



Gunion Fest, March 2014

# Naturalness and the Higgs Boson

James Wells

March 29, 2014

# Great to be back at UC Davis

Was faculty here from 1999-2002. I was very fortunate to have such a tremendous role model at that stage.



(Somehow Jack and others had patience with me.)

Brash young man

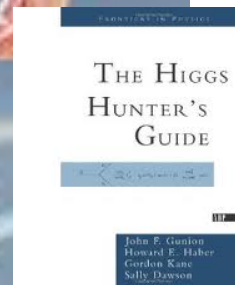
I had certain preconceived notions  
of Jack before arriving....



THE HIGGS  
HUNTER'S  
GUIDE

John F. Ganton  
Howard E. Haber  
Gordon Kane  
Sally Dawson

But over time I saw a different Jack



Over time I also saw his incredible dedication to his physics research

Nothing got in the way.

During a lunch conversation somebody was talking about Kobe Bryant's amazing performance the previous day.

And Jack said ...



Who's Kobe Bryant?





Kobe Bryant: Forward for Los Angeles Lakers.  
NBA MVP and frequent All-Star.  
One of greatest scorers of all time.

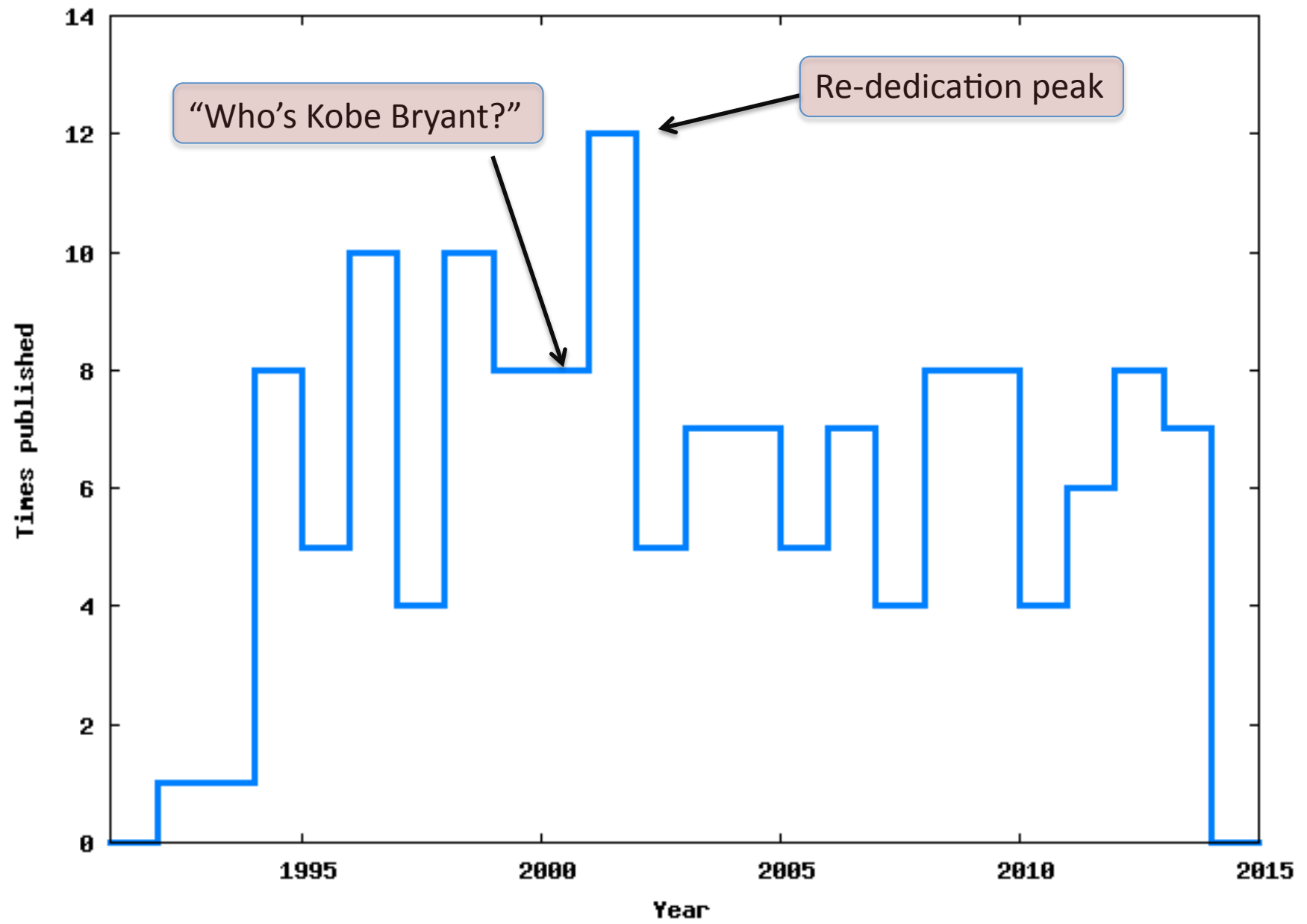
I looked up to Jack and decided to re-dedicated myself.

