# Positive Geometry of Scattering Amplitudes 

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UC Davis, March 2, 2015


## Scattering amplitudes



## Quantum Field Theory (QFT)

\% Our theoretical framework to describe Nature

* Compatible with two principles


## Special relativity

## Quantum mechanics



$$
H(t)|\psi(t)\rangle=i \hbar \frac{\partial}{\partial t}|\psi(t)\rangle
$$

# Standard formulation 

(Dirac, Heisenberg, Pauli; Feynman, Dyson, Schwinger)
$\because$ Fields, Lagrangian, Path integral
$\mathcal{L}=-\frac{1}{4} F_{\mu \nu} F^{\mu \nu}+i \bar{\psi} D \psi-m \bar{\psi} \psi \quad \int \mathcal{D} A \mathcal{D} \psi \mathcal{D} \bar{\psi} e^{i S(A, \psi, \bar{\psi}, J)}$

* Feynman diagrams: pictures of particle interactions Perturbative expansion: trees, loops



## Great success of QFT

\% QFT has passed countless tests in last 70 years

* Example: Magnetic dipole moment of electron

Theory: $\quad g_{e}=2$
1928
Experiment: $\quad g_{e} \sim 2$


## Great success of QFT

$\because$ QFT has passed countless tests in last 70 years

* Example: Magnetic dipole moment of electron

Theory: $g_{e}=2.00232$
1947
Experiment: $g_{e}=2.0023$


## Great success of QFT

$\%$ QFT has passed countless tests in last 70 years
\% Example: Magnetic dipole moment of electron
1957 Theory: $g_{e}=2.0023193$
1972 Experiment: $g_{e}=2.00231931$


## Great success of QFT

\% QFT has passed countless tests in last 70 years

* Example: Magnetic dipole moment of electron

Theory: $g_{e}=2.0023193044$
1990
Experiment: $\quad g_{e}=2.00231930438$


## Dualities

\% At strong coupling: perturbative expansion breaks

\% Surprises: dual to weakly coupled theory

- Gauge-gauge dualities
(Montonen-Olive 1977, Seiberg-Witten 1994)
- Gauge-gravity duality
(Maldacena 1997)



## Motivation

\% Our picture of QFT is incomplete

* Also, tension with gravity and cosmology

If there is a new way of thinking about QFT, it must be seen even at weak coupling
\% Explicit evidence: scattering amplitudes

Hidden simplicity in scattering amplitudes

## Scattering amplitudes

$\because$ Function of spin and external kinematics $\quad \mathcal{M}(p, s, \ldots)$
\% Probability of a given process during a particle collision
\% Experimentalists measure cross-section

$$
\sigma=\int d \Omega|\mathcal{M}|^{2}
$$



## Colliders at high energies

$\therefore$ Proton scattering at high energies


LHC - gluonic factory
$\because$ Needed: amplitudes of gluons for higher multiplicities

$$
g g \rightarrow g g \ldots g
$$

Two helicities: + -


## Early 80s

$\because$ Status of the art: $g g \rightarrow g g g$

Brute force calculation<br>24 pages of result



## New collider

$\because$ 1983: Superconducting Super Collider approved
$\therefore$ Energy 40 TeV : many gluons!

\% Demand for calculations, next on the list: $g g \rightarrow g g g g$

## Parke-Taylor formula


$\because$ Process $g g \rightarrow g g g g$
$\therefore 220$ Feynman diagrams, $\sim 100$ pages of calculations

GLUONIC TWO GOES TO FOUR

: 1985: Paper with 14 pages of result

Stephen J. Parke and T.R. Taylor Fermi National Accelerator Laboratory P.0. Box 500, Batavia, IL 60510 U.S.A.

ABSTRACT
The cross section for two gluon to four gluon scattering is given in a form suitable for fast numerical calculations.

## Parke-Taylor formula


$\because$ Process $g g \rightarrow g g g g$
$\therefore 220$ Feynman diagrams, $\sim 100$ pages of calculations







## Parke-Taylor formula



Our result has succesfully passed both these numerical checks.
Details of the calculation, together with a full exposition of our techniques, will be given in a forthcoming article. Furthermore, we hope to obtain a simple analytic form for the answer, making our result not only an experimentalist's, but also a theorist's delight.

