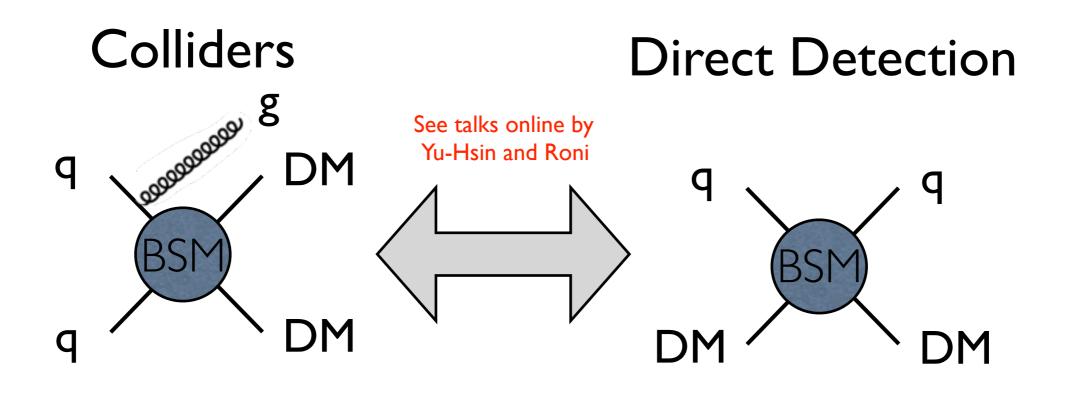
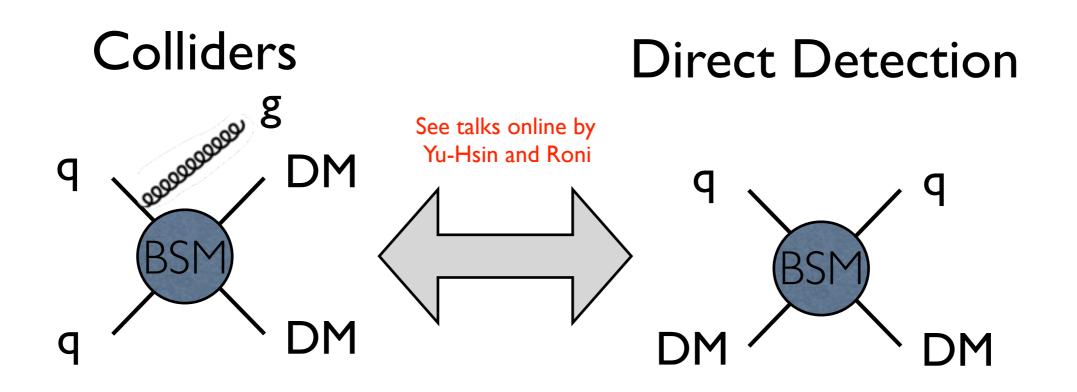
### MULTIPURPOSE MONOJETS ATTHE LHC: NEUTRINOS & DARK MATTER Ian M. Shoemaker April 13, 2012 [arXiv: 1111.5331, PLB in press] Alex Friedland, Michael Graesser, IMS, and Luca Vecchi

[arXiv: 1112.5457] IMS and Luca Vecchi

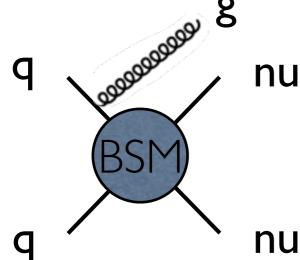
# More complementarity



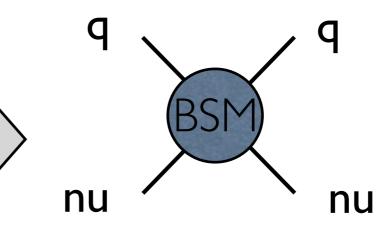
# More complementarity



Colliders



Oscillation experiments



Why do you care?

- Neutrinos certainly exist in Nature.
- Solar neutrino hints.
- Weakly constrained so far.
- The LHC can set world's best limits.

#### Generalizing Fermi

PHYSICAL REVIEW D

**VOLUME 17, NUMBER 9** 

1 MAY 1978

#### Neutrino oscillations in matter

L. Wolfenstein

Carnegie-Mellon University, Pittsburgh, Pennsylvania 15213 (Received 6 October 1977; revised manuscript received 5 December 1977)

The effect of coherent forward scattering must be taken into account when considering the oscillations of neutrinos traveling through matter. In particular, for the case of massless neutrinos for which vacuum oscillations cannot occur, oscillations can occur in matter if the neutral current has an off-diagonal piece connecting different neutrino types. Applications discussed are solar neutrinos and a proposed experiment involving transmission of neutrinos through 1000 km of rock.

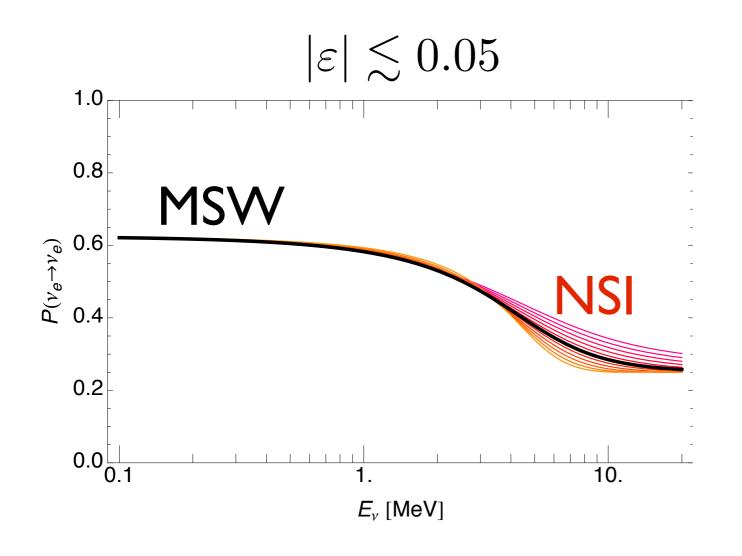
$$\mathcal{L}_{\text{NSI}} = -2\sqrt{2}G_F \epsilon_{\alpha\beta}^{fP} (\overline{\nu}_{\alpha}\gamma^{\rho}\nu_{\beta}) (\overline{f}\gamma_{\rho}Pf)$$

