, ΔB_{cosm} , INFLAT., DE)

THEORETICAL REASONS

•INTRINSIC INCONSISTENCY OF SM AS QFT

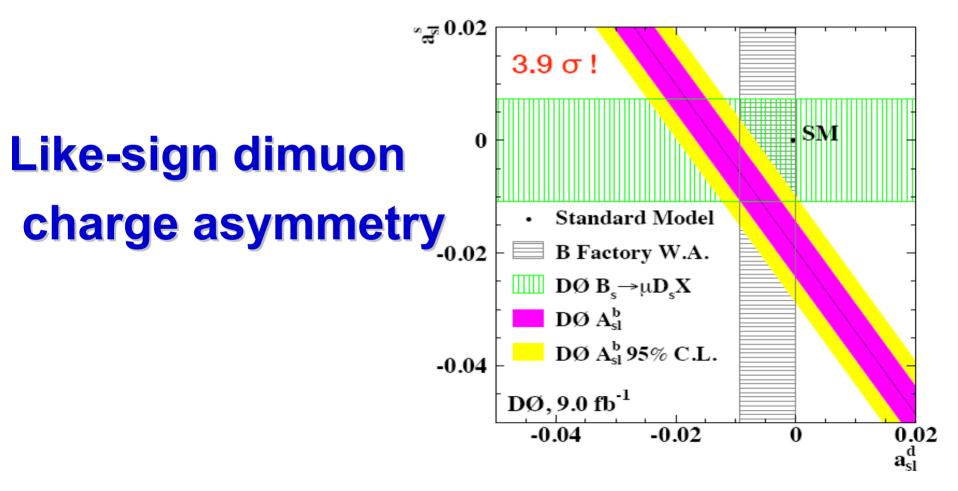
NO (spont. broken gauge theory without anomalies)

•NO ANSWER TO QUESTIONS THAT "WE" CONSIDER "FUNDAMENTAL" QUESTIONS TO BE ANSWERED BY "FUNDAMENTAL" THEORY

YES (hierarchy, unification, flavor)

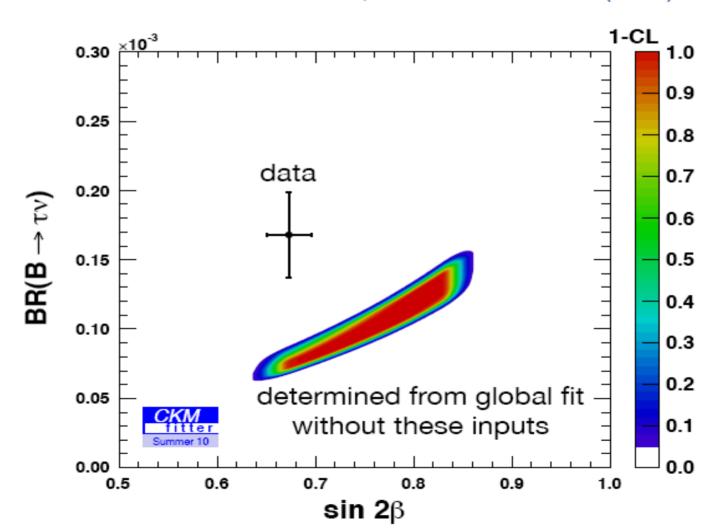
EVIDENCE OF NP ALONG THE HIGH INTENSITY ROAD?

"FLAVOR COLDS for the SM:



But *tension* in the UT fit even neglecting CPV in the B_s mixing

Lenz, Nierste + CKMfitter (2010)

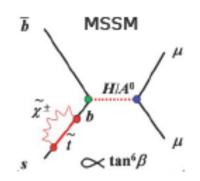


V_{ub} CRISIS

- discrepancies in the determinations of V_{ub} from inclusive semileptonic decays B→X_uIv, exclusive semileptonic decays B→πIv, and leptonic decay B→τν ("V_{ub} crisis")
- large difference of (14.4±2.9)% in the direct CP asymmetries measured in B⁰→K⁺π⁻ vs. B⁺→K⁺π⁰ decays, which is in conflict with the prediction of (2.2±2.4)% from QCD factorization ("B→Kπ puzzle")
- ▶ enhanced B_s→μ⁺μ⁻ branching ratio observed by CDF (but not by LHCb and CMS ⊕)

Rare decays B_{d,s}→µ⁺µ⁻

* interesting rare decays, which can be much enhanced in models with a warped extra dimension or SUSY models with large tanβ



Excess in B_s mode reported by CDF:

$$\mathcal{B}(B_s \to \mu^+ \mu^-) = (1.8^{+1.1}_{-0.9}) \cdot 10^{-8}$$
 SM: $(3.2 \pm 0.2) \cdot 10^{-9}$ SM: $(1.0 \pm 0.1) \cdot 10^{-10}$

Unfortunately no excess seen at LHCb (CMS):

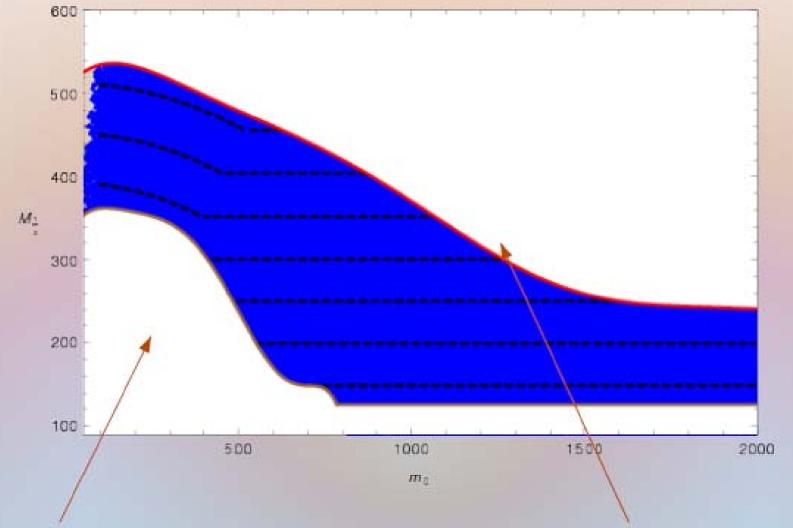
$$\mathcal{B}(B_s \to \mu^+ \mu^-) < 1.5 \ (1.9) \cdot 10^{-8}$$
 (at 95% CL)
 $\mathcal{B}(B_d \to \mu^+ \mu^-) < 5.2 \ (4.6) \cdot 10^{-9}$

NEUBERT EPS11

These bounds to not rule out the CDF result, but without refined LHC measurements the situation is inconclusive!

Relevant Parameter Space for 2 fb⁻¹

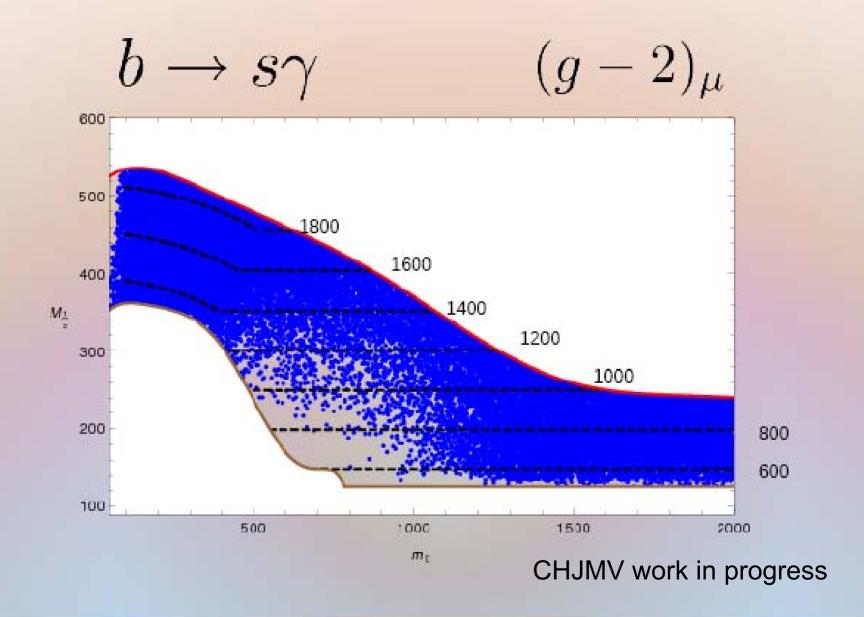
Jones at the EPS-HEP 2011 on the work in progress by Calibbi, Hodgkinson, Jones, A.M. and Vives



ATLAS Collaboration (1102.5290 [hep-ex])

Baer, Barger, Lessa, Tata (1004.3594 [hep-ph])

Flavour 3^o Constraints



The Role of $B_s \rightarrow \mu \mu$

LHCb with 2 fb⁻¹

IMPACT ON THE SUSY PARAMETER SPACE

- Exclusion of BR(B_s -> μ μ) down to 4x10⁻⁹, 95% C.L.
- 3σ evidence of BR(B_s -> μ μ) down to $5x10^{-9}$.
- 5σ discovery of BR(B_s -> $\mu \mu$) down to $9x10^{-9}$.