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\% Within a year they realized

$$
\mathcal{M}_{6}=\frac{\langle 12\rangle^{3}}{\langle 23\rangle\langle 34\rangle\langle 45\rangle\langle 56\rangle\langle 61\rangle}
$$

Spinor-helicity variables

$$
\begin{aligned}
p^{\mu} & =\sigma_{a \dot{a}}^{\mu} \lambda_{a} \tilde{\lambda}_{\dot{a}} \\
\langle 12\rangle & =\epsilon_{a b} \lambda_{a}^{(1)} \lambda_{b}^{(2)} \\
{[12] } & =\epsilon_{\dot{a} \dot{b}} \tilde{\lambda}_{\dot{a}}^{(1)} \tilde{\lambda}_{\dot{b}}^{(2)}
\end{aligned}
$$

(Mangano, Parke, Xu 1987)

## Parke-Taylor formula



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## : Within a year they realized

AN AMPLITUDE FOR n GLUON SCATTERING

$$
\mathcal{M}_{n}=\frac{\langle 12\rangle^{3}}{\langle 23\rangle\langle 34\rangle\langle 45\rangle \ldots\langle n 1\rangle}
$$

## Gauge redundancy

* Where is the problem? Massless particles
\% Particles with spin: gauge redundancy $\epsilon^{\mu} \rightarrow \epsilon^{\mu}+\alpha p^{\mu}$
\% Individual Feynman diagrams not gauge invariant Huge cancellations among diagrams


## Locality and unitarity

$\because$ Redundancy: local interaction picture, off-shell particles
\% Two principles manifest:

I) Locality: particles interact point-like
$\begin{gathered}\text { Amplitude: } \\ \text { only poles }\end{gathered} \quad \frac{1}{P^{2}} \rightarrow \infty \quad P=\sum_{i \in \sigma} p_{i}$
II) Unitarity: sum of probabilities is 1

Amplitude: factorization



## Modern methods for amplitudes

* Lessons from Parke-Taylor calculation:


## Gauge invariance + physical states

\% No fields, Lagrangians or path integrals

* Exploit locality and unitarity: fix the amplitude



## Recursion relations

$\because$ Large class of theories at tree-level

* Tree-level unitarity

$\therefore$ Shift momenta + Cauchy formula

$$
\begin{aligned}
& p_{1} \rightarrow p_{1}+z q \\
& p_{2} \rightarrow p_{2}-z q
\end{aligned}
$$


\% Very efficient method:

$$
\begin{gathered}
g g \rightarrow 4 g \\
220 \\
3
\end{gathered}
$$

| $g g \rightarrow 5 g$ | $g g \rightarrow 6 g$ |
| :---: | :---: |
| 2485 | 34300 |
| 6 | 20 |

## Unitarity methods


(Bern-Dixon-Kosower)

$\because$ Iterative use of the unitary cut

*Generate basis of integrals, fixing coefficients from cuts
\% Tremendous success in calculations in 1990-today

Example: Four point 3-loop amplitudes in supersymmetric Yang-Mills theory and gravity


# Unitarity methods 


(Bern-Dixon-Kosower)
$\because$ QCD background at LHC
\% BlackHat collaboration


* Huge efficiency in NLO calculations


Used by CMS in comparison to data, March 2014

## Toy model

$\because$ This is a great success; is there a deeper structure?

* Time-proven method: study a toy model first

Wish list:

- Four-dimensional interacting theory
- Close to the real world (QCD) as much as possible
- Ability to generate plenty of explicit results


# Maximally supersymmetric Yang-Mills theory in planar limit 

(Brink-Scherk-Schwarz 1977)
© Conformal, convergent series
$\therefore$ Great toy model for QCD

- Tree-level amplitudes identical
- Loop amplitudes simpler, structures similar
- But, no confinement :(
: Past: new methods for amplitudes originated here


## Many faces of the theory

\% Useful playground for many theoretical ideas


## Simple amplitudes

$\because$ Comparison: Feynman diagrams vs unitary methods

$$
g g \rightarrow g g
$$

Number of graphs


87 vs 1


## What is the amplitude?

## New definition of the amplitude

$\because$ Standard: Function consistent with locality and unitarity

* Our goal: Different definition
- No fields, Lagrangians, path integrals
- Unitarity, locality emergent from other principles
- Powerful method for calculations

Prelude

## Volume of polyhedron

$\because$ New kinematical variables - momentum twistors

$$
Z \in \mathbb{C}^{3}
$$

* Tree-level process: $g g \rightarrow 5 g$
* Comparison of two calculations of recursion relations