$$\bigwedge_{\text{Neutrino Flavor}} f = \text{SM fermion}$$

$$P = L, R$$

# Laid the foundation for the MSW effect and pointed out that NSI can modify neutrino propagation.

# Solar Neutrinos and the LHC: a UV-IR duality?



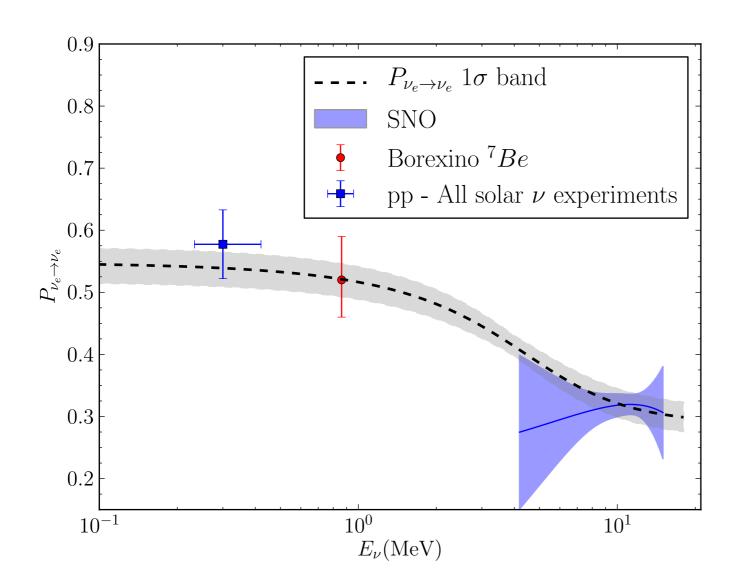
Recently, both SNO and Super-K lowered thresholds to discover the MSW "upturn:"

#### neither see it

Borexino recently targeted 8B neutrinos and also found *no evidence*.

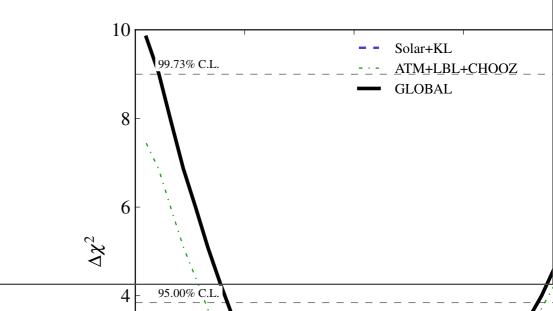
Combined >2σ discrepancy. Palazzo [arXiv:1101.3875]

# A SNO-ball's chance?



Dashed line: best fit to LMA solution

Blue: SNO solar neutrino data





"This could be the discovery of the century. Depending, of course, on how far down it goes."



"This could be the discovery of the century. Depending, of course, on how far down it goes."

#### Not all MET is created equal

Jets + MET searches can bound many invisible things, like ADD gravitons, DM, (sterile) neutrinos, unparticles.

1) Yet, only SM neutrinos can *interfere* with the SM:

 $\sigma(pp/p\overline{p} \to j + \text{MET}) = \sigma_{\text{SM}} + \epsilon \sigma_{\text{int}} + \epsilon^2 \sigma_{\text{NSI}}$ 

2) SM neutrinos have nonzero electroweak charge.

**MadGraph** for parton-level signal.

Pass to **Pythia** for hadronization/ showering and ultimately analysis.

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 $N_{BSM} < 192$ @ 95% CL

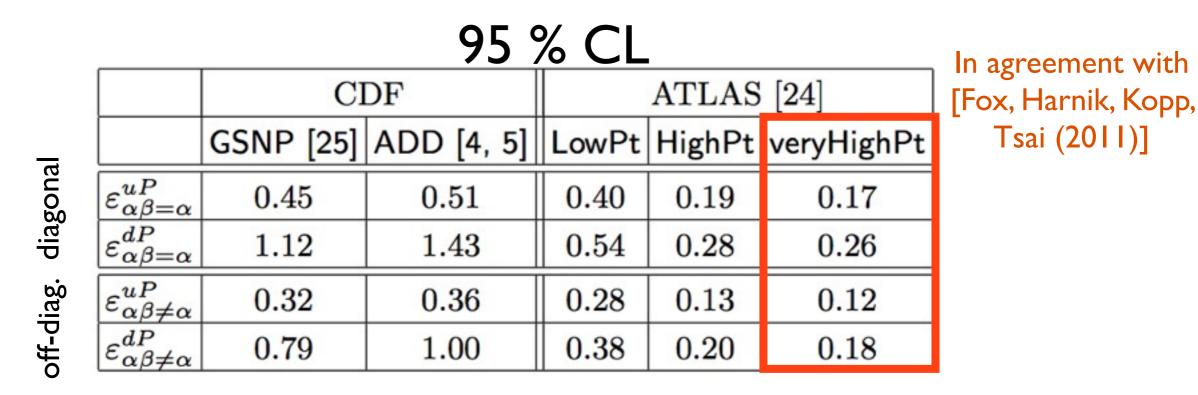
 $\mathcal{L}_{\rm NSI} = -2\sqrt{2}G_F \,\epsilon^{fP}_{\alpha\beta} (\overline{\nu}_{\alpha}\gamma^{\rho}\nu_{\beta}) (\overline{f}\gamma_{\rho}Pf)$ 

		95 % CL					
		CI	ATLAS [24]				
		GSNP [25]	ADD [4, 5]	LowPt	HighPt	veryHighPt	
diagonal	$\varepsilon^{uP}_{\alpha\beta=lpha}$	0.45	0.51	0.40	0.19	0.17	
diag	$\varepsilon^{dP}_{\alpha\beta=lpha}$	1.12	1.43	0.54	0.28	0.26	
liag.	$\varepsilon^{uP}_{\alpha\beta\neqlpha}$	0.32	0.36	0.28	0.13	0.12	
off-diag.	$\varepsilon^{dP}_{\alpha\beta\neqlpha}$	0.79	1.00	0.38	0.20	0.18	

