

LEP Constraints & The Bounds on DM/Higgs Coupling

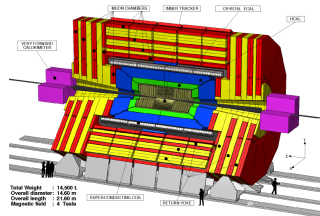
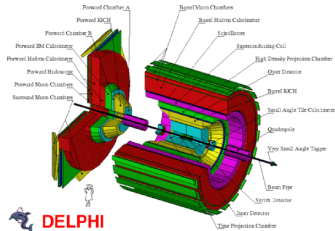
Yuhsin Tsai

Cornell  University

In collaboration with Patrick Fox, Roni Harnik, and Joachim Kopp

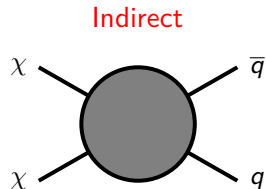
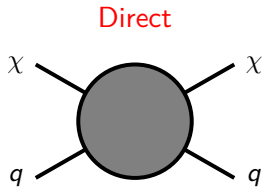
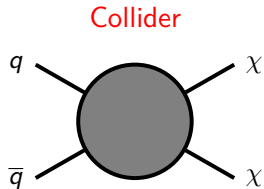
DM in Collision, UC Davis, 12 April 2012

MonoPhoton @ LEP



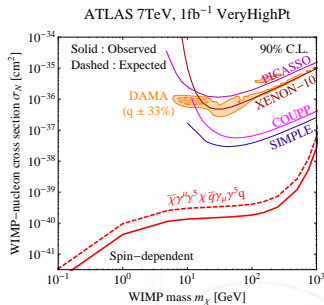
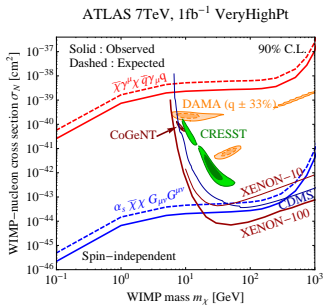
P. J. Fox, R. Harnik, J. Kopp, and YT, 1103.0240

Effective coupling description



- independent of **astrophysical** and **experimental** assumptions.
- good bounds on **light DM**.
- good bounds on **spin dependent** case.

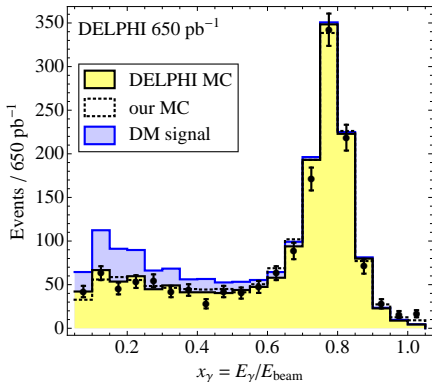
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Mono-photon search at LEP

- In search of Large Extra Dimension (ADD)
- New physics channels: $e\bar{e} \rightarrow \gamma G$

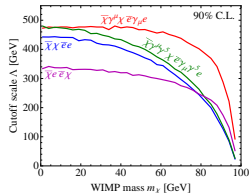
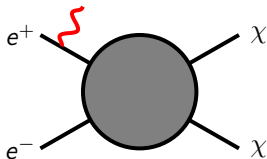


- Experiment: DELPHI
- E_{beam} : 90 – 105 GeV
- Use the cuts in [1],
($E_\gamma \gtrsim 10\text{GeV}$).
- Background: $e^+e^- \rightarrow \gamma\nu\bar{\nu}$
- We use CompHEP.

[1] DELPHI Collaboration, hep-ex/0406019.

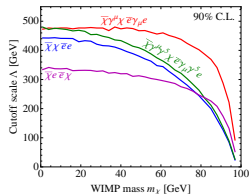
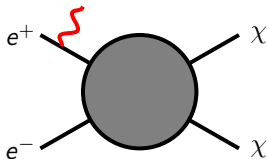
Direct Detection Bound

- Assume DM particle is a **Dirac fermion**.
- Use **shape analysis** (χ^2) to constraint the size the coupling



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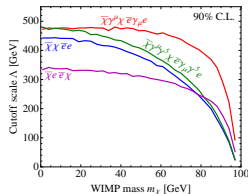
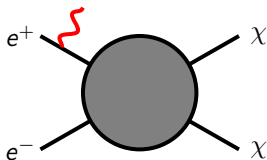
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Need the coupling to quarks!