## Evidence for a new structure


(Arkani-Hamed, Bourjaily, Cachazo, Hodges, JT 2010)

(Arkani-Hamed, Cachazo, Cheung, Kaplan 2009)

## "Conjecture"

Amplitudes are volumes of some regions in some space

## The Amplituhedron

(Arkani-Hamed, JT 2013)

## Strategy

$\because$ Simple intuitive geometric ideas: use equations
$\because$ Generalization: - More complicated geometry

- Higher dimensions
* Same equations persist


## Road to Amplituhedron



## Road to Amplituhedron



## Road to Amplituhedron

Start:
Point inside a convex polygon


## Road to Amplituhedron



## Road to Amplituhedron



## Amplituhedron conjecture

$\because$ Volume of $\mathcal{A}_{n, k, \ell}$ :
Amplitudes in maximally supersymmetric Yang-Mills theory

$$
\ell=0: \text { Amplitudes of gluons in } \mathrm{QCD}
$$

* Consistency check: Locality and Unitarity

number of particles<br>helicity information



\% Explicit checks against reference theoretical data

## Volume of the space

$\because$ Set of inequalities: Volume $=$ differential form
$\therefore$ Simple examples: $\quad x>0: \quad \mathrm{Vol}=\frac{d x}{x}$
$y>0, x>0: \quad \mathrm{Vol}=\frac{d x}{x} \frac{d y}{y} \quad y>x>0: \quad \mathrm{Vol}=\frac{d x}{x} \frac{d y}{y-x}$

* Amplituhedron for amplitude $g g \rightarrow g g$
- Nice interpretation: Configuration of vectors on a plane
- Easy to state, hard to solve - "High school problem"


## High school problem $\quad g g \rightarrow g g$

\% Positive quadrant


## High school problem $\quad g g \rightarrow g g$

\% Positive quadrant
\% Vectors

$$
\vec{a}_{1}=\binom{x_{1}}{y_{1}} \quad \vec{b}_{1}=\binom{z_{1}}{w_{1}}
$$



$$
\operatorname{Vol}(1)=\frac{d x_{1}}{x_{1}} \frac{d y_{1}}{y_{1}} \frac{d z_{1}}{z_{1}} \frac{d w_{1}}{w_{1}}=
$$

## High school problem $\quad g g \rightarrow g g$

\% Positive quadrant
$\therefore$ Vectors

$$
\begin{aligned}
& \vec{a}_{1}=\binom{x_{1}}{y_{1}} \quad \vec{b}_{1}=\binom{z_{1}}{w_{1}} \\
& \vec{a}_{2}=\binom{x_{2}}{y_{2}} \quad \vec{b}_{2}=\binom{z_{2}}{w_{2}}
\end{aligned}
$$


$[\operatorname{Vol}(1)]^{2}=\frac{d x_{1}}{x_{1}} \frac{d y_{1}}{y_{1}} \frac{d z_{1}}{z_{1}} \frac{d w_{1}}{w_{1}} \frac{d x_{2}}{x_{2}} \frac{d y_{2}}{y_{2}} \frac{d z_{2}}{z_{2}} \frac{d w_{2}}{w_{2}}=\square \times \square$

## High school problem $\quad g g \rightarrow g g$

: Positive quadrant
$\therefore$ Vectors

$$
\begin{aligned}
& \vec{a}_{1}=\binom{x_{1}}{y_{1}} \quad \vec{b}_{1}=\binom{z_{1}}{w_{1}} \\
& \vec{a}_{2}=\binom{x_{2}}{y_{2}} \quad \vec{b}_{2}=\binom{z_{2}}{w_{2}}
\end{aligned}
$$


\% Impose:
$\left(\vec{a}_{2}-\vec{a}_{1}\right) \cdot\left(\vec{b}_{2}-\vec{b}_{1}\right) \leq 0 \quad \phi>90^{\circ}$
Subset of configurations allowed

## High school problem $\quad g g \rightarrow g g$

* Positive quadrant
$\therefore$ Vectors

$$
\begin{aligned}
& \vec{a}_{1}=\binom{x_{1}}{y_{1}} \quad \vec{b}_{1}=\binom{z_{1}}{w_{1}} \\
& \vec{a}_{2}=\binom{x_{2}}{y_{2}} \quad \vec{b}_{2}=\binom{z_{2}}{w_{2}}
\end{aligned}
$$