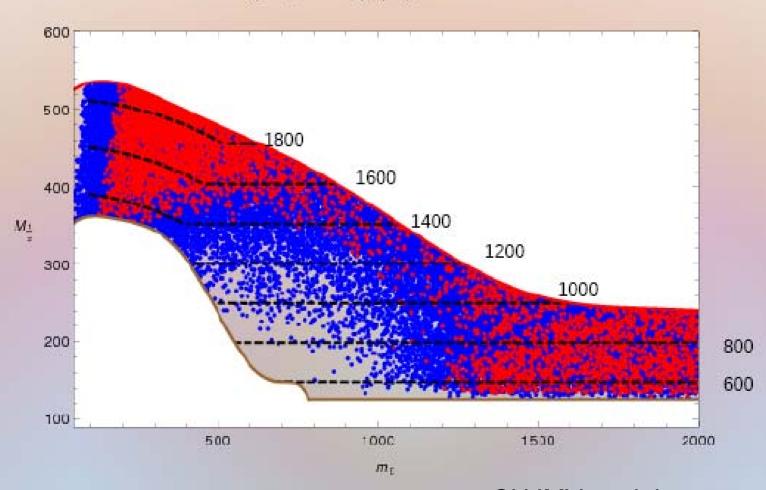
R. Lambert @ Moriond

CDF with 7 fb⁻¹

- BR(B_s ->
$$\mu \mu$$
) = (1.8 ± 1)x10⁻⁸

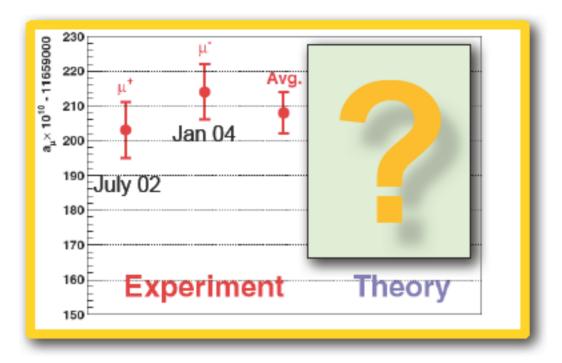
Exclusion due to $B_s \rightarrow \mu \mu$

BR
$$(B_s \to \mu \mu) < 4 \times 10^{-9}$$



CHJMV work in progress

The muon g-2: the experimental result



- **Today**: $a_{\mu}^{EXP} = (116592089 \pm 54_{stat} \pm 33_{sys})x10^{-11}[0.5ppm]$.
- Future: new muon g-2 experiments proposed at:
 - Fermilab (P989), aiming at 0.14ppm STAGE-1 APPROVAL!!
 - J-PARC aiming at 0.1 ppm
 - [D. Hetzog & N. Saito, U.Paris, Feb 2010; B. Lee Roberts & T. Mibe, Tau2010]
- Are theorists ready for this (amazing) precision? [not yet]

The muon g-2: Standard Model vs. Experiment

Adding up all contributions, we get the following SM predictions and comparisons with the measured value:

$$a_{\mu}^{EXP}$$
 = 116592089 (63) × 10⁻¹¹

E821 – Final Report: PRD73 (2006) 072 with latest value of $\lambda = \mu_{\mu}/\mu_{p}$ (CODATA'06)

$a_{\mu}^{\scriptscriptstyle \mathrm{SM}} imes 10^{11}$	$(\Delta a_{\mu} = a_{\mu}^{\text{EXP}} - a_{\mu}^{\text{SM}}) \times 10^{11}$	σ
[1] 116 591 782 (59)	307 (86)	3.6
[2] 116 591 802 (49)	287 (80)	3.6
[3] 116 591 830 (52)	259 (82)	3.2
[4] 116 591 894 (54)	195 (83)	2.4

with
$$a_u^{HHO}(IbI) = 105 (26) \times 10^{-11}$$

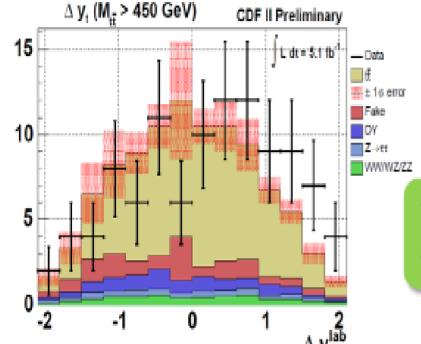
- F. Jegerlehner, A. Nyffeler, Phys. Rept. 477 (2009) 1
- [2] Davier et al, arXiv:1010.4180, Oct 2010 (includes BaBar and KLOE10 2π)
- [3] HLMNT10: Hagiwara et al, Tau 2010, Sep. 2010 (incl BaBar and KLOE10 2π)
- [4] Davier et al, arXiv:1010.4180, Oct 2010, ⊤ data.

Note that the th. error is now about the same as the exp. one

Top anti-Top asymmetry



5.1 fb⁻¹



Events

CDF public note 10436

$$A_{fb} = 0.42 \pm (0.15)^{stat} \pm (0.05)^{syst}$$
 (dilepton final state)

- \checkmark 2.3 σ from the SM prediction.
- \checkmark 3.4 σ in the l+jets topology.

DUPERRIN EPS-HEP 2011

- ✓ axigluons, diquarks, new weak bosons, EDs etc...
- ✓ Or gluon radiations modeling at NLO?

Is it possible that there is "only" a light higgs boson and no NP?

- This is acceptable if one argues that no ultraviolet completion of the SM is needed at the TeV scale simply because there is no actual fine-tuning related to the higgs mass stabilization (the correct value of the higgs mass is "environmentally" selected). This explanation is similar to the one adopted for the cosmological constant
- Barring such wayout, one is lead to have TeV NP to ensure the unitarity of the elw. theory at the TeV scale

% FINE-TUNING FOR THE NEW PHYSICS AT THE ELW. SCALE

- Elementary Higgs →In the MSSM % fine-tuning among the SUSY param. to avoid light SUSY particles which would have been already seen at LEP and Tevatron and now also at LHC
- Elementary Higgs \rightarrow PSEUDO-GOLDSTONE boson in the LITTLE HIGGS model \rightarrow Λ^2 div. cancelled by new colored fermions, new W,Z, γ , 2Higgs doublets... \rightarrow % fine-tuning to avoid too large elw. Corrections
- COMPOSITE HIGGS in a 5-dim. holographic theory (Higgs is a PSEUDO-GOLDSTONE boson and the elw. symmetry breaking is triggered by bulk effects (in 5 dim. the theory is WEAKLY coupled, but in 4 dim. the bulk looks like a STRONGLY coupled sector) → also here % fine-tuning needed to survive the elw. precision tests

The Energy Scale from the "Observational" New Physics

neutrino masses dark matter baryogenesis inflation

NO NEED FOR THE NP SCALE TO BE **CLOSE TO THE ELW. SCALE**

The Ener Scale from the "Theoretical" New Physics

 $A \rightarrow A$ Stabilization of the electroweak symmetry breaking at M_W calls for an ULTRAVIOLET COMPLETION of the SM

already at the TeV scale





CORRECT GRAND UNIFICATION "CALLS" FOR NEW PARTICLES THE ELW. SCALE

THE DM ROAD TO NEW PHYSICS BEYOND THE SM: IS DM A PARTICLE OF THE NEW PHYSICS AT THE ELECTROWEAK ENERGY SCALE?

CONNECTION DM – ELW. SCALE THE WIMP MIRACLE: STABLE ELW. SCALE WIMPs

1) ENLARGEMENT OF THE SM

SUSY $(\mathbf{X}^{\mu}, \theta)$

EXTRA DIM.

LITTLE HIGGS.