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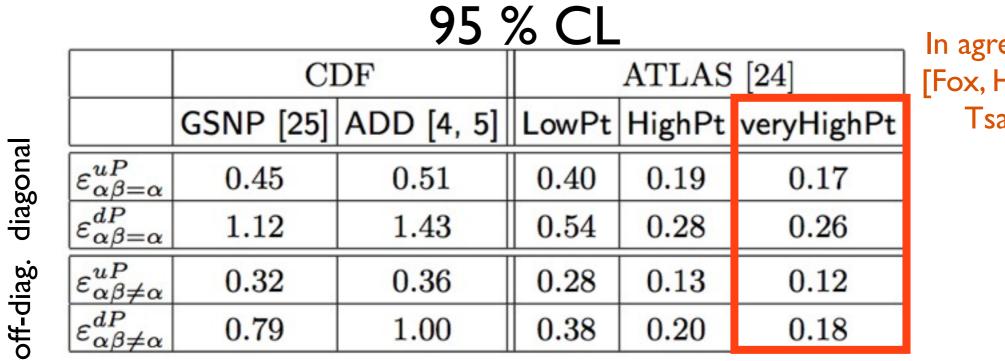
	95 % CL						In agreement with
		CDF		ATLAS [24]			[Fox, Harnik, Kopp,
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•Up-quark couplings more constrained.

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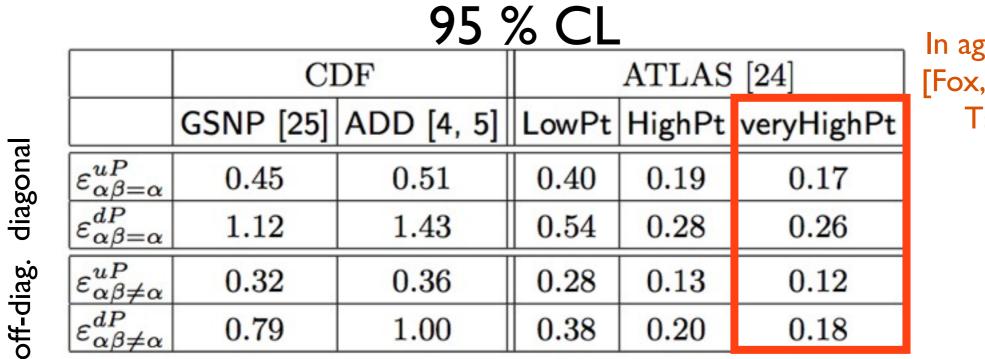


In agreement with [Fox, Harnik, Kopp, Tsai (2011)]

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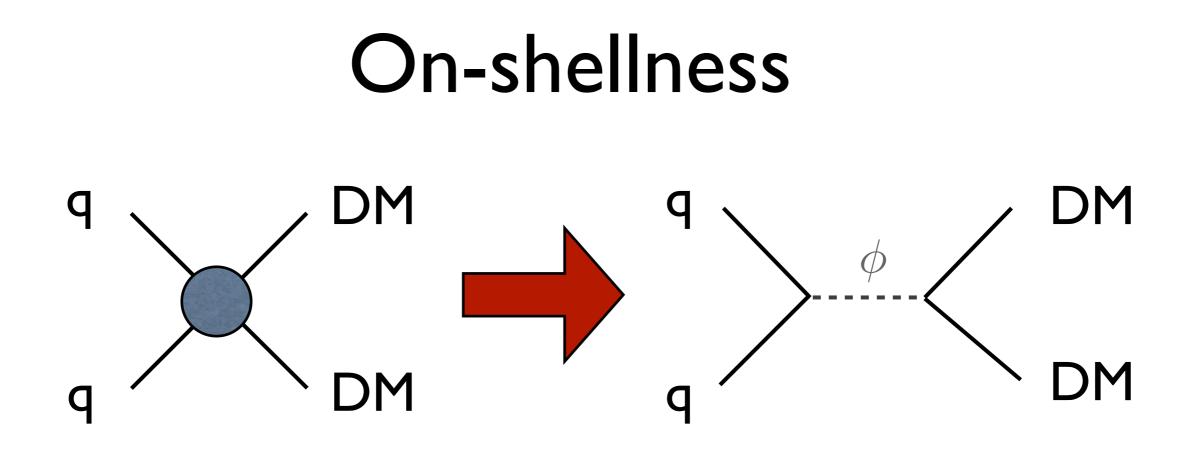


In agreement with [Fox, Harnik, Kopp, Tsai (2011)]

- •Up-quark couplings more constrained.
- •Off-diagonal couplings stronger by  $(\sqrt{2})$ .
- •Most stringent bounds to date on electron and tau-type NSIs.

E.g. previously  $\varepsilon_{uR}^{ee} < 0.7$  from DIS at CHARM (Davidson et al., 2003).

Friday, April 13, 2012

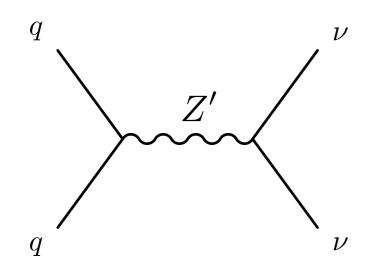


More model-independent, BUT only valid as long as the new physics scale is large compared to LHC energies.

#### 

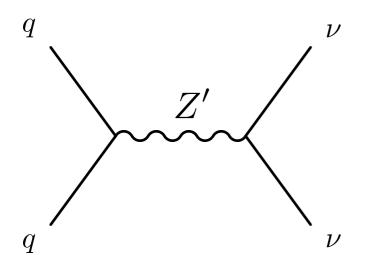
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What if it's not?