$\operatorname{Vol}(2)=\frac{d x_{1}}{x_{1}} \frac{d y_{1}}{y_{1}} \frac{d z_{1}}{z_{1}} \frac{d w_{1}}{w_{1}} \frac{d x_{2}}{x_{2}} \frac{d y_{2}}{y_{2}} \frac{d z_{2}}{z_{2}} \frac{d w_{2}}{w_{2}}\left[\frac{\vec{a}_{1} \cdot \vec{b}_{2}+\vec{a}_{2} \cdot \vec{b}_{1}}{\left(\vec{a}_{2}-\vec{a}_{1}\right) \cdot\left(\vec{b}_{2}-\vec{b}_{1}\right)}\right]$

## High school problem $\quad g g \rightarrow g g$

* Positive quadrant
$\therefore$ Vectors

$$
\begin{aligned}
& \vec{a}_{1}=\binom{x_{1}}{y_{1}} \quad \vec{b}_{1}=\binom{z_{1}}{w_{1}} \\
& \vec{a}_{2}=\binom{x_{2}}{y_{2}} \quad \vec{b}_{2}=\binom{z_{2}}{w_{2}}
\end{aligned}
$$


$\operatorname{Vol}(2)=$ $\square$
$\square$

## High school problem $\quad g g \rightarrow g g$

\% Positive quadrant
$\therefore$ Vectors
$\vec{a}_{1}, \vec{a}_{2}, \vec{a}_{3} \quad \vec{b}_{1}, \vec{b}_{2}, \vec{b}_{3}$
$\because$ Conditions

$$
\begin{aligned}
& \left(\vec{a}_{1}-\vec{a}_{2}\right) \cdot\left(\vec{b}_{1}-\vec{b}_{2}\right) \leq 0 \\
& \left(\vec{a}_{1}-\vec{a}_{3}\right) \cdot\left(\vec{b}_{1}-\vec{b}_{3}\right) \leq 0 \\
& \left(\vec{a}_{2}-\vec{a}_{3}\right) \cdot\left(\vec{b}_{2}-\vec{b}_{3}\right) \leq 0
\end{aligned}
$$


$\operatorname{Vol}(3)=$


## High school problem $\quad g g \rightarrow g g$

\% Positive quadrant
$\because$ Vectors

$$
\vec{a}_{1}, \vec{a}_{2}, \ldots, \vec{a}_{\ell} \quad \vec{b}_{1}, \vec{b}_{2}, \ldots, \vec{b}_{\ell}
$$

$\therefore$ Conditions

$$
\left(\vec{a}_{i}-\vec{a}_{j}\right) \cdot\left(\vec{b}_{i}-\vec{b}_{j}\right) \leq 0
$$

for all pairs $i, j$
Let me know if you solve it!

$\operatorname{Vol}(\ell)=\ldots \ldots$.

## Positivity

$\because$ In the definition of Amplituhedron


Amplituhedron

$$
\left.\begin{array}{c|ll}
\text { Positive matrices: } & \mid & * \\
\text { Minors are positive } & * & *
\end{array} \right\rvert\,>0
$$


\% Positivity: crucial property of geometry

- Locality, unitarity, even planarity derived from it
- Hidden symmetry of this theory (Yangian) manifest


## Question for mathematicians

## How to make big positive matrices?

For the case $\ell=0$ solved by Alexander Postnikov in 2006
Positive Grassmannian $G_{+}(k, n)$

$$
\left(\begin{array}{llllll}
* & * & * & * & * & * \\
* & * & * & * & * & * \\
* & * & * & * & * & *
\end{array}\right) \quad\left|\begin{array}{ccc}
* & * & * \\
* & * & * \\
* & * & *
\end{array}\right|>0
$$

## Gluing procedure

\% Construct big positive matrix from small ones


Gluing preserves positivity of minors

* Arbitrary graph: positive matrix


$$
\left(\begin{array}{lllll}
* & * & * & * & * \\
* & * & * & * & *
\end{array}\right)\left(\begin{array}{llllll}
* & * & * & * & * & * \\
* & * & * & * & * & * \\
* & * & * & * & * & *
\end{array}\right)
$$



* Triangulation of Amplituhedron: set of these diagrams


## Permutations

* "Basis" of these matrices: labeled by permutations

$$
\left(\begin{array}{llll}
* & * & * & * \\
* & * & * & *
\end{array}\right) \leftrightarrow
$$


$(3,4,1,2)$
: Juggling patterns
Allen Knutson (Cornell U.)
1990-1995 world record in juggling (12 balls)

POSITROID VARIETIES I: JUGGLING AND GEOMETRY

ALLEN KNUTSON, THOMAS LAM, AND DAVID E SPEYER



Deligne table

## Other appearances of graphs

: Cluster variables associated with each graph

(Fock-Goncharov 2003)

* Dual graphs: quivers, Seiberg duality, shalow water waves,...