 $(\mathbf{X}^{\mu}, \mathbf{j}^{i})$

SM part + new part

Anticomm. Coord.

New bosonic Coord.

to cancel Λ^2 at 1-Loop

2) **SELECTION** RULE

R-PARITY LSP

KK-PARITY LKP

T-PARITY LTI

→DISCRETE SYMM.

→ STABLE NEW PART.

> 3) FIND REGION (S) PARAM. SPACE WHERE THE "L" NEW PART. IS NEUTRAL + $\Omega_{\rm I}$ h^2 OK

Neutralino spin 1/2

 m_{ISP}

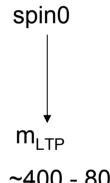
~100 - 200

GeV *

spin1 $m_{\text{\tiny I_KP}}$

~600 - 800

Ge\/



~400 - 800

GeV

But abandoning gaugino-masss unif. → Possible to have m_{ISP} down to 7 GeV

IS THE "WIMP MIRACLE" AN ACTUAL MIRACLE?

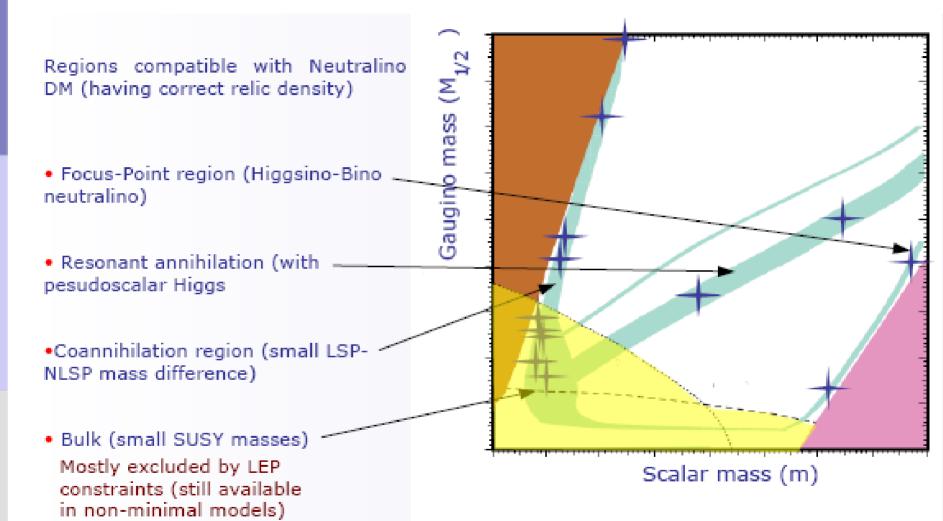
USUAL STATEMENT

Many possibilities for DM candidates, but WIMPs are really special: peculiar coincidence between particle physics and cosmology parameters to provide a VIABLE DM CANDIDATE AT THE ELW. SCALE

HOWEVER

when it comes to quantitatively reproduce the precisely determined DM density → once again the fine-tuning threat...





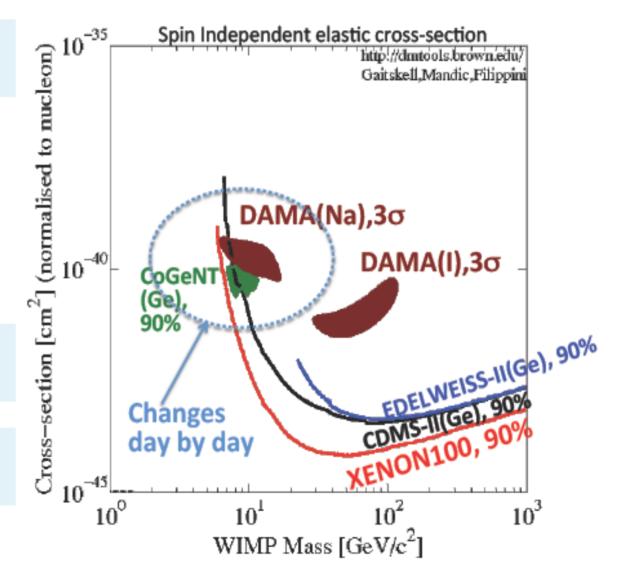
(see e.g., Ellis, Ferstl, Olive)

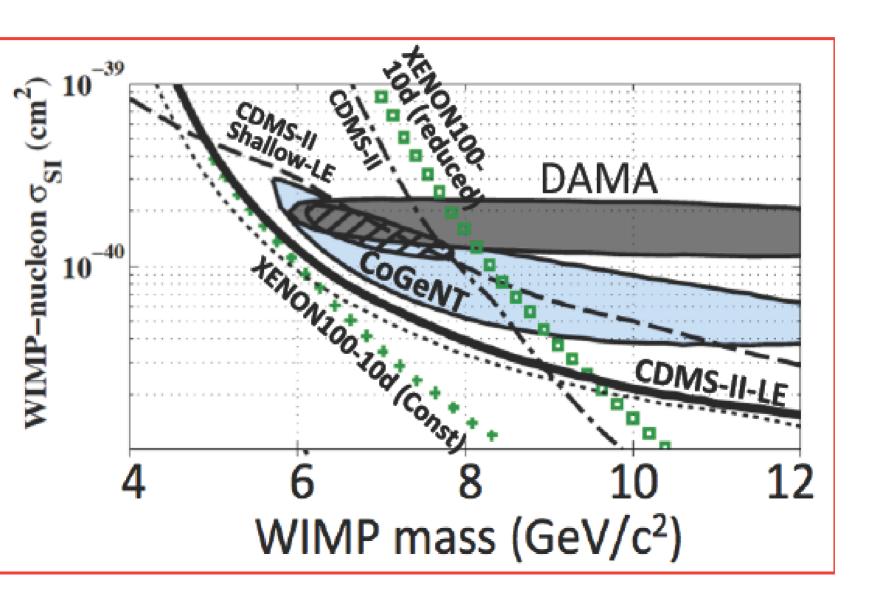
Recent Status

Sorry, We did not plot all the results

CRESST-II: Wait until their FINAL results

Low threshold Analysis by CDMS-II (LE) and XENON-10 (LE)





DM and NON-STANDARD COSMOLOGIES BEFORE NUCLEOSYNTHESIS

 NEUTRALINO RELIC DENSITY MAY DIFFER FROM ITS STANDARD VALUE, i.e. the value it gets when the expansion rate of the Universe is what is expected in Standard Cosmology (EX.: SCALAR-TENSOR THEORIES OF GRAVITY, KINATION, EXTRA-DIM. RANDALL-SUNDRUM TYPE II MODEL, ETC.)

• WIMPS MAY BE "COLDER", i.e. they may have smaller typical velocities and, hence, they may lead to smaller masses for the first structures which form GELMINI, GONDOLO

WHY H ≠ H_{GR}

$$H_{\rm GR}^2 = \frac{1}{3M_p^2} \rho_{\rm tot} \simeq 2.76 \, g_* \frac{T^4}{M_p^2}$$

Change the number of relativistic d.o.f.'s, g_{*};

R. Catena

- Consider a ρ_{tot} not dominated by relativistic d.o.f.'s;
 - Kination
 P. Salati, Phys. Lett. B 571 (2003) 121
- Consider theories where the effective Planck mass is different from the constant M_p:
 - Scalar-Tensor theories
 R. C., N. Fornengo, A. Masiero, M. Pietroni and F. Rosati, Phys. Rev. D 70 (2004) 063519
 - Extradimensions
 L. Randall and R. Sundrum, Phys. Rev. Lett. 83 (1999) 4690

DIRECT AND INDIRECT SEARCHES FOR WIMPs

PROBING NEW PHYSICS AT THE ELW.
 SCALE

 INFORMATION ON THE EVOLUTION OF THE EARLY UNIVERSE BEFORE THE NUCLEOSYNTHESIS TIME, i.e. at times < 1 sec.

