

# Studying challenging theories with the superconformal bootstrap

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1304.1803, 1306.3228, [1312.5344](#), [1404.1079](#), 1408.6522, 1412.7541, [150X.XXXX](#)

with various subsets of:

M. Lemos, P. Liendo, W. Peelaers, L. Rastelli, B. C. van Rees, A. Sen

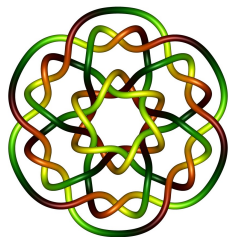
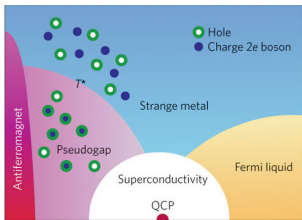
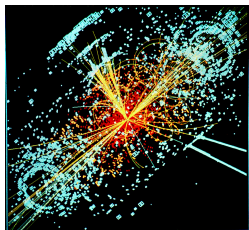
February 19, 2015



<sup>†</sup> Supported by the Frank and Peggy Taplin Fellowship

# Quantum field theory

Quantum field theory is ubiquitous in modern theoretical physics (and mathematics).



However, quantum field theory is not a *technology* – can't just take it off the shelf and turn the crank.

This is more than a technical problem. There are hints that *we are missing something significant*.

# Fields, Lagrangians, path integrals...

QFT is usually formulated as a theory of quantum fields:

$$\varphi(x), \quad \psi_\alpha(x), \quad A_\mu(x), \quad \dots,$$

Write a Lagrangian (subject to some conditions), compute path integral:

$$\mathcal{L}[\varphi] = \partial_\mu \varphi \partial^\mu \varphi + m^2 \varphi + g^2 \varphi^4 + \dots$$

$$Z = \int [D\varphi] e^{\frac{i}{\hbar} S}, \quad S = \int d^D x \mathcal{L}[\varphi(x)].$$

Many subtleties (regularization, renormalizability), but the story is basically established and useful.

Disclaimer: In this talk, QFT  $\equiv$  Lorentz-invariant, unitary QFT

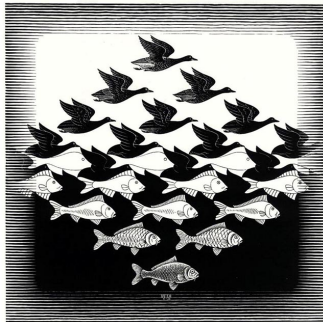
# Fields, Lagrangians, path integrals...

So {quantum field theories} = {(UV complete) Lagrangians}?

# *The* **FIELD** *Is Not Enough* **007<sup>th</sup>**

**Duality:** Some theories admit multiple Lagrangian descriptions.

strongly coupled catfish



weakly coupled goose

weakly coupled catfish

strongly coupled goose

E.g., Electric-Magnetic duality in  $\mathcal{N} = 4$  super Yang-Mills (+ many more...)

Duality connects to deep mathematics

Mirror symmetry (2d)  
Geometric Langlands (4d) [Gukov; Kapustin; Witten]

So {quantum field theories} = {Lagrangians}/Duality

...

Lagrangians like coordinate charts?

# *The* **FIELD** *Is Not Enough* **007<sup>™</sup>**

**Non-Lagrangian theories:** some theories seem to admit *no Lagrangian description*.



Existence deduced indirectly, often using *decoupling limits of string/M-theory*.

Such theories pose a serious conceptual challenge

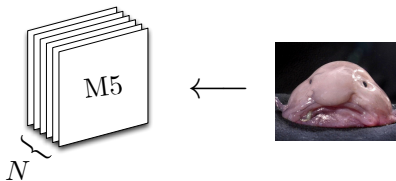
In the rest of this talk, I'm going to describe a conservative approach to understanding a particularly interesting class of non-Lagrangian theories using algebraic methods.



## $(2, 0)$ theory in $d = 6$ [Seiberg (1996)]

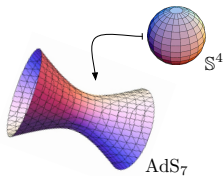
No (interacting) continuum Lagrangian QFTs in  $d > 4$  dimensions.

Nevertheless, *six-dimensional interacting QFTs exist*.



*Conformally invariant:*  $SO(5, 1) \rightarrow SO(6, 2)$

*Maximally supersymmetric:*  $SO(6, 2) \rightarrow OSp(8|4)$



Holographic dual description for  $N \rightarrow \infty$ .  
Can't compute  $1/N$  corrections.

## $(2, 0)$ theory in $d = 6$ [Seiberg (1996)]

These theories appear to play be fairly important:

- ▶ No superconformal symmetry in  $d > 6$  [Nahm (1978)].

(Speculation: no interacting QFT in  $d > 6$ ?).

- ▶ The “theory of M5 branes” (what is M-theory?)