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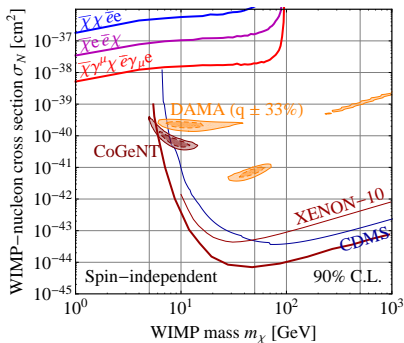
Need the coupling to quarks!

consider two possibilities:

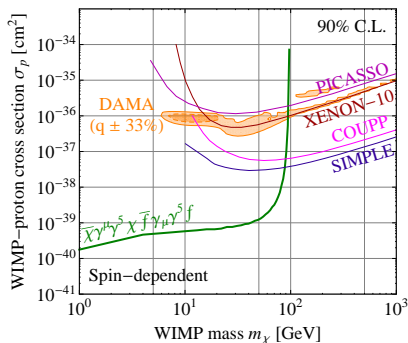
- **Equal Coupling**: to quarks and leptons.
- **Leptophilic**: coupling to leptons only.

Equal couplings

Equal couplings to all SM fermions

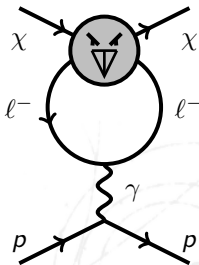
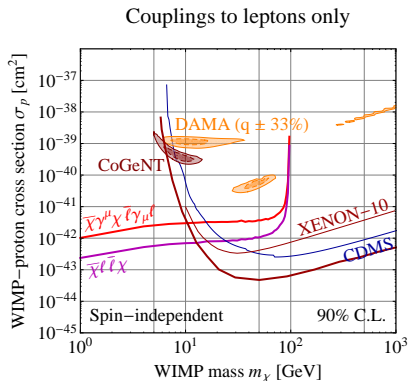


Equal couplings to all SM fermions



Leptophilic: no tree-level coupling to q 's

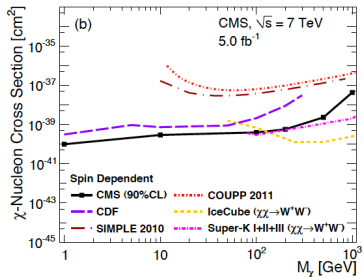
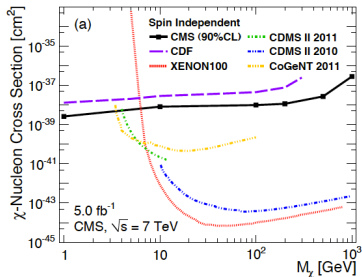
- Get **loop suppression**. \mathcal{O}_A , \mathcal{O}_S vanish at one loop.
- Leptophilic model proposed to explain DAMA or CoGeNT **is ruled out**.



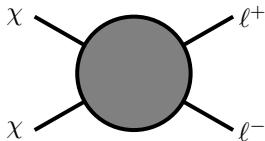
MonoPhoton at LHC

The CMS MonoPhoton search: arXiv:1204.0821

CMS Result: $73 (75.1 \pm 9.4)$



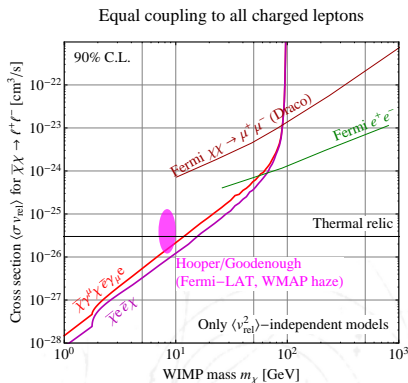
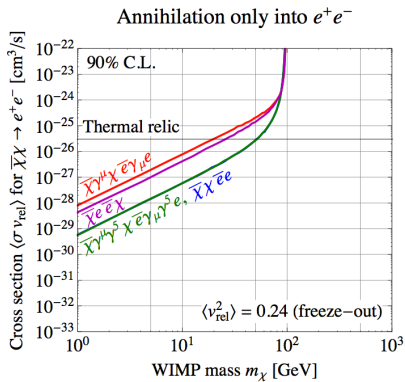
Bounds for Indirect Detections



Two things we can compare to:

- Thermal-relic bound
- Fermi observation bounds.