4. ELW. SYMM. BREAKING STABILIZATION VS. FLAVOR PROTECTION: THE SCALE TENSION

$$M(B_d-B_d) \sim c_{SM} \frac{(y_t V_{tb}^* V_{td})^2}{16 \pi^2 M_W^2} + c_{new} \frac{1}{\Lambda^2}$$

If
$$c_{\text{new}} \sim c_{\text{SM}} \sim 1$$

Isidori

$$\Lambda > 10^4 \text{ TeV for O}^{(6)} \sim (\overline{\text{s}} \text{ d})^2$$

$$[K^0 - \overline{K^0} \text{ mixing }]$$

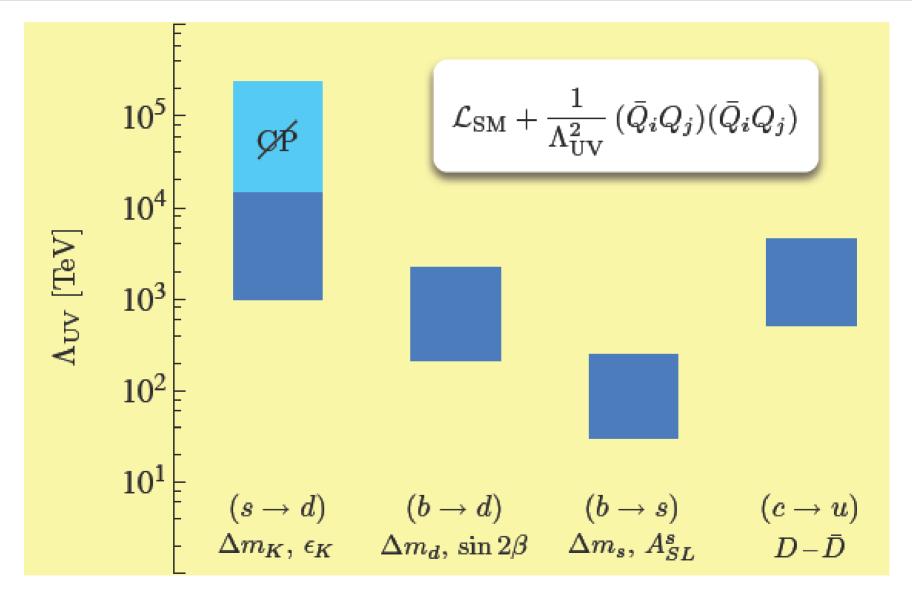
$$\Lambda > 10^3 \text{ TeV for O}^{(6)} \sim (\bar{b} \text{ d})^2$$

$$[B^0 - \bar{B}^0 \text{ mixing }]$$

UV SM COMPLETION TO STABILIZE THE ELW. SYMM. BREAKING: $\Lambda_{IIV} \sim O(1 \text{ TeV})$

Flavor Structure in the SM and Beyond

NEUBERT EPS-HEP 2011



Generic bounds without a flavor symmetry

$$K - \overline{K}$$
 8×10^{-7} 6×10^{-9} $D - \overline{D}$ 5×10^{-7} 1×10^{-7} $B - \overline{B}$ 5×10^{-6} 1×10^{-6} $B_s - \overline{B_s}$ 2×10^{-4} 2×10^{-4}

SMALLNESS OF THE NP COUPLINGS IF THE NP SCALE IS 1 TEV

SMALLNESS OF THE SM COUPLINGS

NIR

THE FLAVOUR PROBLEMS

FERMION MASSES

What is the rationale hiding behind the spectrum of fermion masses and mixing angles (our "Balmer lines" problem)

LACK OF A FLAVOUR "THEORY"

(new flavour – horizontal symmetry, radiatively induced lighter fermion masses, dynamical or geometrical determination of the Yukawa couplings, ...?)

FCNC

Flavour changing neutral current (FCNC) processes are suppressed.

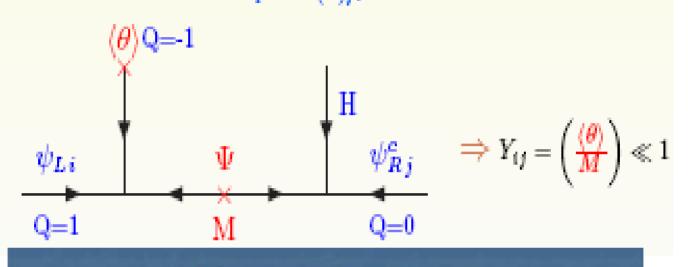
In the SM two nice mechanisms are at work: the GIM mechanism and the structure of the CKM mixing matrix.

How to cope with such delicate suppression if the there is new physics at the electroweak scale?

MSSM X FAMILY SYMM.

- AMBITION: simultaneously accounting for the "correct" SM fermion masses and mixings (SM Flavor Puzzle) and a structure of the SUSY soft breaking masses allowing for adequate FCNC suppression + possible "explanation" of the alleged SM FCNC difficulties (SUSY Flavor Puzzle)
- Mechanism a la Frogatt Nielsen with abelian or non-abelian family symmetry

 Froggatt-Nielsen mechanism and flavour symmetry to understand small Yukawa elements. Example: $U(1)_{fl}$



Yukawa Textures

What we want:

What we want:
$$Y_u \ll \begin{pmatrix} 0 & \varepsilon^3 & \varepsilon^3 \\ \varepsilon^3 & \varepsilon^2 & \varepsilon^2 \\ \varepsilon^3 & \varepsilon^2 & 1 \end{pmatrix} \quad Y_d \ll \begin{pmatrix} 0 & \varepsilon^3 & \varepsilon^3 \\ \varepsilon^3 & \varepsilon^2 & \varepsilon^2 \\ \varepsilon^3 & \varepsilon^2 & 1 \end{pmatrix}$$

$$\varepsilon = 0.05 \quad \varepsilon = 0.15$$

SU(3) Flavour model

ROBERTS, ROMANINO, ROSS, VELASCO-SEVILLA; ROSS, VELASCO-SEVILLA, VIVES

- $Q, L \sim 3$ and $d^c, u^c, e^c \sim 3$; flavon fields: $\theta_3, \theta_{23} \sim \overline{3}, \overline{\theta}_3, \overline{\theta}_{23} \sim 3$
- Family Symmetry breaking: $SU(3) \xrightarrow{(\theta_3)} SU(2) \xrightarrow{(\theta_{23})} \emptyset$

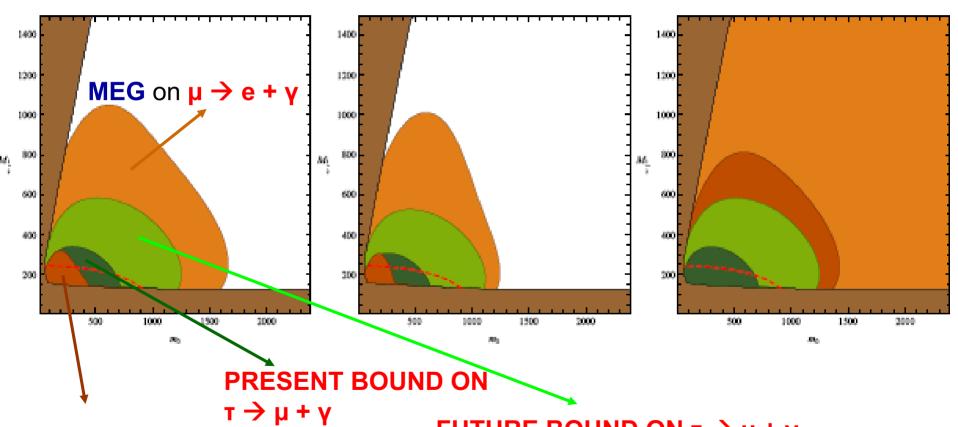
$$\theta_3, \overline{\theta}_3 = \left(\begin{array}{c} 0 \\ 0 \\ a_3 \end{array}\right), \ \ \theta_{23}, \overline{\theta}_{23} = \left(\begin{array}{c} 0 \\ b \\ b \end{array}\right) \text{with} \ \ \left(\frac{a_3}{M}\right) \sim \mathcal{O}(1), \ \left(\frac{b}{M_u}\right) \simeq \left(\frac{b}{M_d}\right)^2 = \varepsilon \sim 0.05.$$

• Yukawa superpotential: $W_Y = H\psi_i\psi_j^{\circ} \left[\theta_3^i\theta_3^j + \theta_{23}^i\theta_{23}^j \left(\theta_3\overline{\theta_3}\right) + \epsilon^{ikl}\overline{\theta}_{23,k}\overline{\theta_3}_{,l}\theta_{23}^j \left(\theta_{23}\overline{\theta_3}\right)\right]$

$$Y^{f} = \begin{pmatrix} 0 & a \, \varepsilon^{3} & b \, \varepsilon^{3} \\ a \, \varepsilon^{3} & \varepsilon^{2} & c \, \varepsilon^{2} \\ b \, \varepsilon^{3} & c \, \varepsilon^{2} & 1 \end{pmatrix} \frac{|a_{3}|^{2}}{M^{2}},$$

O. VIVES

LFV CONSTRAINTS IN THE $M_0 - M_{1/2}$ SUSY PLANE



PRESENT BOUND ON

 $\mu \rightarrow e + v$

FUTURE BOUND ON $\tau \rightarrow \mu + \gamma$ at SUPER B

CALIBBI, JONES, A.M., J-H. PARK, POROD and VIVES

FLAVOR BLINDNESS OF THE NP AT THE ELW. SCALE?