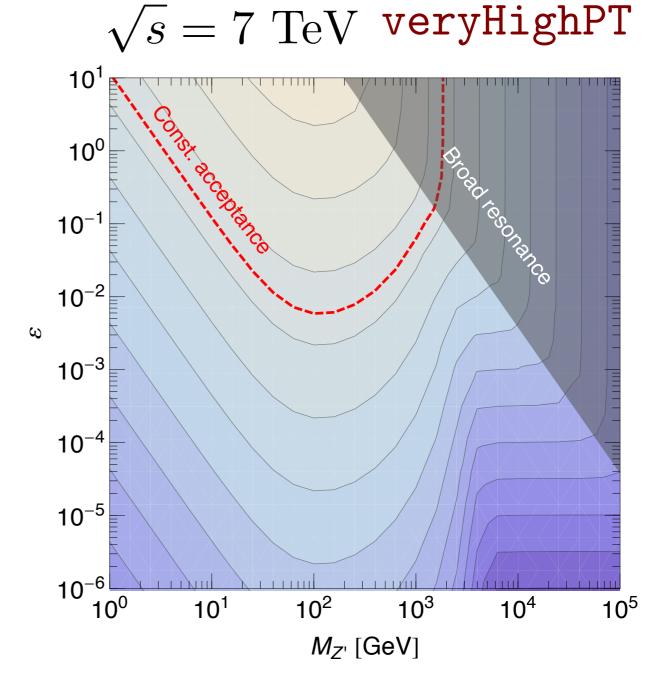


Minimal width, Z' only couples to one quark flavor, chirality and a neutrino pair.

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 $q \qquad \nu$   $Z' \qquad \qquad \nu$   $q \qquad \nu$ 

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10 10<sup>0</sup>  $10^{-1}$ 10<sup>-2</sup> ω  $10^{-3}$ Contact  $10^{-4}$ 10<sup>-5</sup>  $10^{-6}$ 10<sup>2</sup> 10<sup>4</sup> 10<sup>0</sup> 10<sup>3</sup> 10<sup>1</sup> 10<sup>5</sup>  $M_{Z'}$  [GeV]

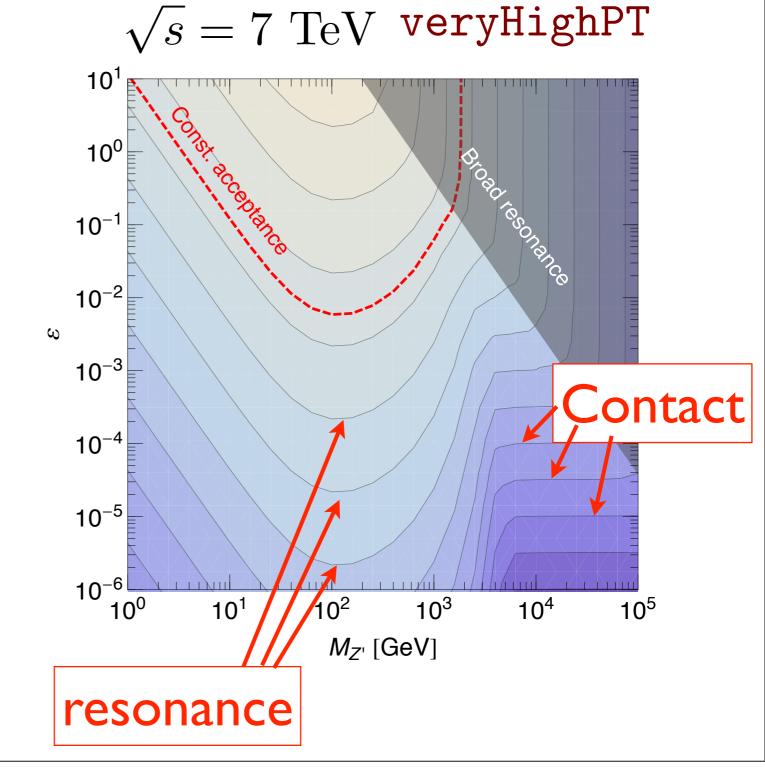
 $\sqrt{s} = 7 {
m TeV} {
m veryHighPT}$ 