No more watching basketball  
and Kobe Bryant.

No more excessive attention to politics.

No more obsessing over  
“Buffy the Vampire Slayer”

(never mind)



Jack's precision calculations prowess went well beyond particle physics.

I was on a faculty committee charged to give an award to a student.

There were three nominations.

Each one looked great, and indistinguishable to me.



# How can we decide who is best?



My strategy, as usual, was ...

... WWJD = What would Jack Do?

So, I asked, “Jack, who’s best?”



Without hesitation Jack said,



1<sup>st</sup>

2<sup>nd</sup>

3<sup>rd</sup>

That ended the discussion for me...

Other committee members needed  
more convincing

Somebody had the brilliant idea to...

... look at the students' GPAs!!

We shuffled through the folders and  
found the data



3.82

3.81

3.77

1 part in  $10^3$  precision calculation  
ability, not just in particle physics!

We all admire Jack's incredible

Energy

Dedication

Consistency

Technical Skill

Creative Insight

Go-to expertise in Higgs Physics etc

Jack's contributions to Higgs physics in particular has been second to none in the world, and continues to have great influence.



# Naturalness and the Higgs Boson

*Audience caution: you may experience slight dizziness, minor irritations at times, and in rare circumstances loss of consciousness.*



# Should we believe in the Higgs boson?

The Higgs boson is a speculative particle explanation for elementary particle masses.

Cons:

1. One particle carries all burdens of mass generation?
2. Fundamental scalar not known in nature.
3. Hasn't been found yet.
4. Too simplistic -- dynamics for vev not built in.
5. Idea not stable to quantum corrections.

Fixing that → New Physics

Pros: **Still consistent with experimental facts!**

# New Physics Ideas and Higgs boson viability

Trying to fix and understand Higgs physics leads to new ideas that have new states/dynamics (susy, extra dimensions,  $t'$ , etc.) and a scalar state that look very similar to the Standard Model Higgs boson.

What have we found so far?

# Physicists Find Elusive Particle Seen as Key to Universe



Pool photo by Denis Ballbouse

Scientists in Geneva on Wednesday applauded the discovery of a subatomic particle that looks like the Higgs boson.