- THREE DECADES OF FLAVOR TESTS (Redundant determination of the UT triangle → verification of the SM, theoretically and experimentally "high precision" FCNC tests, ex. b → s + γ, CP violating flavor conserving and flavor changing tests, lepton flavor violating (LFV) processes, ...) clearly state that:
- A) in the HADRONIC SECTOR the CKM flavor pattern of the SM represents the main bulk of the flavor structure and of (flavor violating) CP violation;
- B) in the LEPTONIC SECTOR: although neutrino flavors exhibit large admixtures, LFV, i.e. non – conservation of individual lepton flavor numbers in FCNC transitions among charged leptons, is extremely small: once again the SM is right (to first approximation) predicting negligibly small LFV

What to make of this triumph of the CKM pattern in hadronic flavor tests?

New Physics at the Elw.
Scale is Flavor Blind
CKM exhausts the flavor
changing pattern at the elw.
Scale

MINIMAL FLAVOR VIOLATION

MFV: Flavor originates only from the SM Yukawa coupl.

New Physics introduces

NEW FLAVOR SOURCES in addition to the CKM pattern. They give rise to contributions which are <10% in the "flavor observables" which have already been observed!

SuperB vs. LHC Sensitivity Reach in testing Λ_{SUSY}

	superB	general MSSM	high-scale MFV
$\left \left(\delta^d_{13} \right)_{LL} \right (LL \gg RR)$	$1.8 \cdot 10^{-2} \frac{m_{\tilde{q}}}{(350 \text{GeV})}$	1	$\sim 10^{-3} rac{(350 { m GeV})^2}{m_{ ilde{q}}^2}$
$ \left(\delta^d_{13}\right)_{LL} (LL \sim RR)$	$1.3 \cdot 10^{-3} \frac{m_{\tilde{q}}}{(350 \text{GeV})}$	1	_
$ \left(\delta^d_{13}\right)_{LR} $	$3.3 \cdot 10^{-3} \frac{m_{\tilde{q}}}{(350 \text{GeV})}$	$\sim 10^{-1} \tan \beta \frac{(350 \text{GeV})}{m_{\tilde{q}}}$	$\sim 10^{-4} \tan \beta \frac{(350 { m GeV})^3}{m_{\tilde{q}}^3}$
$ \left(\delta^d_{23}\right)_{LR} $	$1.0 \cdot 10^{-3} \frac{m_{\tilde{q}}}{(350 \mathrm{GeV})}$	$\sim 10^{-1} \tan \beta \frac{(350 { m GeV})}{m_q}$	$\sim 10^{-3} \tan \beta \frac{(350 { m GeV})^3}{m_q^3}$

SuperB can probe MFV (with small-moderate $tan\beta$) for TeV squarks; for a generic non-MFV MSSM sensitivity to squark masses > 100 TeV !

Ciuchini, Isidori, Silvestrini SLOW-DECOUPLING OF NP IN FCNC

Estimates of error for 2015

	元文		
4.0	1	ş I	

Hadronic matrix	Current lattice	6 TFlop Year	60 TFlop Year	1-10 PFlop Year
element	error	7641	[2011 LHCb]	[2015 SuperB
$f_{+}^{K\pi}(0)$	0.9%	0.7%	0.4%	< 0.1%
1 ₊ (0)	(22% on 1-f ₊)	(17% on 1-f ₊)	(10% on 1-f ₊)	(2.4% on 1-f ₊)
$\mathbf{\hat{B}}_{\mathrm{K}}$	11%	5%	3%	1%
f_B	14%	3.5 - 4.5%	2.5 - 4.0%	1-1.5%
$\mathbf{f}_{\mathrm{B}s}\mathbf{B}_{\mathrm{B}s}^{1/2}$	13%	4 - 5%	3 - 4%	1-1.5%
ξ	5%	3%	1.5 - 2 %	0.5 – 0.8 %
2	(26% on ξ-1)	(18% on ξ-1)	(9-12% on ξ-1)	(3-4% on ξ-1)
$\mathcal{F}_{B \to D/D^*l\nu}$	4%	2%	1.2%	0.5%
	$(40\% \text{ on } 1-\mathcal{F})$	(21% on 1- \mathcal{F})	(13% on 1- \mathcal{F})	(5% on 1- <i>F</i>)
$\mathbf{f}_{\scriptscriptstyle{+}}^{\mathrm{B}\pi},\!\dots$	11%	5.5 - 6.5%	4 - 5%	2-3%
$T_1^{B \to K^*/\rho}$	13%			3-4%

SUSY SEE-SAW

• UV COMPLETION
OF THE SM TO
STABILIZE THE
ELW. SCALE:

LOW-ENERGY
SUSY

 COMPLETION OF THE SM FERMIONIC SPECTRUM TO ALLOW FOR **NEUTRINO MASSES:** NATURALLY SMALL PHYSICAL NEUTRINO MASSES WITH RIGHT-HANDED NEUTRINO WITH A LARGE MAJORANA MASS

SEE-SAW

LFV and NEW PHYSICS

- Flavor in the HADRONIC SECTOR:
 CKM paradigm
- Flavor in the LEPTONIC SECTOR:
 - Neutrino masses and (large) mixings
 - Extreme smallness of LFV in the charged lepton sector of the SM with massive neutrinos

$$I_{i}$$
 suppressed by $(m_{v_{i}}^{2} - m_{v_{k}}^{2})/M_{W}^{2}$

NEW BOUND OF MEG AT THE EPS 2011

The MEG Experiment

$$\mu^+ \to e^+ \gamma$$

SHOWN AT ICHEP 2010

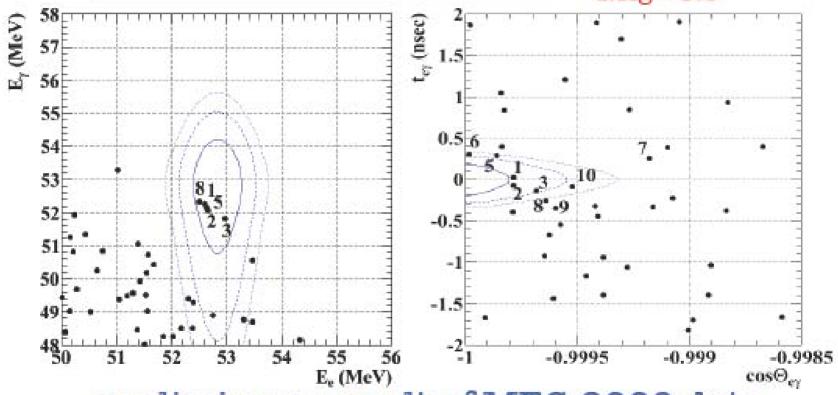


Event distribution after unblinding

BR < 1.5×10-11 @90%CL

6.1×10-12 expected

Nsig = 3.0

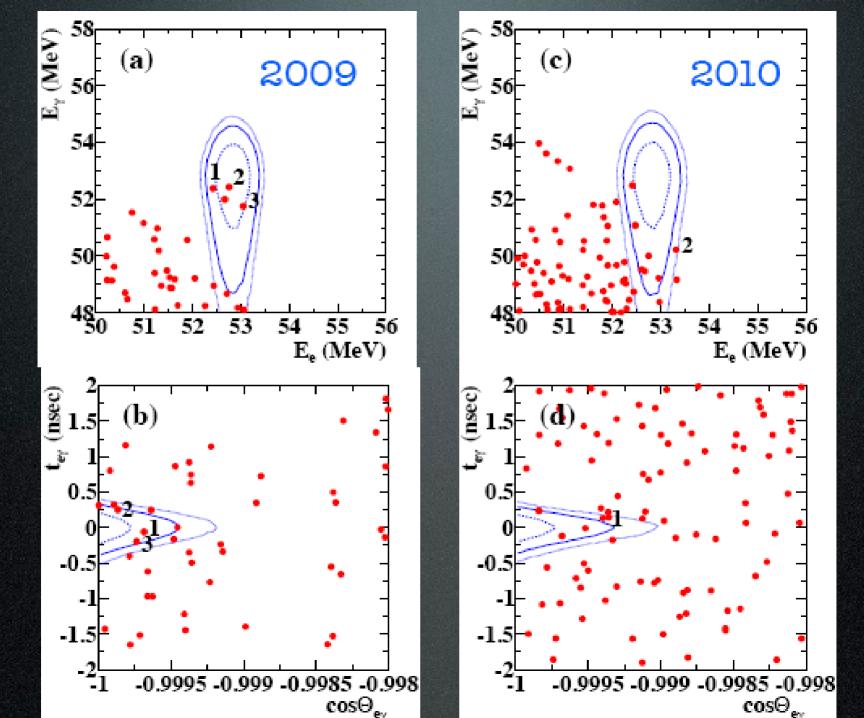


preliminary result of MEG 2009 data

Blue lines are 1(39.3 % included inside the region w.r.t. analysis window), 1.64(74.2%) and 2(86.5%) sigma regions.

