

Positive Geometry of Scattering Amplitudes

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Scattering amplitudes





Quantum Field Theory (QFT)

Our theoretical framework to describe Nature

Compatible with two principles

Special relativity



Quantum mechanics

$$H(t)|\psi(t)\rangle = i\hbar\frac{\partial}{\partial t}|\psi(t)\rangle$$

Standard formulation

(Dirac, Heisenberg, Pauli; Feynman, Dyson, Schwinger)

Fields, Lagrangian, Path integral

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + i\overline{\psi}D\psi - m\overline{\psi}\psi$$

 $\mathcal{D}A \mathcal{D}\psi \mathcal{D}\overline{\psi} \ e^{iS(A,\psi,\overline{\psi},J)}$

Feynman diagrams: pictures of particle interactions
 Perturbative expansion: trees, loops





QFT has passed countless tests in last 70 years

Example: Magnetic dipole moment of electron

Theory: $g_e = 2$ Experiment: $g_e \sim 2$



1928







QFT has passed countless tests in last 70 years

Example: Magnetic dipole moment of electron

1947

Theory: $g_e = 2.00232$ Experiment: $g_e = 2.0023$









QFT has passed countless tests in last 70 years

Example: Magnetic dipole moment of electron
 1957 Theory: $g_e = 2.0023193$ 1972 Experiment: $g_e = 2.00231931$









QFT has passed countless tests in last 70 years

* Example: Magnetic dipole moment of electron Theory: $g_e = 2.0023193044$ 1990 Experiment: $g_e = 2.00231930438$









Dualities

At strong coupling: perturbative expansion breaks





- Surprises: dual to weakly coupled theory
 - Gauge-gauge dualities
 (Montonen-Olive 1977, Seiberg-Witten 1994)
 - Gauge-gravity duality (Maldacena 1997)



Motivation

- Our picture of QFT is incomplete
- Also, tension with gravity and cosmology

If there is a new way of thinking about QFT, it must be seen even at weak coupling

Explicit evidence: scattering amplitudes

Hidden simplicity in scattering amplitudes

Scattering amplitudes

- * Function of spin and external kinematics $\mathcal{M}(p, s, ...)$
- Probability of a given process during a particle collision
- Experimentalists measure cross-section

$$\sigma = \int d\Omega \left| \mathcal{M} \right|^2$$



Colliders at high energies

Proton scattering at high energies





LHC - gluonic factory

Needed: amplitudes of gluons for higher multiplicities

 $qq \rightarrow qq \dots q$ Two helicities: + -







• Status of the art: $gg \rightarrow ggg$

Brute force calculation 24 pages of result



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 $(k_1 \cdot k_4)(\epsilon_2 \cdot k_1)(\epsilon_1 \cdot \epsilon_3)(\epsilon_4 \cdot \epsilon_5)$

New collider

- 1983: Superconducting Super Collider approved
- Energy 40 TeV: many gluons!





✤ Demand for calculations, next on the list: $gg \rightarrow gggg$



- Process $gg \rightarrow gggg$
- ✤ 220 Feynman diagrams, ~100 pages of calculations

GLUONIC TWO GOES TO FOUR

1985: Paper with 14 pages of result

Stephen J. Parke and T.R. Taylor Fermi National Accelerator Laboratory P.O. Box 500, Batavia, IL 60510 U.S.A.

ABSTRACT

The cross section for two gluon to four gluon scattering is given in a form suitable for fast numerical calculations.



* Process $gg \rightarrow gggg$

* 220 Feynman diagrams, \sim 100 pages of calculations





Our result has succesfully passed both these numerical checks. Details of the calculation, together with a full exposition of our techniques, will be given in a forthcoming article. Furthermore, we hope to obtain a simple analytic form for the answer, making our result not only an experimentalist's, but also a theorist's delight.



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Within a year they realized

$$\mathcal{M}_6 = \frac{\langle 12 \rangle^3}{\langle 23 \rangle \langle 34 \rangle \langle 45 \rangle \langle 56 \rangle \langle 61 \rangle}$$

Spinor-helicity variables

$$p^{\mu} = \sigma^{\mu}_{a\dot{a}}\lambda_{a}\tilde{\lambda}_{\dot{a}}$$
$$\langle 12 \rangle = \epsilon_{ab}\lambda^{(1)}_{a}\lambda^{(2)}_{b}$$
$$[12] = \epsilon_{\dot{a}\dot{b}}\tilde{\lambda}^{(1)}_{\dot{a}}\tilde{\lambda}^{(2)}_{\dot{b}}$$

(Mangano, Parke, Xu 1987)



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Within a year they realized

AN AMPLITUDE FOR n GLUON SCATTERING

$$\mathcal{M}_n = \frac{\langle 12 \rangle^3}{\langle 23 \rangle \langle 34 \rangle \langle 45 \rangle \dots \langle n1 \rangle}$$

STEPHEN J. PARKE and T. R. TAYLOR Fermi National Accelerator Laboratory P.O. Box 500, Batavia, IL 60510.