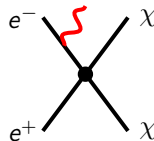
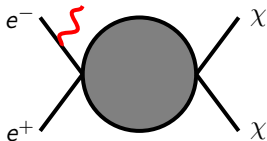
Bounds for Indirect Detections



What happens if the mediator is light?

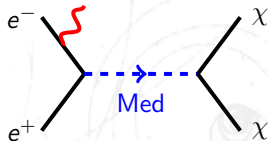
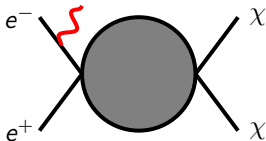
- When it is **heavy**, we consider **contact operators** only:

$$m_{\text{Med}} \gg 2 E_{\text{beam}}$$

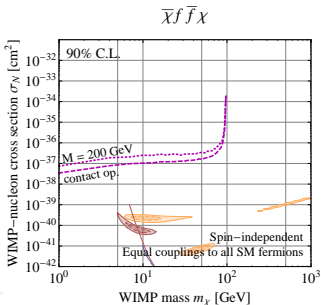
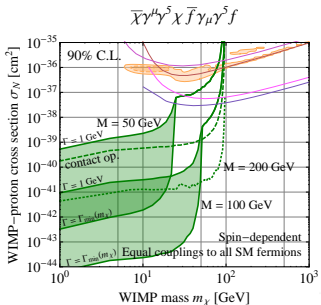
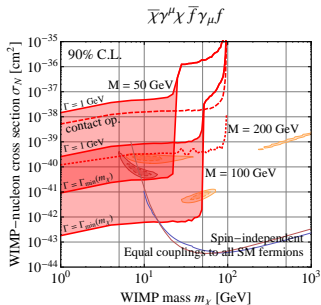
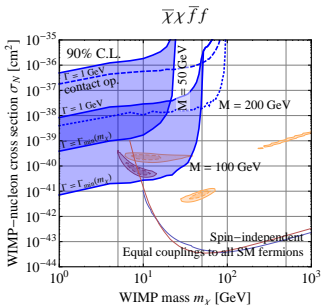


- When it can be **on-shell**, the kinematics is important:

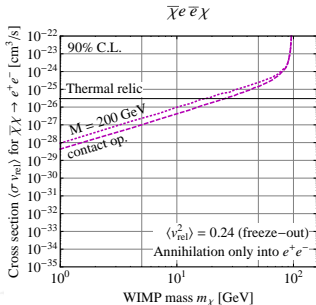
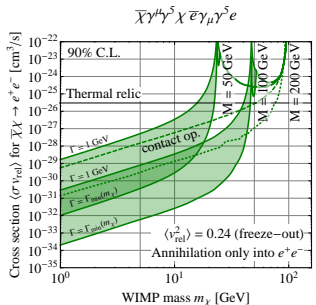
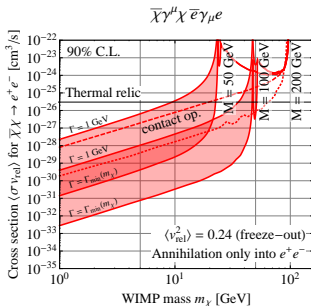
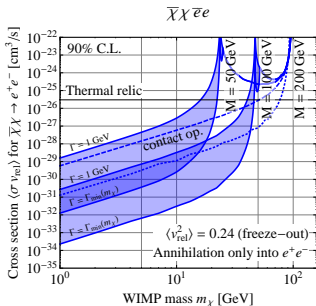
$$m_{\text{Med}} \ll 2 E_{\text{beam}}$$



Direct Detection with light mediator



Indirect Detection with Light Mediators



DM coupling to Higgs

Based on: P. Fox et al. 1109.4398

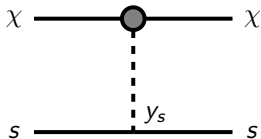
See also:

M. Pospelov et al. 1109.4872; I. Low et al. 1110.4405; C. Englert et al. 1111.1719; O. Lebedev et al. 1111.4482; A. Djouadi et al. 1112.3299.

