## Challenges for theoretical cosmology

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Fundamental Questions in Cosmology
UC Davis May 2013

Multipole moment, $\ell$


## Challenges for Cosmic Inflation (eternal inflation)

"Anything that can happen will happen infinitely many times" (A. Guth)

1) Measure Problems
2) Problems defining probabilit
3) Problems/hidden assumptions re initial conditions
$\rightarrow$ problem claiming generic predictions about state
$\rightarrow$ cannot claim "solution to cosmological
problems"
$\rightarrow$ Related to $2^{\text {nd }}$ law, low $S$ start
4) Yet, Successful fits to data

## Slow rolling of inflaton



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Multipole moment, $\ell$


Challenges: Answer these questions re your theories \& beliefs:

1) Do you predict the observed state of the universe to be likely or natural? (And do you care?)
2) Do you treat infinities rigorously?
3) Do you require a probability tooth fairy?

## 1) Do you predict the observed state of the universe to be likely or natural? (And do you care?)

- Beware hidden assumptions about initial conditions (often related to $2^{\text {nd }}$ law: $S>0 \rightarrow S$ initially small $\rightarrow$ starting in limited part of phase space)

Gibbons \& Turok
Carroll \& Tam
Shiffren \& Wald
Penrose

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(just as true of cyclic models)

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In general: Need a quantitative theory for your starting point (inflation, cyclic, whatever) to make this claim.

Attempts I know to create this rigor have led to surprises.

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| $\mathbf{X}$ | Y |
| :--- | :--- |
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| X | $\mathbf{Y}$ |
| :--- | :--- |
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| Volume of inflated <br> regions | Probability for starting <br> inflation |
| Entropy | Probability of starting <br> a cyclic universe |
| Number of observers <br> (in my theory) who see <br> a universe like ours | The infinitely many <br> other observers who <br> see something totally <br> different |

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Need more rigor:

- Hernley, AA \& Dray (2013) $\leftrightarrow$ Guth toy model
- AA \& Sorbo (2004)

Increasing the level of rigor usually reveals significant hidden assumptions that amount to tuning of initial conditions.

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Non-Quantum probabilities in a toy model:

$$
\begin{aligned}
& U=A \otimes B \quad A:\left\{|1\rangle^{A},|2\rangle^{A}\right\} \quad B:\left\{|1\rangle^{B},|2\rangle^{B}\right\} \\
& U:\{|1\rangle\rangle,|12\rangle,|21\rangle,|22\rangle\} \quad|i j\rangle \equiv|i\rangle^{A}|j\rangle^{B}
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Possible Measurements $\leftrightarrows$ Projection operators:
Measure A only:

$$
\hat{P}_{i}^{A}=\left(|i\rangle^{A A}\langle i|\right) \otimes \mathbf{1}^{B}=[|i 1\rangle\langle i 1|+|i 2\rangle\langle i 2|]
$$

Measure $B$ only:

$$
\hat{P}_{i}^{B}=\left(|i\rangle^{B B}\langle i|\right) \otimes \mathbf{1}^{A}=[|1 i\rangle\langle 1 i|+|2 i\rangle\langle 2 i|]
$$

Measure entire $U$ :

$$
\hat{P}_{i j} \equiv|i j\rangle\langle i j|
$$

## 3) Do yoı BUT: It is impossible to construct a projection operator for the case where you do not know whether it is <br> Non-Quan $A$ or $B$ that is being measured.

Could Write

$$
U=A^{\prime} \quad \hat{P}_{i}=p_{A} \hat{P}_{i}^{A}+p_{B} \hat{P}_{i}^{B}
$$

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BUT: It is impossible to construct a projection operator for the case where you do not know whether it $\mathrm{j} \quad$ Classical $A$ or $B$ that is being measured.