(Kodama-Williams 2011)


## On-shell diagrams

## * Same diagrams: radically different interpretation

$\because$ Physical on-shell processes: product of 3pt amplitudes


* Detailed study of this connection

```
arXiv:1212.5605 [pdf, other]
```


## Scattering Amplitudes and the Positive Grassmannian

Nima Arkani-Hamed, Jacob La. Bourjaily, Freddy Cachazo, Alexander B. Goncharov, Alexander Postnikov, Jaroslav Trnka Comments: 158 pages, 264 figures
Subjects: High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Combinatorics (math.CO)

## Physics vs geometry

* Dynamical particle interactions in 4-dimensions

$\because$ Static geometry in high dimensional space



## At the intersection

*Fascinating connection between fields which have


Walter Burke Institute Workshop
Grassmannian Geometry
of Scattering Amplitudes


California Institute of Technology
December 8-12, 2014
Invited speakers: Nima Arkani-Hamed, Till Bargheer, Zvi Bern, Jacob Bourjaily, Johannes Broedel, Lance Dixon, Nick Early, Davide Forcella, Sebastian Franco, Johannes Broedel, Lance Dixon, Nick Early, Davide Forcella, Sebastian Franco,
Daniele Galloni, Song He, Johannes Henn, Andrew Hodges, Thomas Lam, Daniele Galloni, Song He, Johannes Henn, Andrew Hodges, Thomas Lam,
Sangmin Lee, Tomasz Lukowski, Lionel Mason, Timothy Olson, David Speyer Matthias Staudacher, Anastasia Volovich, Lauren Williams, Dan Xie.

Organizers: Hirosi Ooguri, Jaroslav Trnka https://burkeinstitute.caltech.edu/workshops/Grassmannian2014


## Back to Parke-Taylor formula

: Scattering $\quad g g \rightarrow g g g \ldots g g$
in our toy model

$$
\mathcal{M}_{n}^{\text {tree }}=\frac{\langle 12\rangle^{3}}{\langle 23\rangle\langle 34\rangle\langle 45\rangle \ldots\langle n 1\rangle}
$$

## Back to Parke-Taylor formula

$\because$ Scattering $\quad g g \rightarrow g g g \ldots g g$
in our toy model
$\mathcal{M}_{n}=\mathcal{M}_{n}^{\text {tree }}\left\{\begin{array}{l}1 \\ 1\end{array}+\right.$
Tree-level
(1985)

## Back to Parke-Taylor formula

$\because$ Scattering $\quad g g \rightarrow g g g \ldots g g$
in our toy model


## Back to Parke-Taylor formula

$\because$ Scattering $\quad g g \rightarrow g g g \ldots g g$
in our toy model


|  |  |  |
| :---: | :---: | :---: |
|  |  |  |
|  | $\frac{5}{5}$ | $\frac{5+5}{5+5}$ |
| $\begin{aligned} & \frac{\pi}{x}= \\ & \frac{\pi}{4}= \end{aligned}$ |  | of result |

## Back to Parke-Taylor formula

$\because$ Scattering $g g \rightarrow g g g \ldots g g$
in our toy model



## Back to Parke-Taylor formula

$\because$ Scattering $\quad g g \rightarrow g g g \ldots g g$
in our toy model


## Back to Parke-Taylor formula

$\because$ Scattering $g g \rightarrow g g g \ldots g g$
in our toy model


## Outlook: Beyond the toy model

* Amplituhedron: Geometric picture for amplitudes
* Next step: non-planar, gravity, string amplitudes, QCD
: Evidence beyond the toy model
- On-shell diagram: non-planar, no susy
(Arkani-Hamed, Bourjaily, Cachazo, Postnikov, JT 2014)

- Amplituhedron-type construction beyond the planar limit
(Arkani-Hamed, Bourjaily, Cachazo, JT 2014) (Bern, Hermann, Litsey, Stankowicz, JT 2014)
- Connection to EFTs (NL $\sigma$ M, DBI, Galileon) via soft limits
(Cheung, Kampf, Novotny, JT 2014)

Thank you for your attention