By DENNIS OVERBYE

Published: July 4, 2012 |  122 Comments

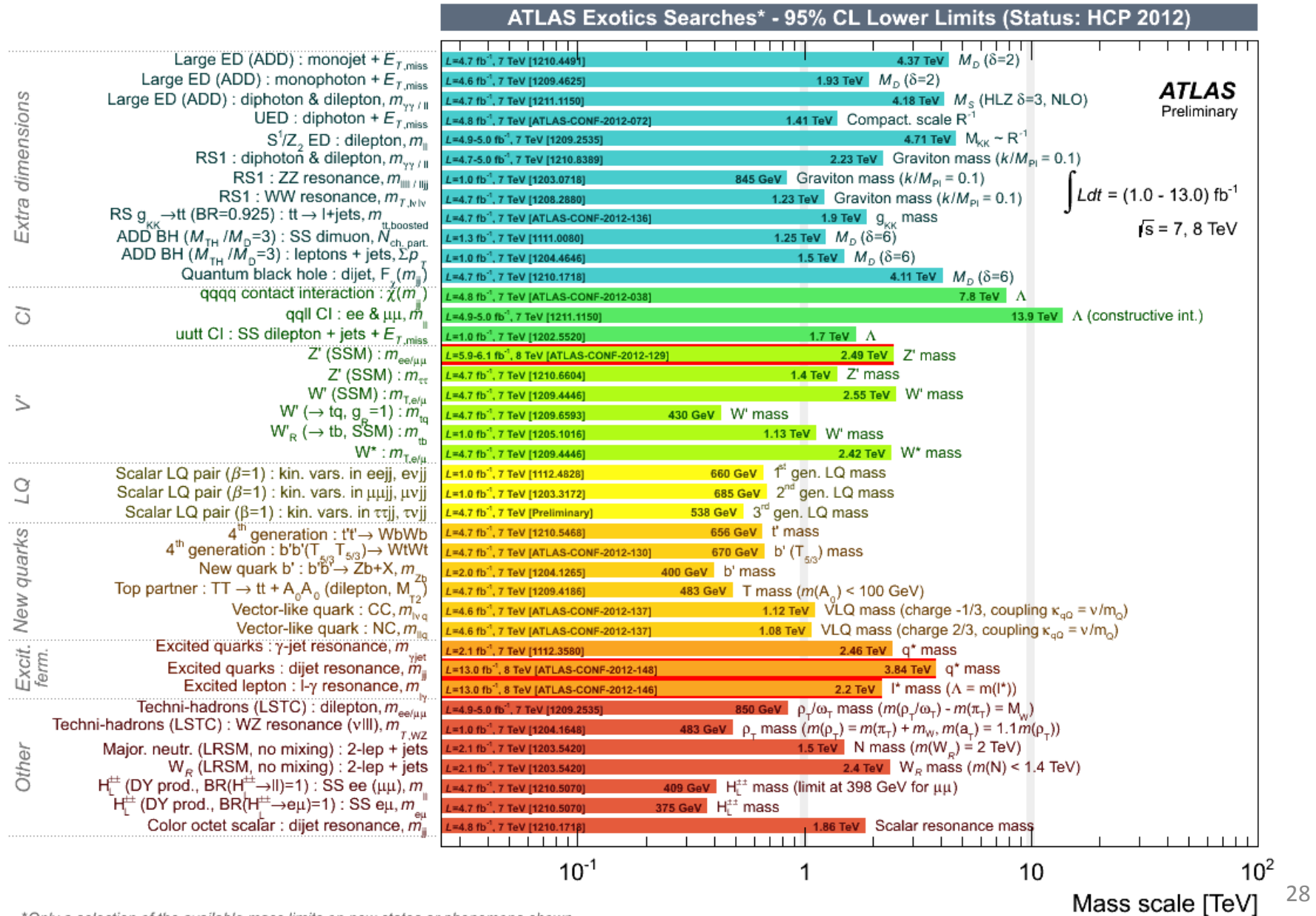
mass = 126 GeV

New York Times

# But nothing else has been found so far....

4/24/13

AtlasSearches\_exotics\_hcp12.png (1096x827)



\*Only a selection of the available mass limits on new states or phenomena shown




# Losing the Naturalness Religion

Starting to hear many more comments like:

“Quadratic divergence Naturalness problem is just philosophical – not really a data-driven concern.”

“Dimensional regularization has no quadratic divergence Naturalness problem, so maybe it doesn’t exist”


$$\rightarrow m_W^2 \left( \frac{1}{4-n} - \gamma_E + \ln 4\pi + 1 - \ln \frac{m_W^2}{\mu^2} \right) + \dots$$

(Note, there is no  $\Lambda^2$  cutoff funny business – only  $1/(4-n)$ )

# Sensitivity to higher physical scales persists

However, all it takes is for any massive particle to interact with the Higgs and there is a real physical quantum correction to contend with.



$$\Delta m_H^2 \propto \lambda_\Phi m_\Phi^2 \ln m_\Phi$$

It is inconceivable to me that there is nothing else between “here” ( $10^2$  GeV) and the Planck scale ( $10^{18}$  GeV). And if there is another scalar (even if exotically charged!) there is no simple symmetry to forbid it from coupling to the Higgs boson.