Could Write

$$
\hat{P}_{i}=p_{A}^{2} \hat{P}_{i}^{A}+p_{B}^{\widehat{A}} \widehat{P_{i}^{B}}
$$

$$
\hat{P}_{i} \hat{P}_{j} \neq \delta_{i j} \hat{P}_{j}
$$

## Does not

represent a
hents $\longleftrightarrow$ Projection operators:
quantum

## measurement

Measure o only:

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## Classical Probabilities

$$
U=A
$$

Does not represent a quantum

## measurement

Measure $D$ only:
Measure entire $U$ :

Could Write

$$
\hat{P}_{i}=p_{A}^{*} \hat{P}_{i}^{A}+p_{B}^{-} \widehat{P_{i}^{B}}
$$

to measure
A, B

$$
\hat{P}_{i} \hat{P}_{j} \neq \delta_{i j} \hat{P}_{j}
$$

$$
\text { hents } \leftrightarrow \mathrm{p} \text { Page: The }
$$

multiverse requires this (are you in pocket universe A or B?)
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## Classical Probabilities

$$
U=A
$$

Does not represent a quantum measurement Measure donly:

Measure entire $U$ :

Could Write

$$
\hat{P}_{i}=p_{A}^{4} \hat{P}_{i}^{\Lambda}+p_{B}^{-} \widehat{P_{i}^{B}}
$$

to measure
A, B

$$
\hat{P}_{i} \hat{P}_{j} \neq \delta_{i j} \hat{P}_{j}
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Page: The
multiverse requires this (are you in pocket universe A or B?)

- All everyday probabilities are quantum probabilities
- One should not use ideas from everyday probabilities to justify probabilities that have been proven to have no quantum origin

AA \& D. Phillips 2012

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## Page: The

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## All everyday probabilities are quantun

 probabilities- One should not use ideas from everyday probabilities to justify probabilities that have been proven to have no quantum origin

AA \& D. Phillips 2012


## Quantum effects in a billiard gas



$$
\left(\begin{array}{cccccccc}
0^{\circ} & 8_{0}^{0} & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & \varepsilon_{0} & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0
\end{array}\right)
$$

## Quantum effects in a billiard gas




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$$
\Delta b=\delta x_{\perp}+\frac{\delta p_{\perp}}{m} \Delta t
$$



## Quantum effects in a billiard gas



$$
\Delta b=\delta x_{\perp}+\frac{\delta p_{\perp}}{m} \Delta t=\sqrt{2}\left(a+\frac{\hbar}{2 a} \frac{l}{m \bar{v}}\right) \quad \psi \propto \exp \left(\frac{-x^{2}}{2 a^{2}}\right)
$$

$\Delta \bar{b}$


## Quantum effects in a billiard gas



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$$
\begin{aligned}
& \Delta b=\delta x_{\perp}+\frac{\delta p_{\perp}}{m} \Delta t=\sqrt{2}\left(a+\frac{\hbar}{2 a} \frac{l}{m \bar{v}}\right) \\
& \quad \min 2^{3 / 2}\left(\frac{\hbar l}{2 m \bar{v}}\right) \equiv \sqrt{l \lambda_{d B} / 2}
\end{aligned}
$$



## Quantum effects in a billiard gas

After $n$ collisions:

$$
\Delta b_{n}=\Delta b(1+2 l / r)^{n}
$$



## Quantum effects in a billiard gas


$n_{Q}$ is the number of collisions so that $\Delta b_{n_{Q}}=r$
(full quantum uncertainty as to which is the next collision)

$$
n_{Q}=-\frac{\log \left(\frac{\Delta b}{r}\right)}{\log \left(1+\frac{2 l}{r}\right)}
$$

## $n_{Q}$ for a number of physical systems

(all units MKS)

|  | $r$ | $l$ | $m$ | $\bar{v}$ | $\lambda_{d B}$ | $\Delta b$ | $n_{Q}$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Air |  |  |  |  |  |  |  |
| Water |  |  |  |  |  |  |  |
| Billiards |  |  |  |  |  |  |  |
| Bumper <br> Car |  |  |  |  |  |  |  |

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| Air |  |  |  |  |  |  |  |
| Water |  |  |  |  |  |  |  |
| Billiards |  |  |  |  |  |  |  |
| Bumper <br> Car | 1 | 2 | 150 | 0.5 | $1.4 \times 10^{-36}$ | $3.4 \times 10^{-18}$ | 25 |