For each plot, cut on other variables for roughly 90% window is applied.

Numbers in figures are ranking by Last/(Lnmp+Lag). Same numbered dots in the right and the left figure are an identical event.



MEG summary

- 2009+2010 data consistent w/ no signal
- New physics is now constrained by 5× tighter upper limit:
 BR < 2.4×10⁻¹² @90% C.L.
 (Preprint will be posted at arXiv today)
- MEG is accumulating more data this and next year to reach O(10⁻¹³) sensitivity;
 So stay tuned!
- Detector improvements/upgrades

T. MORI AT THE EPS-HEP 2011

SUSY SEESAW: Flavor universal SUSY breaking and yet large lepton flavor violation

Borzumati, A. M. 1986 (after discussions with W. Marciano and A. Sanda)

$$L = f_l \overline{e}_R L h_1 + f_v \overline{v}_R L h_2 + M v_R v_R$$

$$\tilde{L}_{\nu_R} \longrightarrow (m_{\tilde{L}}^2)_{ij} \Box \frac{1}{8\pi^2} (3m_0^2 + A_0^2) (f_v^{\dagger} f_v)_{ij} \log \frac{M}{M_G}$$

Non-diagonality of the slepton mass matrix in the basis of diagonal lepton mass matrix depends on the unitary matrix U which diagonalizes (f_v ⁺ f_v)

How Large LFV in SUSY SEESAW?

- 1) Size of the Dirac neutrino couplings f_v
- 2) Size of the diagonalizing matrix U

In MSSM seesaw or in SUSY SU(5) (Moroi): not possible to correlate the neutrino Yukawa couplings to know Yukawas;

In **SUSY SO(10)** (A.M., Vempati, Vives) at least one neutrino Dirac Yukawa coupling has to be of the **order of the top Yukawa coupling** one large of O(1) f_v

U ____ two "extreme" cases:

- a) U with "small" entries U = CKM;
- b) **U with "large" entries** with the exception of the 13 entry **U = PMNS** matrix responsible for the diagonalization of the neutrino mass matrix

THE STRONG ENHANCEMENT OF LFV IN SUSY SEESAW **MODELS CAN OCCUR EVEN IF THE MECHANISM** RESPONSIBLE FOR SUSY **BREAKING IS ABSOLUTELY** FLAVOR BLIND

LFV in SUSYGUTs with SEESAW



Scale of pearance of the SUSY soft breaking terms resulting from the spontaneous breaking of supergravity

Low-energy SUSY has "memory" of all the multi-step RG occurring from such superlarge scale down to M_W

potentially large <u>LFV</u>

Barbieri, Hall; Barbieri, Hall, Strumia; Hisano, Nomura,

Yanagida; Hisano, Moroi, Tobe Yamaguchi; Moroi; A.M.,, Vempati, Vives;

Carvalho, Ellis, Gomez, Lola; Calibbi, Faccia, A.M, Vempati

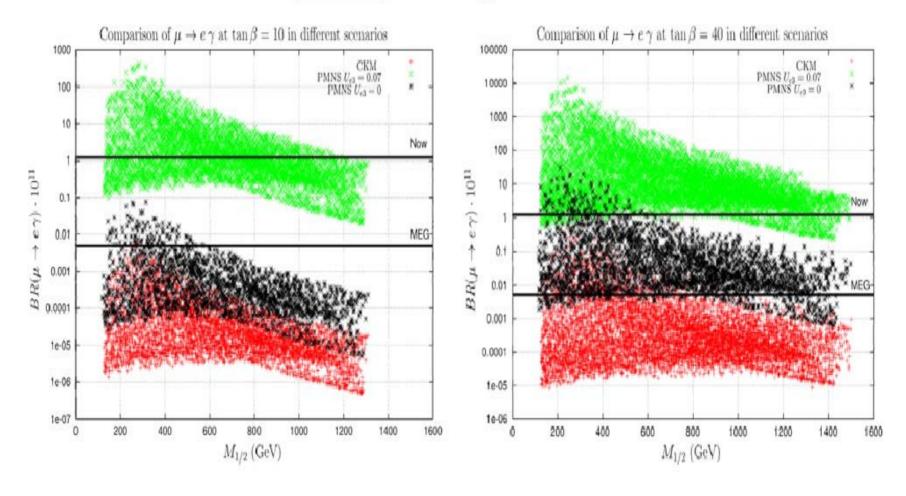
LFV in MSSMseesaw: μ e γ Borzumati, A.M.

τ μγ Blazek, King;

General analysis: Casas Ibarra; Lavignac, Masina, Savoy; Hisano, Moroi, Tobe, Yamaguchi; Ellis, Hisano, Raidal, Shimizu; Fukuyama, Kikuchi, Okada; Petcov, Rodejohann, Shindou, Takanishi; Arganda, Herrero; Deppish, Pas, Redelbach, Rueckl; Petcov, Shindou

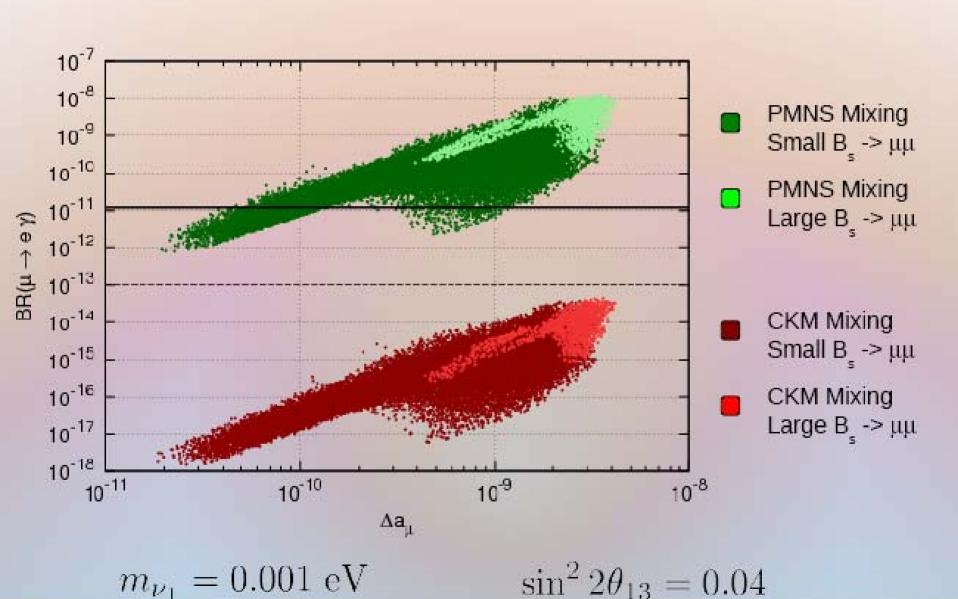
$\mu \rightarrow e+\gamma$ in SUSYGUT: past and future