*Implicit Postulate of Absolute Naturalness:*

A large hierarchy in QFT (even "technically natural") requires further explanation by further dynamics or an additional principle. Either way: "new physics".

Higgs without yet discovering supporting entourage (susy, Xdim, composite states, etc.) is disheartening to some.

Anguished query from the youngsters:

Should I still take this "philosophical concept" of Naturalness seriously as a guiding principle?

## Showing that Naturalness Principle is Effective

Naturalness has been the oxygen of “Beyond the Standard Model Physics”.

We can show the concept’s effectiveness (if it is) by waiting for new dynamics to arise at LHC that stabilizes the Higgs boson to these Naturalness-voiding quadratic divergences.

Or we can try to put it on trial now.



Naturalness defense attorney<sup>32</sup>



# QED Application

Let's test the principle of Naturalness as a guide by applying it to the past.

Example: The early days of Quantum Electrodynamics.

*Specifically, why is the electron so light??*

(I wonder: Why didn't "they" ask that question more earnestly?)

# Quantum Electrodynamics

$$\mathcal{L} = \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\gamma^\mu(\partial_\mu - ieA_\mu)\psi + m_e\bar{\psi}\psi$$

A,F contain the photon and  $\psi$  is the electron ( $m_e=5\times 10^{-4}$  GeV).

Extraordinary theory:

1. Relativistic invariant
2. Dirac equation built in
3. Massless photon – electromagnetic radiation
4. Electromagnetic gauge invariance
5. Renormalizable! (infinities easily handled)
6. Fits the low-energy data very well

Subsequent theories had to live up to QED (e.g., renormalizability of the massive weak interactions, etc.)

# The electron mass

Hard to imagine QED being criticized from the perspective of the 1940s. But what if they took naturalness seriously?

**Naturalist:** Why is the electron mass so small?

**Skeptic:** Small compared to what?

**Naturalist:** Newton's gravity scale  $(G_N)^{-1/2} = 10^{18}$  GeV

**Skeptic:** What's gravity got to do with a little particle's mass?

**Naturalist:**  $G_N$  is a dimensionful scale parameter in the action of natural law just like the electron mass is. How can we have such a large hierarchy between dimensionful numbers?

**Skeptic:** Nobody understands gravity. It's not renormalizable. It's too remote. We don't understand it.

**Naturalist:** Precisely! We don't understand it, so let's try.

**Skeptic:** Maybe Dirac was right with his Large Number Hypothesis. The universe spots us one very large number and we just have to live with it. I'm willing to live with that, especially if it involves mysterious and remote gravity!

**Naturalist:** Ok. I'll give you that for now. Who am I to argue with Dirac. But what about the electron to proton mass ratio  $m_p/m_e = 10^3$ .

**Skeptic:**  $10^3$  is no big deal. That can just be an accident. People do get struck by lightning you know, and that's much more rare.

**Naturalist:** Ok, then what about Fermi's new theory of  $\beta$  decay. His constant is orders of magnitude larger yet than the proton!

**Skeptic:** Hmm. (Stalling) Remind me about Fermi's Theory....

E. Fermi, Zeitschrift für Physik, 88 (1934) 161.

## Versuch einer Theorie der $\beta$ -Strahlen. I<sup>1)</sup>.

Von E. Fermi in Rom.

Mit 3 Abbildungen. (Eingegangen am 16. Januar 1934.)

Eine quantitative Theorie des  $\beta$ -Zerfalls wird vorgeschlagen, in welcher man die Existenz des Neutrinos annimmt, und die Emission der Elektronen und Neutrinos aus einem Kern beim  $\beta$ -Zerfall mit einer ähnlichen Methode behandelt, wie die Emission eines Lichtquants aus einem angeregten Atom in der Strahlungstheorie. Formeln für die Lebensdauer und für die Form des emittierten kontinuierlichen  $\beta$ -Strahlenspektrums werden abgeleitet und mit der Erfahrung verglichen.