- ▶  $d \leq 4$  landscape populated by compactifications.

[Gaiotto (2008)]

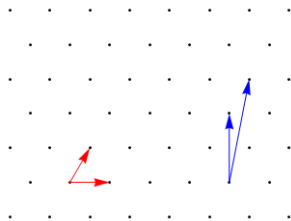
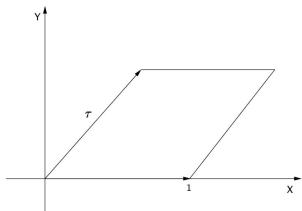
[Bah, Beem, Bobev, Wecht (2011, 2012); Beem, Dimofte, Pasquetti (2012)]

[...]

*Explains duality* in lower dimensions.

# “Explaining” duality in four dimensions [Witten (1995)]

$$(2, 0)_N \text{ on } \mathbb{R}^4 \times T^2 \xrightarrow{\mathbb{R}} SU(N) \mathcal{N} = 4 \text{ SYM on } \mathbb{R}^4$$
$$\text{Modular parameter of } T^2 \longrightarrow \tau = \frac{4\pi i}{g^2} + \frac{\theta}{2\pi} .$$



$$T^2_{\tau} \equiv T^2_{\tau'} , \quad \tau' = \frac{a\tau + b}{c\tau + d} \quad \text{where} \quad ad - bc = 1 .$$

Modular invariance  $\longrightarrow$  S-duality

## So what is the $(2, 0)_N$ theory?

Liberal: Low energy limit of  $N$  coincident M5 branes.

Conservative: List of local operators with superconformally-covariant correlation functions.

Moderate: Mostly conservative, but occasionally cross the aisle.

# The conservative approach



# Consequences of conformal symmetry

Operators in *conformal families*:  $\{\mathcal{O}_{\Delta,\ell}, \partial\mathcal{O}_{\Delta,\ell}, \partial^2\mathcal{O}_{\Delta,\ell}\}$

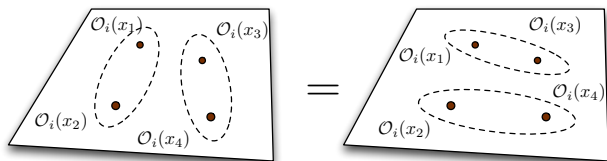
Algebraic structure: *convergent OPE*

$$\mathcal{O}_j(y) \mathcal{O}_i(x) = \sum_k c_{ij}^k(x-y) \mathcal{O}_k(x)$$

Coefficients functions fixed by *three-point functions of primaries*.

$n$ -point functions determined from spectrum and three-point functions (*CFT data*)

# Consequences of conformal symmetry



$$u = \frac{x_{12}^2 x_{34}^2}{x_{14}^2 x_{23}^2}$$

$$v = \frac{x_{13}^2 x_{24}^2}{x_{14}^2 x_{23}^2}$$

CFT data is nontrivially constrained by *crossing symmetry*:

$$\sum_{\mathcal{O}_j} c_{iij}^2 G_{\Delta_j, \ell_j}(u, v) = \sum_{\mathcal{O}_j} c_{iij}^2 G_{\Delta_j, \ell_j}(v, u) .$$

One equation for each four-point function.

Conformal bootstrap: just solve these equations!

[Ferrara, Gatto, Grillo 1971-1975; Polyakov 1974]

[N.B. need *infinite number of conformal families*]

# 21st century bootstrap: convex optimization

In  $d \geq 3$ , no major progress until [Rattazzi, Rychkov, Tonni, Vichi (2008)] – numerical approach.

Roughly speaking, the technology is as follows:

- ▶ Rewrite crossing symmetry as *sum rule with positive coefficients*:

$$\sum_{\mathcal{O}_i} c_i^2 \left( G_{\Delta_i, \ell_i}(u, v) - G_{\Delta_i, \ell_i}(v, u) \right) = 1 .$$

- ▶ Make assumptions about spectrum – limits the basis of functions on LHS.
- ▶ Prove that no solution can exist:

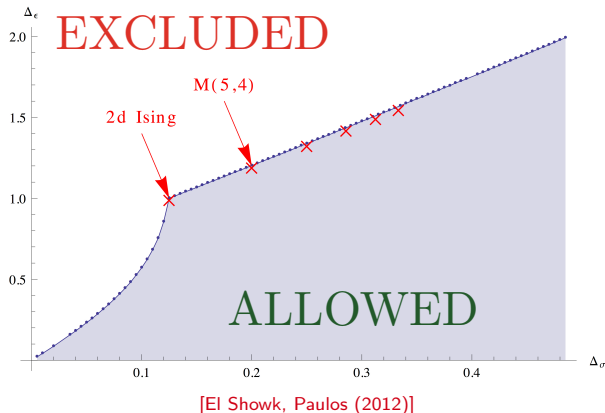
$$\text{i.e., } f_i''(0) > 0 \implies \sum_i c_i^2 f_i(x) \neq 1 .$$

- ▶ Keyword: “convex optimization”



# 21st century bootstrap: convex optimization

In  $d \geq 3$ , no major progress until [Rattazzi, Rychkov, Tonni, Vichi (2008)] – numerical approach.