 $\mu
ightarrow e \, \gamma$ in the U_{e3} = 0 PMNS case

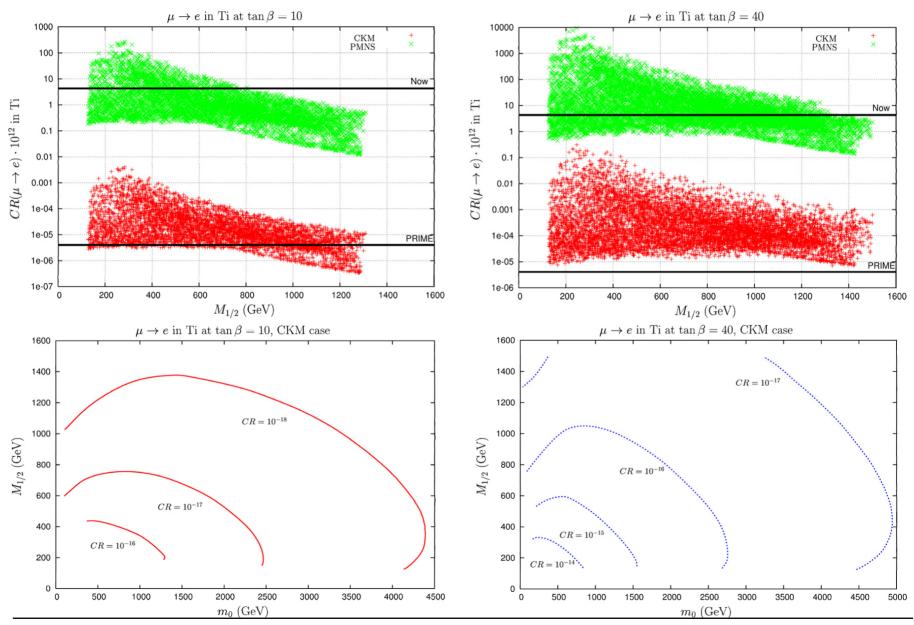


CFMV

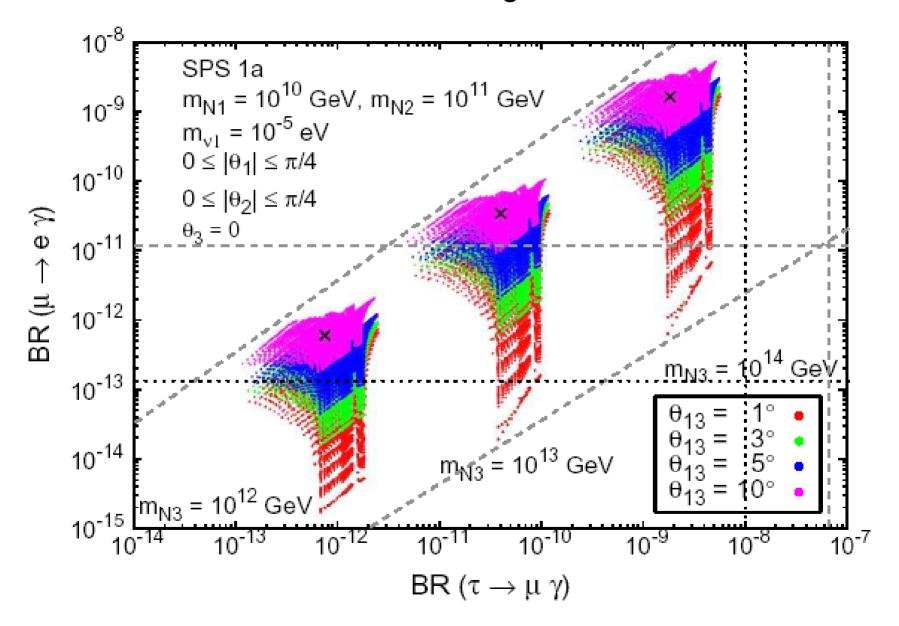
Comparing CKM and PMNS



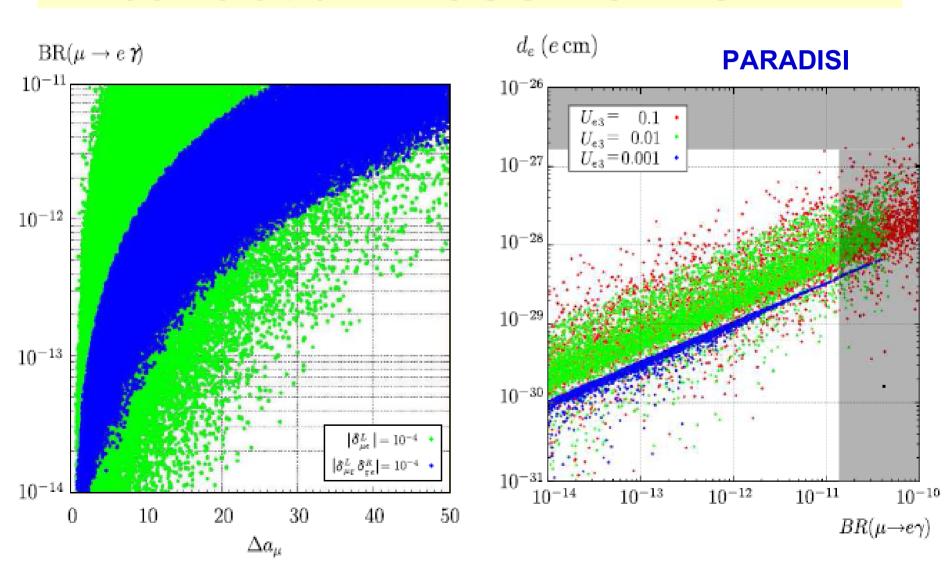
$\mu ightarrow e ext{ in Ti}$ and **PRISM/PRIME** conversion experiment



Antusch, Arganda, Herrero, Teixeira



LFV, g - 2, EDM: a promising correlation in SUSY SEESAW



DEVIATION from μ - e UNIVERSALITY

A.M., Paradisi, Petronzio

• Denoting by $\Delta r_{NP}^{e-\mu}$ the deviation from $\mu - e$ universality in $R_{K,\pi}$ due to new physics, i.e.:

$$R_{K,\pi} = R_{K,\pi}^{SM} \left(1 + \Delta r_{K,\pi NP}^{e-\mu} \right),$$

• we get at the 2σ level:

$$-0.063 \le \Delta r_{KNP}^{e-\mu} \le 0.017 \text{ NA48/2}$$

$$-0.0107 \le \Delta r_{\pi NP}^{e-\mu} \le 0.0022$$
 PDG

Presently: error on R_K down to the **1% level** (KLOE (09) and NA48 (07 data);using 40% of the data collected in 08, NA62 is now decreasing the uncertainty at the **0.7% level Prospects**: Summer conf. we'll have the result concerning the 40% data analysis by NA62 and when the analysis of the whole sample of data is accomplished **the stat. uncertainty will be < 0.3%**

HIGGS-MEDIATED LFV COUPLINGS

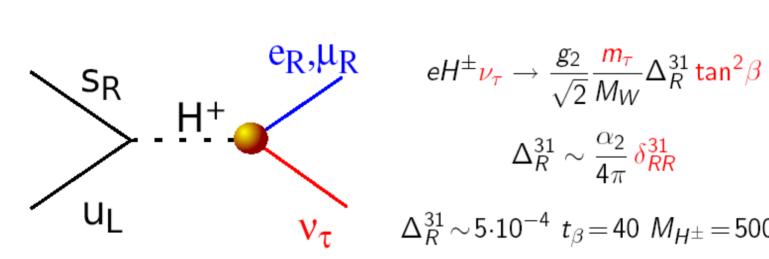
- When non-holomorphic terms are generated by loop effects (HRS corrections)
- And a source of LFV among the sleptons is present
- Higgs-mediated (radiatively induced) H-lepton-lepton LFV couplings arise Babu, Kolda; Sher; Kitano, Koike, Komine,

Okada; Dedes, Ellis, Raidal; Brignole, Rossi; Arganda, Curiel, Herrero, Temes; Paradisi;

Brignole, Rossi

H mediated LFV SUSY contributions to R_k

$$R_K^{LFV} = \frac{\sum_i K \to e\nu_i}{\sum_i K \to \mu\nu_i} \simeq \frac{\Gamma_{SM}(K \to e\nu_e) + \Gamma(K \to e\nu_\tau)}{\Gamma_{SM}(K \to \mu\nu_\mu)} \ , \ \ i = e, \mu, \tau$$



$$eH^\pm
u_ au
ightarrow rac{g_2}{\sqrt{2}}rac{m_ au}{M_W}\Delta_R^{31} an^2eta \ \Delta_R^{31}\sim rac{lpha_2}{A_{TT}}\delta_{RR}^{31}$$

$$V_{\tau}$$
 $\Delta_R^{31} \sim 5.10^{-4} t_{\beta} = 40 M_{H^{\pm}} = 500 \text{GeV}$