Sehen wir zunächst von den Relativitätskorrekturen und der Spinwirkung ab, so ist wohl die einfachst mögliche Wahl von (9) die folgende:

$$H = g \{ Q \psi(x) \varphi(x) + Q^* \psi^*(x) \varphi^*(x) \}, \quad (10)$$

wo  $g$  eine Konstante mit den Dimensionen  $L^5 M T^{-2}$  darstellt;  $x$  reprä-

Aus den Daten der Tabelle 2 kann man eine, wenn auch sehr grobe, Abschätzung der Konstante  $g$  gewinnen. Nimmt man etwa an, daß in den Fällen wo (50) gleich Eins wird, man  $\tau F(\eta_0) = 1$  hat (d. h., in Sekunden, = 3600), so bekommt man aus (45):

$$g = 4 \cdot 10^{-50} \text{ cm}^3 \cdot \text{erg}.$$

Dieser Wert gibt natürlich nur die Größenordnung von  $g$ .

Theory  
innovation

Experimental  
savvy

# Fermi's Constant in the day...

$$g = 4 \times 10^{-50} \text{ cm}^3 \text{ erg} = 3.25 \times 10^{-6} \text{ GeV}^{-2}$$

This translates into scale  $M_F = g^{-1/2} = 555 \text{ GeV}$ .

$M_F/m_e = 10^6$  Now that's starting to be concerning!

Today, we quote Fermi's constant as

$G_F = 1.166 \times 10^{-5} \text{ GeV}^{-2}$  where the normalization is set by

$$\frac{G_F}{\sqrt{2}} \left[ \bar{u} \gamma^\mu \frac{(1 - \gamma_5)}{2} d \right] \left[ \bar{e} \gamma_\mu \frac{(1 - \gamma_5)}{2} \nu \right]$$

# *Agreed: small $m_e$ not Natural*

Now what do we do?

We cogitate... We wait... We ponder...

And then a Realistic Intellectual Leap (RIL) happens.

*RIL #1:  $m_e$  is closer zero than it is to  $M_F$ .*

Skeptic: What does that do for us?

Answer: Let's try to start with a theory that forbids electron mass and see if we have an idea to let in a little bit of mass later.



# Forbidding the electron mass

It's obvious that the problem is that  $\psi\psi$  is gauge invariant. How can we make it non-gauge invariant.

We cogitate... We wait... We ponder...

*RIL #2: The representation structure of the Lorentz group allows us to write QED in two component spinors:*

$$\mathcal{L} = \frac{1}{4}F_{\mu\nu}F^{\mu\nu} + i\psi_L^\dagger\gamma^\mu(\partial_\mu - ieA_\mu)\psi_L + \psi_R^\dagger\gamma^\mu(\partial_\mu - ieA_\mu)\psi_R + m_e(\psi_L^\dagger\psi_R + \psi_R^\dagger\psi_L)$$

Laporte & Uhlenbeck (1931):

$$mc\chi_\ell - \left( \frac{h}{i} \partial_{\dot{\ell}} + \phi_{\dot{\ell}} \right) \psi_{\dot{\ell}} = 0$$
$$mc\psi_{\dot{m}} + \left( \frac{h}{i} \partial_{\dot{m}}^\lambda + \phi_{\dot{m}}^\lambda \right) \chi_\lambda = 0$$

Staring at this we see a qualitative difference between the mass term and the photon interaction term. Mass requires both right and left components, whereas photon int. does not.

→ **RIL #3:** *The electron mass can be forbidden by assigning  $\psi_L$  different properties than  $\psi_R$ , thereby disallowing  $\psi_L^\dagger \psi_R + \psi_R^\dagger \psi_L$  as an invariant of the theory.*

This is the introduction of chirality into the theory.

$$\mathcal{L} = \frac{1}{4}F_{\mu\nu}F^{\mu\nu} + i\psi_L^\dagger\gamma^\mu(\partial_\mu - ieA_\mu)\psi_L \\ + \psi_R^\dagger\gamma^\mu(\partial_\mu - ieA_\mu)\psi_R + m_e(\psi_L^\dagger\psi_R + \psi_R^\dagger\psi_L)$$

$\psi_L$  differently than  $\psi_R$ . The electric charge for both we must keep at  $-1$ , but we must assign different charges for each under the new symmetry  $G$ . A simple concrete start to this would be to let  $G$  be some new abelian group  $U(1)'$  and assign  $\psi_R$  double the charge of  $\psi_L$ . In other words, our spectrum is

$$\text{Under } U(1)_{EM} \times U(1)' : \\ \psi_L = (-1, -1), \quad \psi_R = (-1, -2)$$

With these charge assignments the  $\bar{\psi}\psi$  term is no longer allowed.