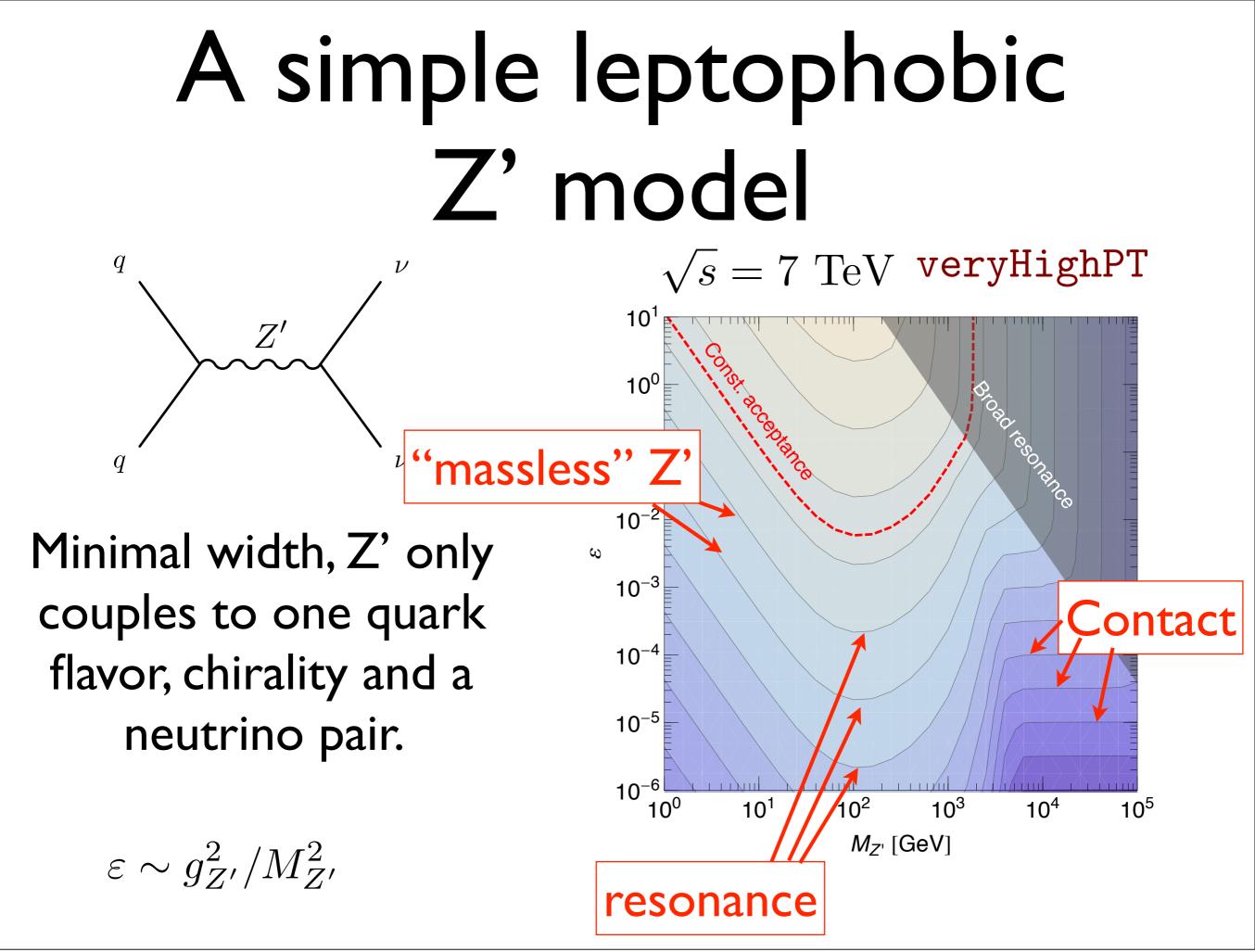
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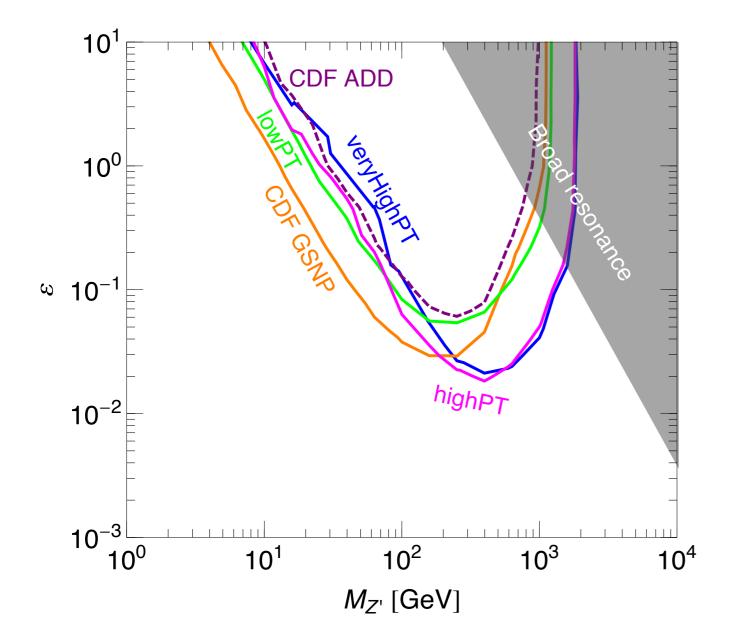
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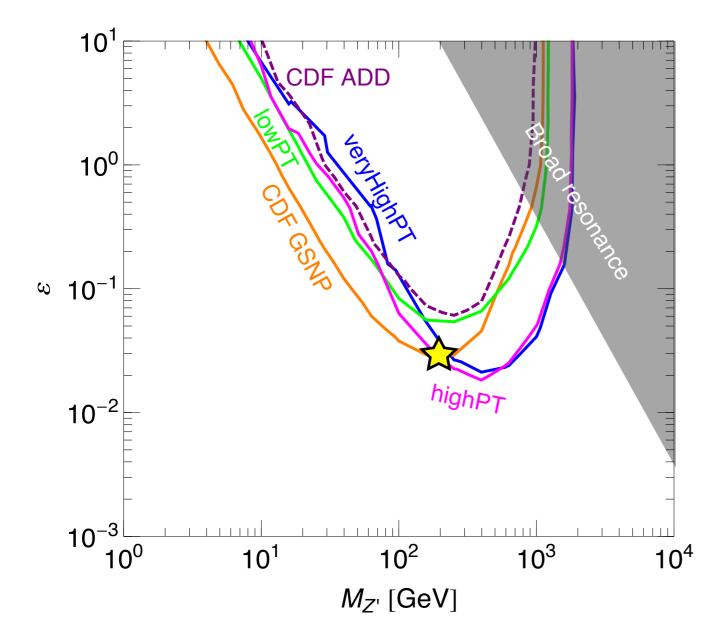