Gauge redundancy

- Where is the problem? Massless particles
- * Particles with spin: gauge redundancy $\epsilon^{\mu} \rightarrow \epsilon^{\mu} + \alpha p^{\mu}$
- Individual Feynman diagrams not gauge invariant Huge cancellations among diagrams

Locality and unitarity

- Redundancy: local interaction picture, off-shell particles
- Two principles manifest:



I) Locality: particles interact point-like

 $\begin{array}{ll} \text{Amplitude:} & 1\\ \text{only poles} & \frac{1}{P^2} \to \infty & P = \sum_{i \in \sigma} p_i \end{array}$

II) Unitarity: sum of probabilities is 1

Amplitude: factorization



Modern methods for amplitudes

Lessons from Parke-Taylor calculation:

Gauge invariance + physical states

- No fields, Lagrangians or path integrals
- Exploit locality and unitarity: fix the amplitude



Recursion relations

(Britto-Cachazo-Feng-Witten 2005)

- Large class of theories at tree-level



Shift momenta + Cauchy formula $p_1 \rightarrow p_1 + zq$ $p_2 \rightarrow p_2 - zq$



34300

20

2485

6

Very efficient method: $gg \to 4g \qquad gg \to 5g \qquad gg \to 6g$ 220 Feynman diagrams 3 **Recursion** relations

Unitarity methods



(Bern-Dixon-Kosower)

Iterative use of the unitary cut



- Generate basis of integrals, fixing coefficients from cuts
- Tremendous success
 in calculations in 1990-today

Example: Four point 3-loop amplitudes in supersymmetric Yang-Mills theory and gravity



Unitarity methods



(Bern-Dixon-Kosower)

QCD background at LHC

BlackHat collaboration



Huge efficiency in NLO calculations



Used by CMS in comparison to data, March 2014

Toy model

- This is a great success; is there a deeper structure?
- Time-proven method: study a toy model first <u>Wish list:</u>
 - Four-dimensional interacting theory
 - Close to the real world (QCD) as much as possible
 - Ability to generate plenty of explicit results

Maximally supersymmetric Yang-Mills theory in planar limit

(Brink-Scherk-Schwarz 1977)

- Conformal, convergent series
- Great toy model for QCD
 - <u>Tree-level amplitudes identical</u>
 - Loop amplitudes simpler, structures similar
 - But, no confinement :(
- Past: new methods for amplitudes originated here

Many faces of the theory

Useful playground for many theoretical ideas



Simple amplitudes

Comparison: Feynman diagrams vs unitary methods

 $gg \rightarrow gg$ Number of graphs $gg \rightarrow gg$ Number of graphs graphs

What is the amplitude?

New definition of the amplitude

Standard: Function consistent with locality and unitarity

Our goal: Different definition

- No fields, Lagrangians, path integrals
- Unitarity, locality emergent from other principles
- Powerful method for calculations

Prelude



Volume of polyhedron

(Hodges 2009)

- * New kinematical variables momentum twistors $Z \in \mathbb{C}^3$
- * Tree-level process: $gg \rightarrow 5g$
- Comparison of two calculations of recursion relations



Evidence for a new structure



"Conjecture"

Amplitudes are volumes of *some regions* in *some space*

The Amplituhedron

(Arkani-Hamed, JT 2013)

Strategy

- Simple intuitive geometric ideas: use equations
- Generalization:
 More complicated geometry
 - Higher dimensions
- Same equations persist











Amplituhedron conjecture

* Volume of $A_{n,k,\ell}$:

Amplitudes in maximally supersymmetric Yang-Mills theory *n* number of particles

k helicity information

number of loops

 $\ell = 0$: Amplitudes of gluons in QCD

Consistency check: Locality and Unitarity

Explicit checks against reference theoretical data

Volume of the space

- Set of inequalities: Volume = differential form
- * Simple examples: x > 0: $\operatorname{Vol} = \frac{dx}{x}$ y > 0, x > 0: $\operatorname{Vol} = \frac{dx}{x}\frac{dy}{y}$ y > x > 0: $\operatorname{Vol} = \frac{dx}{x}\frac{dy}{y-x}$
- * Amplituhedron for amplitude $gg \rightarrow gg$
 - Nice interpretation: Configuration of vectors on a plane
 - Easy to state, hard to solve "High school problem"

Positive quadrant











Positive quadrant

Vectors

$$\vec{a}_1 = \begin{pmatrix} x_1 \\ y_1 \end{pmatrix} \quad \vec{b}_1 = \begin{pmatrix} z_1 \\ w_1 \end{pmatrix}$$
$$\vec{a}_2 = \begin{pmatrix} x_2 \\ y_2 \end{pmatrix} \quad \vec{b}_2 = \begin{pmatrix} z_2 \\ w_2 \end{pmatrix}$$



- Positive quadrant
- * Vectors $\vec{a}_1, \vec{a}_2, \vec{a}_3$ $\vec{b}_1, \vec{b}_2, \vec{b}_3$
- Conditions

$$(\vec{a}_1 - \vec{a}_2) \cdot (\vec{b}_1 - \vec{b}_2) \le 0$$
$$(\vec{a}_1 - \vec{a}_3) \cdot (\vec{b}_1 - \vec{b}_3) \le 0$$
$$(\vec{a}_2 - \vec{a}_3) \cdot (\vec{b}_2 - \vec{b}_3) \le 0$$