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| :--- | :---: | :---: | :---: | :---: | :--- | :--- | :--- |
| Air |  |  |  |  |  |  |  |
| Water |  |  |  |  |  |  |  |
| Billiards | 0.029 | 1 | 0.16 | 1 | $6.6 \times 10^{-34}$ | $5.1 \times 10^{-17}$ | 8 |
| Bumper <br> Car | 1 | 2 | 150 | 0.5 | $1.4 \times 10^{-36}$ | $3.4 \times 10^{-18}$ | 25 |



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| :--- | :---: | :---: | :---: | :---: | :--- | :--- | :---: |
| Air |  |  |  |  |  |  |  |
| Water | $3.0 \times 10^{-10}$ | $5.4 \times 10^{-10}$ | $3 \times 10^{-26}$ | 460 | $7.6 \times 10^{-12}$ | $1.3 \times 10^{-10}$ | 0.6 |
| Billiards | 0.029 | 1 | 0.16 | 1 | $6.6 \times 10^{-34}$ | $5.1 \times 10^{-17}$ | 8 |
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| :--- | :---: | :---: | :---: | :---: | :--- | :--- | :---: |
| Air | $1.6 \times 10^{-10}$ | $3.4 \times 10^{-7}$ | $4.7 \times 10^{-26}$ | 360 | $6.2 \times 10^{-12}$ | $2.9 \times 10^{-9}$ | -0.3 |
| Water | $3.0 \times 10^{-10}$ | $5.4 \times 10^{-10}$ | $3 \times 10^{-26}$ | 460 | $7.6 \times 10^{-12}$ | $1.3 \times 10^{-10}$ | 0.6 |
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| Billiards | 0.029 | 1 | 0.16 | 1 | $6.6 \times 10^{-34}$ | 5.2 | Quantum <br> at every <br> collision |
| Bumper <br> Car | 1 | 2 | 150 | 0.5 | $1.4 \times 10^{-36}$ | 3. | 3. |




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(all units MKS)


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(independent of "interpretation")

## $n_{Q}$ for a number of physical systems

(all units MKS)


# An important role for Brownian motion: Uncertainty in neuron transmission times 



## Analysis of coin flip

$$
\begin{aligned}
& \delta t_{f}=\delta t_{n} \times\left(\frac{v_{h}}{v_{h}+v_{f}}\right) \\
& \delta t_{t}=\sqrt{2} \delta t_{f} \\
& f=\frac{4 v_{f}}{\pi d}
\end{aligned}
$$

$$
\delta N=f \delta t_{t}=0.5
$$

Using:

Coin diameter $=d$


$$
\begin{aligned}
& \delta t_{n} \approx 1 \mathrm{~ms} \quad v_{h}=v_{f}=5 \mathrm{~m} / \mathrm{s} \\
& d=0.01 \mathrm{~m}
\end{aligned}
$$

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## 50-50 coin flip probabilities are <br> a derivable quantum result

Using:

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NB: Coin flip is "at the margin" of classical vs quantum control: Increasing $d$ or deceasing $v_{h}$ can reduce $\delta N$ substantially

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- Which theories really do require classical probabilities not yet resolved rigorously.


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## Challenges for Cosmologists:

1) Find a foundation for inflation (or an alternative theory) that can be *well* tested with modern data. Meet the "Challenges for theorists"
2) Only then can we claim to resolve the famous cosmological puzzles (Monopoles already OK).
3) Still, already have great narrative about the origin of perturbations. (Should we be happy with that?)
4) Run risk of being "stuck" like standard model of particle physics has been (so far).


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4) Run risk of being "stuck" like standard model of particle physics has been (so far). We can do better!

Challenges: Answer these questions re your theories \& beliefs:

1) Do you predict the observed state of the universe to be likely or natural? (And do you care?)
2) Do you treat infinities rigorously?
3) Do you require a probability tooth fairy?

Multipole moment, $\ell$


Challenges: Answer these questions re your theories \& beliefs:

1) Do you predict the observed state of the universe to be likely or natural2 (And do you care?1

## YES

2) Do you treat infinities rigorously?
3) Do you require a probability tooth fairy?