# Beating the LHC from beyond the grave



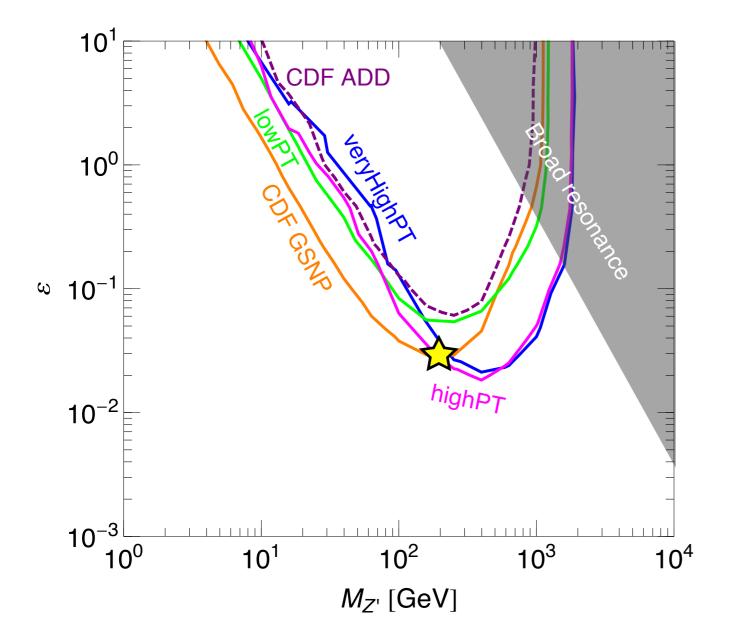
# Beating the LHC from beyond the grave



-Tevatron data is more constraining for

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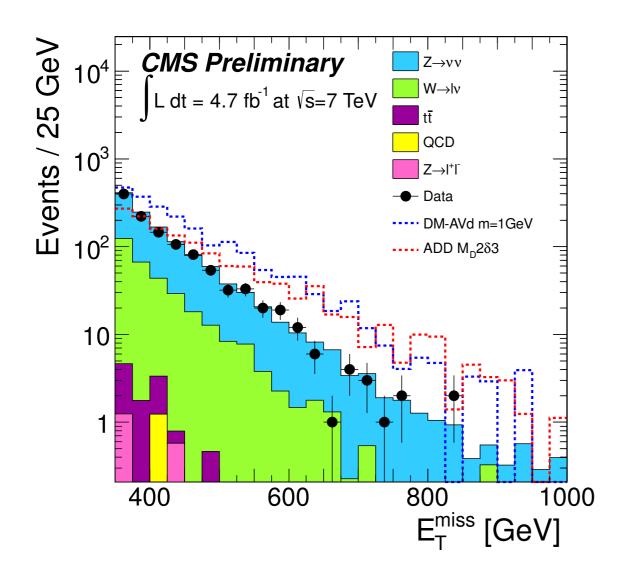
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-Would a yet softer cut yield better bounds?

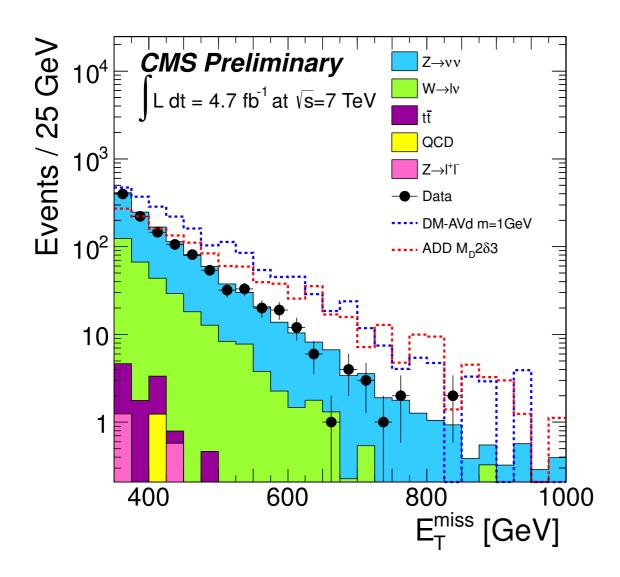
### Latest CMS results



#### EXO-11-059

Background process	Events
$Z \rightarrow \nu \bar{\nu}$	$900\pm94$
W+jets	$312\pm35$
tī	$8\pm8$
$Z(\ell \ell)$ +jets	$2\pm 2$
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Single t	$1\pm 1$
Total background	$1224\pm101$
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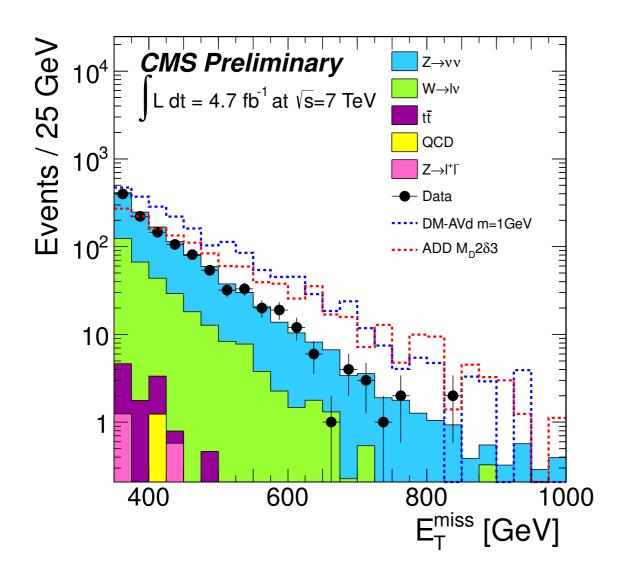


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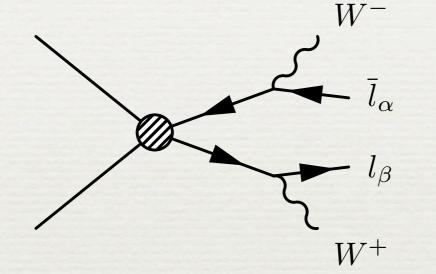
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I/fb ATLAS veryHighPt: $\sigma_{BSM} < 0.015 \text{ pb}$ 4.7/fb CMS: $\sigma_{BSM} < 0.018 \text{ pb}$ 

#### Multileptons vs. monojets

NSI also produce...

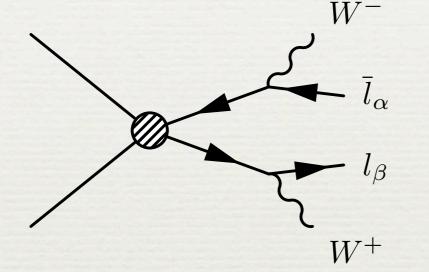


Not as constraining as monojets

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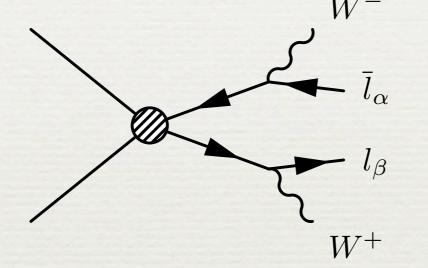
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- Needs very high luminosity (~10 fb<sup>-1</sup>) to compete with monojets.

- Clean lepton final states offer a probe on NSI with different systematics than monojets

#### z=0.3

## Part Two: Dark Matter



Friday, April 13, 2012

### Validity of Effective Description

Suppose by analogy to the NSI case, we wish to constrain DM-quark interactions of the form:

$$\mathcal{O} = \frac{\overline{q}\gamma^{\mu}q\overline{X}\gamma_{\mu}X}{\Lambda^2}$$

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Can we be more rigorous than  $E \sim \Lambda$ ?