The next step is to regain the electron mass through some other means.

After time I think it is inevitable that people would think of

**RIL #4:** *The electron mass can be reintroduced by adding a condensing scalar field  $\Phi$  with charges  $(0, 1)$  under  $U(1)_{EM} \times U(1)'$  that allows the interaction  $y_e \psi_L^\dagger \Phi \psi_R + h.c.$ , which leads to the electron mass  $m_e = y_e \langle \Phi \rangle$ .*

This is basically the start of the Higgs discussion, but certainly by 1950 with Ginzburg-Landau theory it was in the air to have a complex scalar function order parameter interact with QED!

RIL #4 might look to be the least plausible RIL to some, but I think it is might be the most plausible.

Building on Landau's theory of phase transitions (1937), Ginzburg-Landau theory of superconductivity (1950) looks like the Higgs boson.

power series functional expansion of  $\Psi$  was presented that looks very similar to the Higgs potential of today, and its minimization enables  $|\Psi|^2$  to obtain a non-zero value (vacuum expectation value) below a critical temperature. Furthermore,  $\Psi$  was recognized to have a phase symmetry, equivalent to transformations under  $U(1)_{EM}$ , and the energy density of the system contained the term

$$\frac{1}{2m} \left| -i\hbar \text{grad } \Psi - \frac{e}{c} \mathbf{A} \Psi \right|^2 \quad (10)$$

where it was noted that gauge invariance required replacing  $i\hbar \text{grad}$  to  $i\hbar \text{grad} - (e/c)\mathbf{A}$ , which is what we today call the covariant derivative.

At this point it's all chug and crank. RIL#5 below would be obvious, and hardly an "intellectual leap" given RIL#1-#4:

**RIL #5:** *To avoid massless states in conflict with experiment, the new symmetry  $U(1)'$  should be a local gauge symmetry analogous to  $U(1)_{EM}$ . When  $U(1)'$  is broken by  $\langle\Phi\rangle$  the new photon gets a mass,  $m_{A'}^2 = \frac{1}{2}g'^2\langle\Phi\rangle^2$  and there remains a physical propagating scalar boson with mass  $m_\varphi^2 = 2\lambda\langle\Phi\rangle^2$ .*

This is the most obvious RIL, and easiest to achieve, but its consequences are huge. We started with a non-empirical philosophical criterion (Naturalness) and are led to our first "real" physics prediction.

*RIL #6: Researchers trying to renormalize this theory would discover that  $U(1)'$  has chiral anomalies, and would be forced to add cancelling exotics.*

Now, as it stands our theory is sick, because there are anomalies. The  $U(1)'$ -graviton-graviton and  $U(1)'^3$  anomaly cancelation conditions are not met among the fermions. This necessitates additional fermions in the spectrum that are charged under  $U(1)'$ . There are many possibilities that could be written down. In an exhaustive table among these possibilities would be the following exotic fermions:

$$\text{Exotics : } 6Q'_{1/3} + 3Q'_{4/3} + 3Q'_{-2/3} + 1Q'_{-1}$$

where  $nQ'_q$  means  $n$  copies of fermions with charge  $q$ . These add to our original fermions  $1Q'_{-1} + 1Q'_{-2}$ . These

$$\text{Exotics : } 6Q'_{1/3} + 3Q'_{4/3} + 3Q'_{-2/3} + 1Q'_{-1}$$

where  $nQ'_q$  means  $n$  copies of fermions with charge  $q$ . They add to our original fermions  $1Q'_{-1} + 1Q'_{-2}$ . These charge assignments and multiplicities are exactly those of the Standard Model fermions under twice hypercharge