- Positive quadrant
- * Vectors $\vec{a}_1, \vec{a}_2, \dots, \vec{a}_\ell \quad \vec{b}_1, \vec{b}_2, \dots, \vec{b}_\ell$
- Conditions

 $(\vec{a}_i - \vec{a}_j) \cdot (\vec{b}_i - \vec{b}_j) \le 0$

for all pairs i, jLet me know if you solve it!



$$\operatorname{Vol}\left(\ell\right) = \dots$$

Positivity

In the definition of Amplituhedron $\begin{array}{c} \mathcal{Y} = \mathcal{C} \cdot Z \\ \mathcal{Y} = \mathcal{Y} \\ \mathcal{Y} = \mathcal{Y} = \mathcal{Y} \\ \mathcal{Y} = \mathcal{Y} \\$

Positivity: crucial property of geometry

- Locality, unitarity, even planarity derived from it
- Hidden symmetry of this theory (Yangian) manifest

Question for mathematicians

How to make big positive matrices? For the case $\ell = 0$ solved by Alexander Postnikov in 2006

Positive Grassmannian $G_+(k, n)$

Gluing procedure

Construct big positive matrix from small ones

(* * *) (* * *)

Gluing preserves positivity of minors

Arbitrary graph: positive matrix



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 Triangulation of Amplituhedron: set of these diagrams

Permutations

"Basis" of these matrices: labeled by permutations

 $\begin{pmatrix} * & * & * & * \\ & * & * & * \end{pmatrix} \longleftrightarrow$



 \leftrightarrow (3, 4, 1, 2)

Juggling patterns

Allen Knutson (Cornell U.) 1990-1995 world record in juggling (12 balls)

POSITROID VARIETIES I: JUGGLING AND GEOMETRY

ALLEN KNUTSON, THOMAS LAM, AND DAVID E SPEYER







Deligne table

Other appearances of graphs

Cluster variables associated with each graph

 $\begin{array}{c} f_{1} \\ f_{4} \\ f_{6} \\ f_{3} \\ f_{3} \end{array} \leftrightarrow \begin{array}{c} f_{1} \\ f_{4} \\ f_{4} \\ f_{6} \\ f_{7} \\$

(Fock-Goncharov 2003)

Dual graphs: quivers, Seiberg duality, shalow water waves,...





(Kodama-Williams 2011)

On-shell diagrams

Same diagrams: radically different interpretation

Physical on-shell processes: product of 3pt amplitudes



Detailed study of this connection

arXiv:1212.5605 [pdf, other]

Scattering Amplitudes and the Positive Grassmannian

Nima Arkani-Hamed, Jacob L. Bourjaily, Freddy Cachazo, Alexander B. Goncharov, Alexander Postnikov, Jaroslav Trnka Comments: 158 pages, 264 figures Subjects: High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Combinatorics (math.CO)



Physics vs geometry

Dynamical particle interactions in 4-dimensions





 Static geometry in high dimensional space



Real process at LHC

At the intersection

Fascinating connection between fields which have





 CENTRE
DE RECHERCHES
MATHÉMATIQUES
 Positive Grassmannians: Applications to integrable systems and
super Yang-Mills scattering amplitudes
July 27-31, 2015

 MOME
 SCHEDULE

 REGISTER
 CONMODATION

* Scattering $gg \rightarrow ggg \dots gg$ in our toy model

$$\mathcal{M}_{n}^{tree} = \frac{\langle 12 \rangle^{3}}{\langle 23 \rangle \langle 34 \rangle \langle 45 \rangle \dots \langle n1 \rangle}$$

* Scattering $gg \rightarrow ggg \dots gg$ in our toy model

$$\mathcal{M}_{n} = \mathcal{M}_{n}^{tree} \left\{ \begin{array}{c} 1 \\ 1 \end{array} + \\ Tree-level \end{array} \right\}$$

(1985)

* Scattering $gg \rightarrow ggg \dots gg$ in our toy model $\mathcal{M}_n = \mathcal{M}_n^{tree} \left\{ \begin{array}{c} 1 \\ \end{array} + \sum_{ij} \\ \end{array} \right\}_{ij}$ Tree-level One-loop (1994)(1985)





* Scattering $gg \rightarrow ggg \dots gg$ in our toy model





(Arkani-Hamed, Bourjaily, Cachazo, JT 2010)



Outlook: Beyond the toy model

- Amplituhedron: Geometric picture for amplitudes
- Next step: non-planar, gravity, string amplitudes, <u>QCD</u>
- Evidence beyond the toy model
 - On-shell diagram: non-planar, no susy (Arkani-Hamed, Bourjaily, Cachazo, Postnikov, JT 2014)



- Amplituhedron-type construction beyond the planar limit (Arkani-Hamed, Bourjaily, Cachazo, JT 2014) (Bern, Hermann, Litsey, Stankowicz, JT 2014)
- Connection to EFTs (NLσM, DBI, Galileon) via soft limits (Cheung, Kampf, Novotny, JT 2014)

Thank you for your attention