Im
$$(a_J)$$
  
Partial-wave unitarity  
Re $(a_J)$   
physical  
 $\mathcal{M} = 16\pi \sum_j (2j+1) P_j (\cos \theta) a_J(s)$ 

Unitarity:  $(\operatorname{Re}(a_J))^2 + (\operatorname{Im}(a_J) - 1/2)^2 \le 1/4$ 

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$$a_J$$
)  
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Unitarity:  $(\operatorname{Re}(a_J))^2 + (\operatorname{Im}(a_J) - 1/2)^2 \le 1/4$   
E.g. Higgs mass (Lee, Quigg, Thacker 1977)  
 $W^+W^- \to W^+W^ m_h \le \sqrt{\frac{8\pi}{5\sqrt{2}G_F}} \approx 780 \text{ GeV}$ 

### Effective dark matter interactions

### Assume heavy particles can be integrated out:

$$\mathcal{O} = \frac{\overline{q}\gamma^{\mu}q \ \overline{X}\gamma^{\mu}X}{\Lambda^2}$$

$$\Rightarrow \mathcal{M}(q\overline{q} \to X\overline{X}) = 2\sqrt{N_c}\frac{s}{\Lambda^2}$$

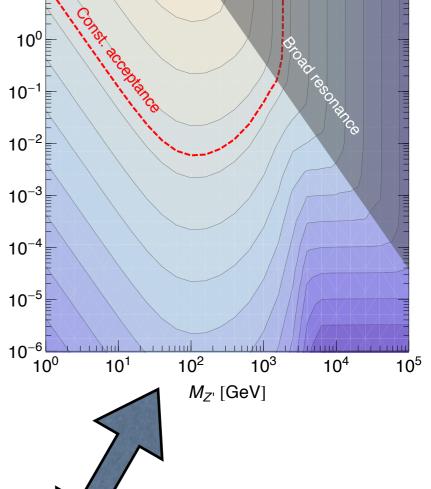
Unitarity implies:  $\Lambda \gtrsim 2 \text{ TeV}$ 

# Effective dark matter $\sqrt{s} = 7 \text{ TeV veryHighPT}$ interactions

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ω

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$$\mathcal{O} = \frac{\overline{q}\gamma^{\mu}q \ X\gamma^{\mu}X}{\Lambda^2}$$
$$\overline{q} \rightarrow X\overline{X} \rightarrow (-1)$$

$$\Rightarrow \mathcal{M}(q\overline{q} \to X\overline{X}) = \begin{cases} \text{Said differently, an effective} \\ \text{DM theory with a 500 GeV} \\ \text{cutoff is consistent if:} \\ s \lesssim 1.3 \text{ TeV} \end{cases}$$

Accessible Z's

With a Z' coupling to DM and a single quark flavor there are 4 parameters in the full parameter space:  $(m_X, m_{Z'}, \sqrt{g_X g_q}, \Gamma_{Z'})$ 

light DM

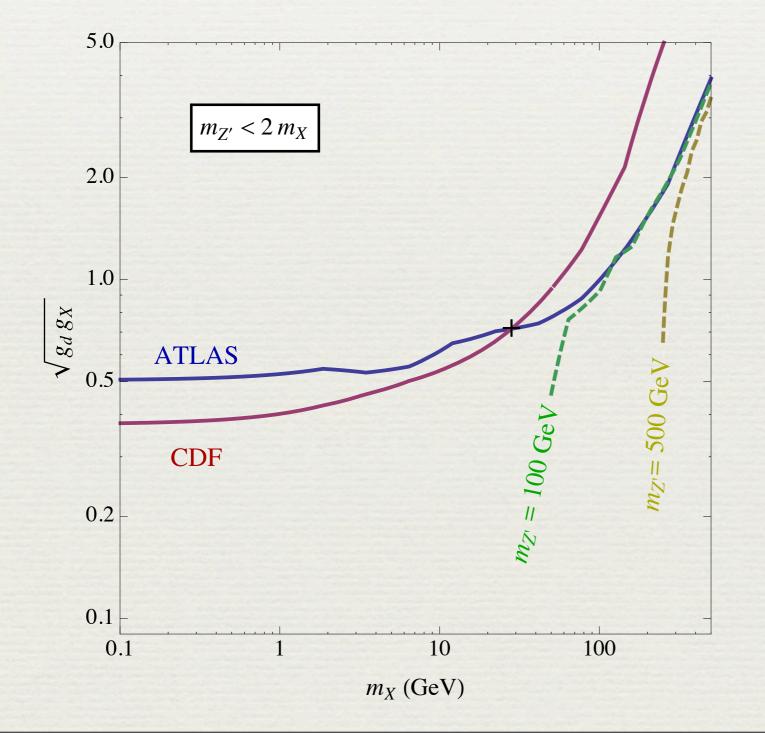
 $\sigma \propto g_a^2 \times \mathrm{BR}\left(Z' \to \overline{X}X\right)$ 

light Z'