| Field   | $SU(3)$  | $SU(2)_L$ | $T^3$   | $\frac{Y}{2}$  | $Q = T^3 + \frac{Y}{2}$                                     |
|---|----------|-----------|---|----------------|---|
| $g_\mu^a$ (gluons)  | <b>8</b> | <b>1</b>  | 0   | 0              | 0   |
| $(W_\mu^\pm, W_\mu^0)$                                    | <b>1</b> | <b>3</b>  | $(\pm 1, 0)$  | 0              | $(\pm 1, 0)$  |
| $B_\mu^0$   | <b>1</b> | <b>1</b>  | 0   | 0              | 0   |
| $Q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$          | <b>3</b> | <b>2</b>  | $\begin{pmatrix} \frac{1}{2} \\ -\frac{1}{2} \end{pmatrix}$ | $\frac{1}{6}$  | $\begin{pmatrix} \frac{2}{3} \\ -\frac{1}{3} \end{pmatrix}$ |
| $u_R$   | <b>3</b> | <b>1</b>  | 0   | $\frac{2}{3}$  | $\frac{2}{3}$   |
| $d_R$   | <b>3</b> | <b>1</b>  | 0   | $-\frac{1}{3}$ | $-\frac{1}{3}$  |
| $E_L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}$        | <b>1</b> | <b>2</b>  | $\begin{pmatrix} \frac{1}{2} \\ -\frac{1}{2} \end{pmatrix}$ | $-\frac{1}{2}$ | $\begin{pmatrix} 0 \\ -1 \end{pmatrix}$                     |
| $e_R$   | <b>1</b> | <b>1</b>  | 0   | -1             | -1  |
| $\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$   | <b>1</b> | <b>2</b>  | $\begin{pmatrix} \frac{1}{2} \\ -\frac{1}{2} \end{pmatrix}$ | $\frac{1}{2}$  | $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$                      |
| $\Phi^c = \begin{pmatrix} \phi^0 \\ \phi^- \end{pmatrix}$ | <b>1</b> | <b>2</b>  | $\begin{pmatrix} \frac{1}{2} \\ -\frac{1}{2} \end{pmatrix}$ | $-\frac{1}{2}$ | $\begin{pmatrix} 0 \\ -1 \end{pmatrix}$                     |



## But Naturalness is not yet fully solved! (never ending pursuit...)

A more worrisome objection is that we have replaced one problem with Absolute Naturalness (the electron mass) for another Naturalness problem, and so have not achieved our objective. If the  $\Phi$  mass is  $-\mu^2 \sim m_F^2$  (natural) then  $y_e \sim 10^{-6}$  (unnatural), or phrased differently, if  $-\mu^2 \sim m_e^2$  (natural) then  $-\mu^2/M_F^2 \sim 10^{-6}$  (unnatural). Either way there is a problem with respect to Absolute Naturalness. Thus, all we have done is recast the Naturalness problem rather than solved it.

### Responses:

- (1) Working to solve Naturalness (not necessarily solving it) fruitful.
- (2) Theory space of where to go next is infinite. Naturalness provides discrete path(s) and gateway(s) through which to investigate.
- (3) We are not done in our Naturalness quest: this is the next door.

## Conclusion 1/2

Thus, it is plausible that if Naturalness were religiously held to by researchers in the 1940's/50's, we would have been led to

1. Exotic gauge symmetry and exotic massive gauge boson
2. Higgs boson and Higgs mechanism
3. Exotic charged fermion scenarios, including identifying full SM content as one viable prediction

These are correct and extraordinarily fruitful results from merely taking Naturalness seriously.

But it would nevertheless take many decades to experimentally find all these implications.

(Depressed unimaginative skeptics would gloat during that time.)<sub>50</sub>

## Conclusion 2/2

Same situation today? Naturalness for Higgs leads to many exotic implications.

Analogies: Supersymmetry & Composite Higgs & Etc

1. Solves Naturalness problem
2. Exotics are required for self-consistency
3. Criticism that exchanges one problem (quadratic instability) for another problem (SUSY:  $\mu$  term problem; Xdim: size of compact dimensions)
4. No theorem for how long it will take to discover it (could be decades, and decades more)

Historical considerations suggest taking Naturalness seriously has value. Present circumstances may not warrant its abandonment.