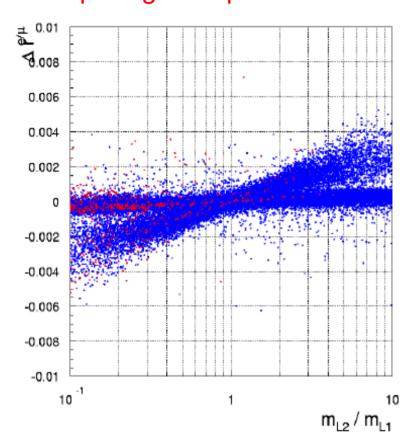
$$\Delta au_{\!K\,SUSY}^{\mathrm{e}-\mu} \simeq \left(rac{m_K^4}{M_{H^\pm}^4}
ight)\!\left(rac{m_ au^2}{m_{\mathrm{e}}^2}
ight) |\Delta_R^{31}|^2 \! an^6\! eta pprox 10^{-2}$$

Extension to B \rightarrow I_V deviation from universality Isidori, Paradisi

LFU breaking occurs in a **LF** conserving case because of the splitting in slepton masses

LFU breaking occurs with LFV

A.M., PARADISI. PETRONZIO



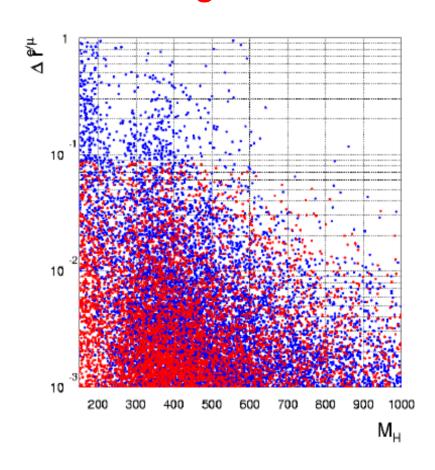


Figure 2: Left: $\Delta r_K^{e/\mu}$ as a function of the mass splitting between the second and the first (left-handed) slepton generations. Red dots can saturate the $(g-2)_{\mu}$ discrepancy at the 95% C.L., i.e. $1 \times 10^{-9} < (g-2)_{\mu} < 5 \times 10^{-9}$. Right: $\Delta r_K^{e/\mu}$ as a function of M_{H^+} .

SUSY GUTS

UV COMPLETION
 OF THE SM TO
 STABILIZE THE
 ELW. SCALE:

LOW-ENERGY SUSY TREND OF
UNIFICATION OF
THE SM GAUGE
COUPLINGS AT
HIGH SCALE:

GUTs

Large v mixing ← large b-s transitions in SUSY GUTs

```
In SU(5) d_R \longrightarrow I_L connection in the 5-plet Large (\Delta^l_{23})_{LL} induced by large f_v of O(f_{top}) is accompanied by large (\Delta^d_{23})_{RR}
```

In SU(5) assume large f_v (Moroi) In SO(10) f_v large because of an underlying Pati-Salam symmetry

(Darwin Chang, A.M., Murayama)

See also: Akama, Kiyo, Komine, Moroi; Hisano, Moroi, Tobe, Yamaguchi, Yanagida; Hisano, Nomura; Kitano, Koike, Komine, Okada

FCNC HADRON-LEPTON CONNECTION IN SUSYGUT

soft SUSY breaking terms arise at a scale > M_{GUT} , they have to respect the underlying quark-lepton GU symmetry constraints on δ^{quark} from LFV and constraints on δ^{lepton} from hadronic FCNC

Ciuchini, A.M., Silvestrini, Vempati, Vives PRL 2004

general analysis Ciuchini, A.M., Paradisi, Silvestrini, Vempati, Vives NPB 2007

For previous works: Baek, Goto, Okada, Okumura PRD 2001;

Hisano, Shimizu, PLB 2003;

Cheung, Kang, Kim, Lee PLB 2007

Borzumati, Mishima, Yamashita hep-ph 0705:2664

For recent works: Goto, Okada, Shindou, Tanaka PRD 2008;

Ko, J-h. Park, Yamaguchi arXiv:0809:2784

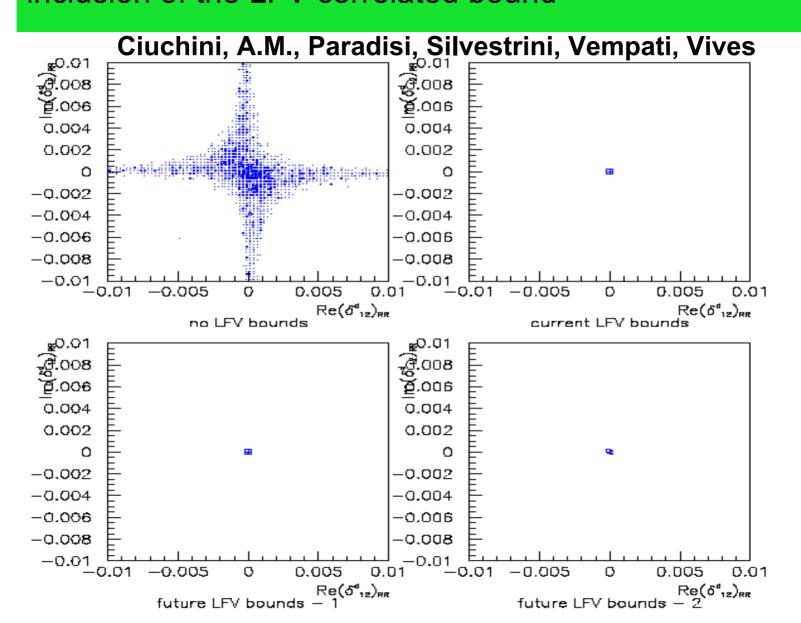
GUT-RELATED SUSY SOFT BREAKING TERMS

$$\begin{split} m_Q^2 &= m_{\tilde{e^c}}^2 = m_{\tilde{u^c}}^2 = m_{10}^2 \\ m_{\tilde{d^c}}^2 &= m_L^2 = m_{\overline{5}}^2 \\ A_{ij}^e &= A_{ji}^d \,. \end{split}$$

SU(5) RELATIONS

	Relations at weak-scale	Relationss at $M_{\rm GUT}$
(1)	$(\delta^u_{ij})_{\mathrm{RR}}~\approx~(m_{e^c}^2/m_{u^c}^2)~(\delta^l_{ij})_{\mathrm{RR}}$	$m_{u^c_0}^2 = m_{e^c_0}^2$
(2)	$(\delta_{ij}^q)_{ m LL} \approx (m_{e^c}^2/m_Q^2) (\delta_{ij}^l)_{ m RR}$	$m_{Q_0}^2 = m_{e^c_0}^2$
(3)	$(\delta^d_{ij})_{\mathrm{RR}}~pprox~(m_L^2/m_{d^c}^2)~(\delta^l_{ij})_{\mathrm{LL}}$	$m_{d^c_0}^2 = m_{L_0}^2$
(4)	$(\delta_{ij}^d)_{ m LR} \approx (m_{L_{avg}}^2/m_{Q_{avg}}^2) (m_b/m_{ au}) (\delta_{ij}^l)_{ m LR}^{\star}$	$A^e_{ij_0} = A^d_{ji_0}$

Bounds on the hadronic $(\delta_{12})_{RR}$ as modified by the inclusion of the LFV correlated bound



3 QUESTIONS

- Are we sure that there is new physics (NP) at the TeV scale? YES (barring an antropic approach)
- If yes, are we sure that LHC will see something "new", i.e. beyond the SM with its "standard higgs boson"? YES
- If there is new physics at the TeV scale, what can flavor and DM physics tell to LHC and viceversa? (or, putting it in a less politically correct fashion: if LHC starts seeing some new physics signals, are flavor and DM physics still a valuable road to NP, or are they definitely missing that train? NO, actually to catch the "right train" it is highly desirable, though maybe strictly not necessary, to make use of all the three roads at the same time

A FUTURE FOR FLAVOR PHYSICS IN OUR SEARCH BEYOND THE SM?

- The traditional competition between direct and indirect (FCNC, CPV) searches to establish who is going to see the new physics first is no longer the priority, rather
- COMPLEMENTARITY between direct and indirect searches for New Physics is the key-word
- Twofold meaning of such complementarity:
- i) synergy in "reconstructing" the "fundamental theory" staying behind the signatures of NP;
- ii) coverage of complementary areas of the NP parameter space (ex.: multi-TeV SUSY physics)