q  $Z^{\prime*}$ 

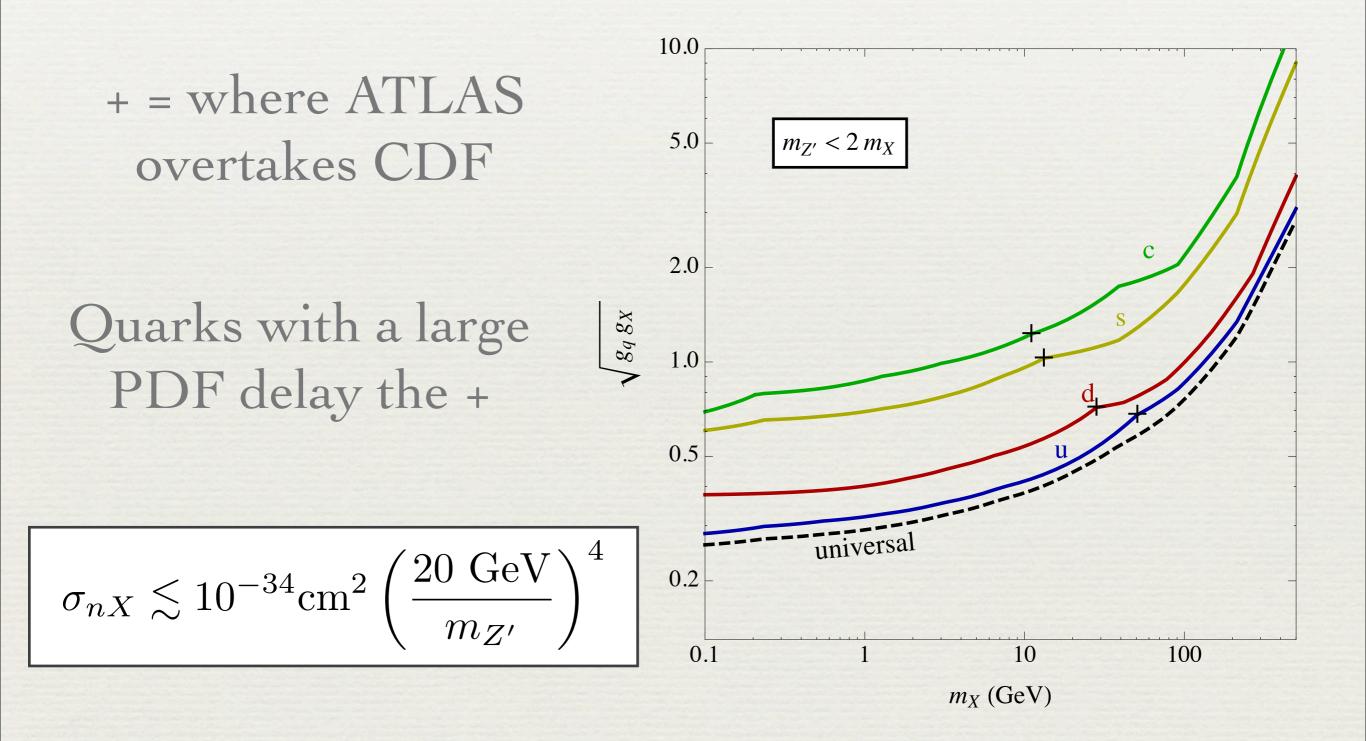
 $\sigma \propto g_q^2 g_X^2$ 

## Tevatron wins even in the off-shell regime

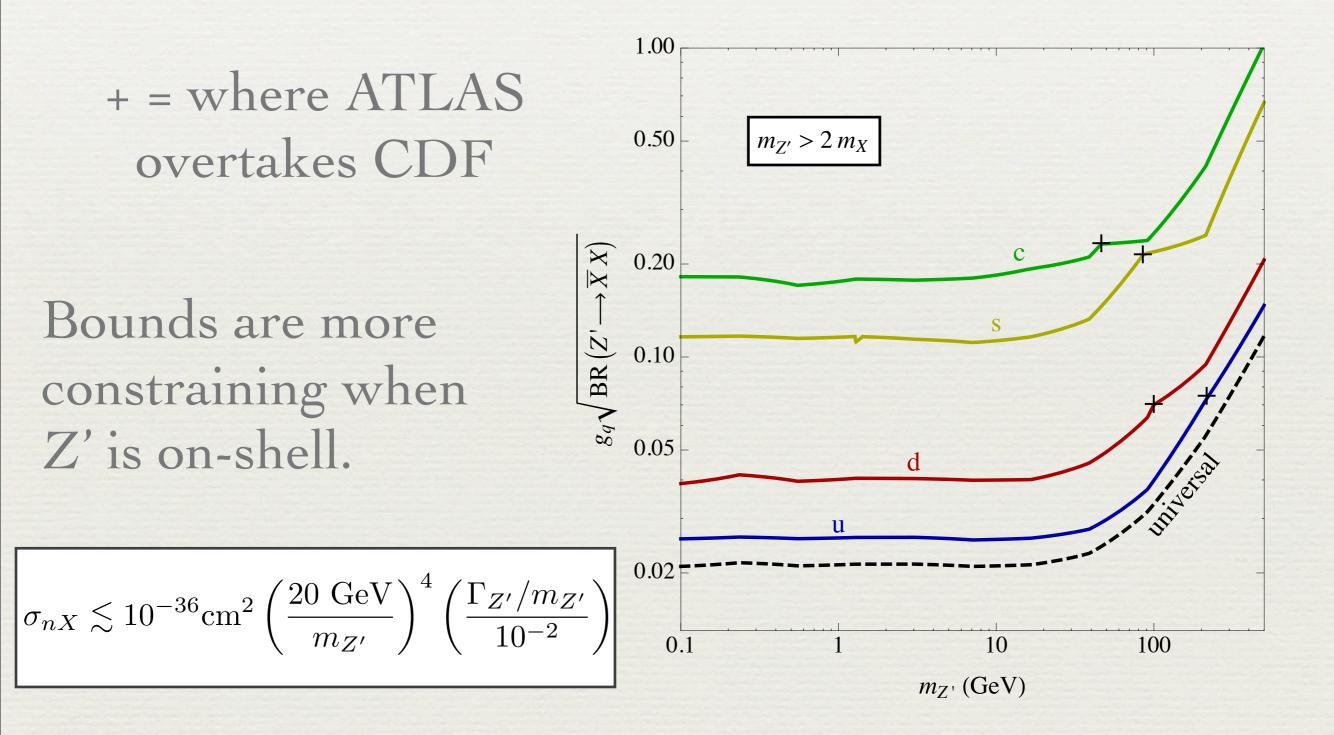


Soft cuts are good for light particles.

#### Upper bounds on quark-DM coupling



#### Quark-DM bounds



Friday, April 13, 2012

#### Take aways

- Neutrinos aren't just a background. Could hide new physics.
- \* Solar data give hints. Test this hypothesis at the LHC.
- The Tevatron 1/fb reigns supreme at low masses.
   Who can beat them?