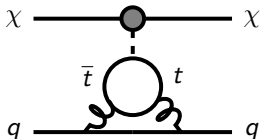
Direct detection through the Higgs portal

Benefits from the additional suppressions

yukawa suppression



loop suppression

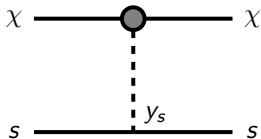


Get the $\chi - h - \chi$ bound from:

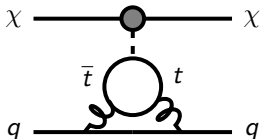
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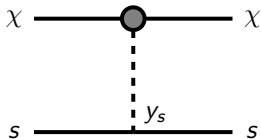
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- mono-jet search? \Rightarrow light mediator, not good.

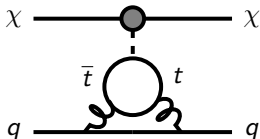
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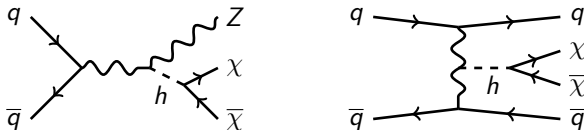
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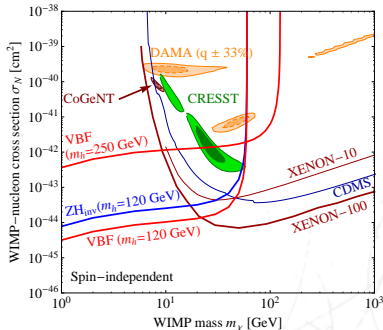
Get the $\chi - h - \chi$ bound from:

- mono-jet search? \Rightarrow light mediator, not good.
- other higgs searches? e.g. invisible Higgs search.

Invisible-higgs search at LHC : Future bounds



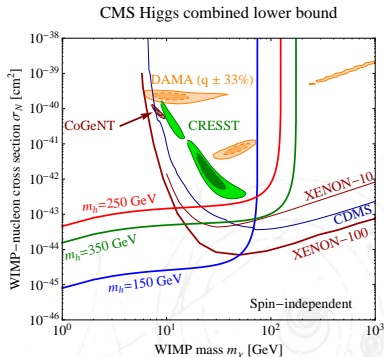
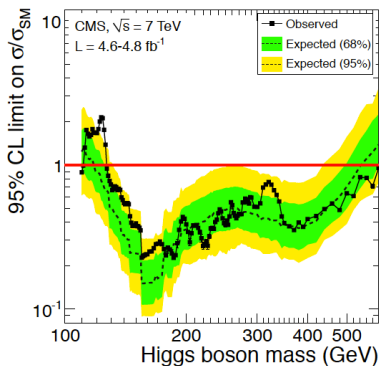
ATLAS 30 fb⁻¹ upper bound (projected)



Invisible-higgs search at LHC : Current bounds

Current LOWER bounds under the assumption:

- Higgs was missed at LHC due to its large invisible width.



Conclusion

Mono-photon search at LEP gives:

- important direct detection bounds
- good constrains on indirect detections
- interesting bounds for light mediators

Direct detection bounds for DM coupling through the Higgs:

- get upper bounds from the invisible Higgs search
- get lower bounds from the current Higgs search

From raw data to direct detection bounds

$$\frac{dR}{dE_R} = N_T \frac{\rho_0}{m_\chi} \int_{v_{\min}}^{v_{\text{esc}}} d^3v \frac{d\sigma}{dE_R} v f(v)$$

- DM density $\rho_0 \sim 0.3 \text{ GeV cm}^{-3}$
- Recoil energy $E_R = E_{\text{obv}} / \text{quenching}$ $q_{\text{Na}} = 0.3 \pm 0.1$,
 $q_{\text{I}} = 0.09 \pm 0.03$
- Velocity distribution $f(v)$ Maxwell-Boltzman
- Escape velocity $v_{\text{esc}} \sim 650 \text{ km s}^{-1}$
- v_{\min} , (in-)elastic scattering?
- Spin Independent σ , Spin Dependent $\sigma(v)$.
- XXXX