Theories on boundary can have CFT data *systematically reconstructed*.

# Consequences of supersymmetry [Beem, et. al. (2013)]

*Superconformal families:*  $\{\mathcal{O}, Q\mathcal{O}, \dots, Q^{16}\mathcal{O}, \text{descendants}\}$

Interesting supersymmetric operators:  $Q^n \mathcal{O}_{BPS} = 0, n < 16$ .

Expect infinitely many BPS operators

(cp. ordinary CFT: only conserved currents)

OPE algebra admits *truncation* involving *only BPS operators*

$$“\mathcal{O}_{BPS}^{(i)} \times \mathcal{O}_{BPS}^{(j)} = \sum_k \mathcal{O}_{BPS}^{(k)} .”$$

BPS algebra *much simpler* than full operator algebra.

Disclaimer: actual truncation very complicated

## Consequences of supersymmetry [Beem, et. al. (2013)]

For  $(2, 0)$  theories, truncation requires operators lie in  $\mathbb{C}_{[z, \bar{z}]^2}^2 \subset \mathbb{R}^6$ :

$$\mathcal{O}_i(z)\mathcal{O}_j(w) \sim \sum_k \frac{c_{ij}^k \mathcal{O}_k(w)}{(z-w)^{h_i+h_j-h_k}}$$

Known as *chiral algebras* (or *vertex algebras*) – appear in 2d CFT.

Crossing symmetry is nontrivial, but *tractable*.

analogy: complex analysis vs. real analysis

## (2, 0) chiral algebra [Beem, Rastelli, van Rees (2014)]

Know enough about BPS spectrum to *solve chiral algebra bootstrap* completely.

$$\text{BPS chiral algebra} = W_N \text{ algebra}$$

First calculable correlation functions in the (2, 0) theory at finite  $N$

At large  $N$  we can algebraically verify predictions from holography:

$$c_{k_1 k_2 k_3} = \frac{2^{2\alpha-2}}{(\pi N)^{3/2}} \Gamma\left(\frac{k_1 + k_2 + k_3}{2}\right) \left( \frac{\Gamma(\frac{k_{123}+1}{2})\Gamma(\frac{k_{231}+1}{2})\Gamma(\frac{k_{312}+1}{2})}{\sqrt{\Gamma(2k_1-1)\Gamma(2k_2-1)\Gamma(2k_3-1)}} \right).$$

(Research project): Finite  $N$  – *quantum gravity corrections* in M-theory.

# Analytic $\implies$ Numerical [Beem, Rastelli, van Rees (2013)]

BPS correlators alone are a big improvement, but can we do more?

Analytic results for BPS operators sets the stage for numerical analysis:

Solution to chiral algebra bootstrap

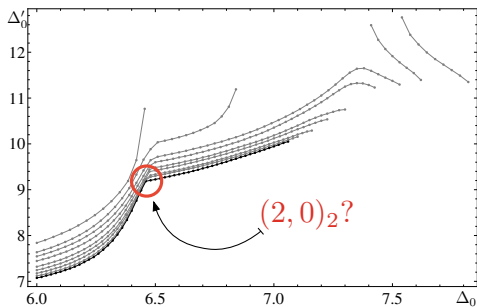


CFT data for BPS operators



Crossing symmetry for non-BPS operators

# Cornering the $(2, 0)_2$ theory [Beem, Lemos, Rastelli, van Rees (in progress)]



Numerical bootstrap results for  $(2, 0)$  theory.

“Interesting point” on the boundary seems to correspond to the  $(2, 0)_2$  theory.

$(2, 0)$  theory = Ising model for the 21st century?

# Where we are

- ▶ Rich algebraic structures connected to SCFTs.
  - $6d \mathcal{N} = (2, 0) \implies$  chiral algebra [Beem, Rastelli, van Rees (2014)]
  - $4d \mathcal{N} \geq 2 \implies$  chiral algebra [Beem et. al. (2013)]
  - $3d \mathcal{N} = 4 \implies$  deformation quantization [Beem, Peelaers, Rastelli (in progress)]
  
- ▶ Can compute BPS correlators in non-Lagrangian theories.  
[Beem, Rastelli, van Rees (2014); Beem, Peelaers, Rastelli, van Rees (2014)]
  
- ▶ Strong indications  $(2, 0)$  theory numerically accessible.  
[Beem, Lemos, Rastelli, van Rees (in progress)]

## Future directions

- ▶ Explore “protected” algebraic structures –  
*many connections to interesting mathematics.*
- ▶ Are numerically accessible theories analytically special?
- ▶ Right mathematical framework for the bootstrap?





Thanks!