Annihilation cross sections

σ_S and σ_A are velocity suppressed:

$$\sigma_{SV_{rel}} = \beta (m_\chi^2 - m_\ell^2) v_{rel}^2,$$

$$\sigma_{VV_{rel}} = \frac{1}{6} \beta \left(24(2m_\chi^2 + m_\ell^2) + \frac{8m_\chi^4 - 4m_\chi^2 m_\ell^2 + 5m_\ell^4}{m_\chi^2 - m_\ell^2} v_{rel}^2 \right),$$

$$\sigma_{AV_{rel}} = \frac{1}{6} \beta \left(24m_\ell^2 + \frac{8m_\chi^4 - 22m_\chi^2 m_\ell^2 + 17m_\ell^4}{m_\chi^2 - m_\ell^2} v_{rel}^2 \right),$$

$$\sigma_{tV_{rel}} = \frac{1}{24} \beta \left(24(m_\chi + m_\ell)^2 + \frac{(m_\chi + m_\ell)^2 (8m_\chi^2 - 16m_\chi m_\ell + 11m_\ell^2)}{m_\chi^2 - m_\ell^2} v_{rel}^2 \right),$$

$$\beta = \frac{1}{8\pi \Lambda^4} \sqrt{1 - \frac{m_\ell^2}{m_\chi^2}}.$$

Few remarks about the loop calculation

(show this if people stay awake)

The **loop-suppressed** cross section is

$$\sigma_{1\text{-loop}} \simeq \frac{4\alpha^2 \mu_p^2}{18^2 \pi^3 \Lambda^4} \cdot \left[\sum_{\ell=e,\mu,\tau} f(q^2, m_\ell) \right]^2$$

where $f(q^2, m_\ell) =$

$$\frac{1}{q^2} \left[5q^2 + 12m_\ell^2 + 6(q^2 + 2m_\ell^2) \sqrt{1 - \frac{4m_\ell^2}{q^2}} \coth^{-1} \left(\sqrt{1 - \frac{4m_\ell^2}{q^2}} \right) - 3q^2 \ln \left(\frac{m_\ell^2}{\Lambda_{\text{ren}}^2} \right) \right]$$

- Take the most conservative case (the largest σ):
 $v_\chi = v_{\text{esc}} = 500 \text{ km/sec}$, scattering angle 180° .
- This gives $q^2 = -4\mu_p^2 v_\chi^2$.
- Take the cutoff Λ_{ren} from the loop integral the same as the operator cutoff Λ .

Direct Detection bounds w/ light mediator

$$\mathcal{A} \propto \frac{g_e g_\chi}{q^2 - M^2 + iM\Gamma} = \frac{M^2}{q^2 - M^2 + iM\Gamma} \frac{g_e g_\chi}{M^2} \equiv (R \times \Lambda)^{-2} = \Lambda_{\text{exp}}^{-2}.$$

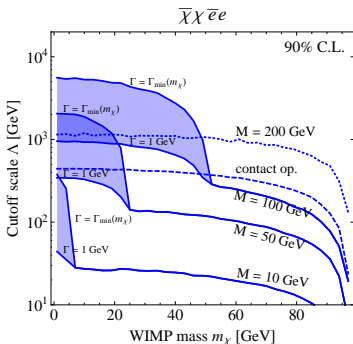
Λ : the cutoff in the plot. Λ_{exp} : the collider constrained cutoff.

S-channel

- $M \gg 2E_{\text{beam}}$: $\Lambda = \Lambda_{\text{exp}}$.
- $M > 2E_{\text{beam}}$: $\Lambda \sim \frac{M}{\sqrt{M^2 - q^2}} \Lambda_{\text{exp}}$.
- $2m_\chi < M < 2E_{\text{beam}}$: $\Lambda \sim \left(\frac{M}{\Gamma}\right)^{\frac{1}{4}} \Lambda_{\text{exp}}$.
- $M < 2m_\chi$: $\Lambda \sim \frac{M}{\sqrt{q^2 - M^2}} \Lambda_{\text{exp}}$.

T-channel

- For any M : $\Lambda = \frac{M}{\sqrt{|q|^2 + M^2}} \Lambda_{\text{exp}}$.



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