# A New Method for Resolving Combinatorial Ambiguities at Hadron Colliders 

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## The LHC Era is here!

- The current 7 TeV run will finish in 2011 with $\approx 1 \mathrm{fb}^{-1}$
- Potential for new physics (cf. Workshop on Topologies for Early LHC Searches, SLAC, Sept. 22-26, 2010)
- One hope is for the direct production of dark matter



## Motivation

- Collider signals of SUSY, UED
- Cascade decay chains ending in LSP, LKP
$-Z_{2}$ symmetry ensures the LSP, LKP is stable
- If neutral, is usually a good dark matter candidate
$-Z_{2}$ symmetry ensures the kinematics of the event are not readily measurable
- Many kinematic methods have been developed to aid determination of the LSP, LKP mass


## Motivation

- Kinematic edges
- Thresholds and maximums of invariant mass distributions provide algebraic expressions relating underlying cascade decay chain masses
- Polynomial method
- Solve momentum conservation equations
- Use non-linear constraints to solve for LSP, LKP momenta
- $\mathrm{m}_{\mathrm{T} 2}, \mathrm{~m}_{\mathrm{CT} 2}$ method
$-\mathrm{m}_{\mathrm{T} 2}$ Kink
- Sub-system $\mathrm{m}_{\mathrm{T} 2}$

Baer, Chen, Paige, Tata, Hinchliffe Allanach, Lester, Parker, White
Barr, Gripaios, Nojiri, Cheng, Gunion, Han, Marandella, McElrath, ...
$-m_{\mathrm{T} 2}$ as a discovery variable

## Motivation

- Need the LSP, LKP mass to determine the mass scale
- Differences in masses are easy
- Usual procedure to solve for LSP, LKP mass:

1. Use cuts to isolate some collection of objects (e.g. 4 jets inclusive +0 leptons + MET), eliminate background
2. Hypothesize an underlying decay chain topology (e.g. pair-produced gluinos decaying to neutralinos via squarks)
3. Assign objects to decay chains (ordered, if possible)
4. Apply your favorite mass reconstruction technique

## Motivation

| Method | Decay chain assignment? <br> (Combinatorial ambiguity) | Ordering required? <br> (Permutative ambiguity) |
| :---: | :---: | :---: |
| $\mathrm{H}_{\mathrm{T}}, M_{\text {eff }}$ | No | No |
| Kinematic edges | Yes | Yes/No |
| Polynomial | Yes | Yes |
| $m_{T 2}$ | Yes | No |
| $m_{C T 2}$ | Yes | No |
| Subsystem $m_{T 2}$ | Yes | Yes |

## Motivation

- Decay chain assignments must deal with combinatorial ambiguities

- For pair-produced gluinos decaying to LSPs via squarks, we have 4 quarks that can be grouped into 3 pair-pair combinations
- Important note: this combinatorial ambiguity is present even if dealing with only signal events


## Motivation

- Other possibilities: in busy cascade decay chains, it is important to associate particles correctly
- Usually, there are additional tools to eliminate wrong combinations



## Outline

- Motivation
- Our simple models
- Brief review of basic kinematics
- The well-studied hemisphere method
- The new $p_{T}$ v. M method
- Comparison
- Conclusions


## Models

- Consider gluino pair-production in a 7 TeV LHC or 14 TeV LHC, both decay identically via on-shell or offshell squarks to LSP neutralino

|  | Gluino Mass | Squark Mass | Neutralino Mass | Kinematic Edge |
| :--- | :--- | :--- | :--- | :--- |
| Model A | 600 GeV | 400 GeV | 100 GeV | 433 GeV |
| Model B | 600 GeV | 800 GeV | 100 GeV | 500 GeV |

- Generate 100,000 events using

MadGraph/MadEvent 4.4.26, decay using BRIDGE

- No ISR/FSR, hadronization
- Only consider parton level


## Kinematic edge in invariant mass

- An oft-used feature of cascade decay chain kinematics is the invariant mass edge
- On-shell squark - characteristic triangular shape

$$
\left.m_{q q}\right|_{\mathrm{edge}}=\sqrt{\frac{\left(m_{\tilde{g}}^{2}-m_{\tilde{q}}^{2}\right)\left(m_{\tilde{q}}^{2}-m_{\hat{\chi}_{1}^{0}}^{2}\right)}{m_{\tilde{q}}^{2}}}
$$

- Off-shell squark - small number of events near edge

$$
\left.m_{q q}\right|_{\text {edge }}=m_{\tilde{g}}-m_{\tilde{\chi}_{1}^{0}}
$$

## Kinematic edge in invariant mass

## On-shell squark scenario



## Kinematic edge in invariant mass

 Off-shell squark scenario

## The Hemisphere Method

- Basic idea: divide an event into hemispheres where each decay chain falls entirely into separate hemispheres
- Two steps
- Step 1: Choose 2 seeds
- These are the central axes for the hemisphere clusters
- Step 2: Cluster remaining objects with the given seeds
- Figure of merit is minimum pdR: roughly, a momentumweighted angular separation

$$
\mathrm{pdR} \equiv(|\Delta \mathrm{p}| \Delta R) \text {, where } \Delta R \equiv \mathrm{~V}\left[(\Delta \phi)^{2}+(\Delta \eta)^{2}\right]
$$

## The Hemisphere Method

- Our implementation

1. Choose highest $p_{T}$ object as seed 1
2. PDR1. Of the remaining objects, choose the maximum pdR object as seed 2

PDR2. Of the remaining objects, choose the maximum invariant mass object as seed 2
3. For each remaining objects, calculate pdR w.r.t seed 1 and w.r.t. seed 2 . Cluster the object with whichever seed has the smaller pdR, i.e. the closer seed in momentum-weighted angle space.

## The Hemisphere Method

## - Our cuts

Cut 1. The highest $p_{T}$ object must have $p_{T} \geq 200 \mathrm{GeV}$
PDR1 Cut 2. The minimum pdR between seed 1 and seed 2 must be 1800 GeV

PDR2 Cut 2. The invariant mass of seed 1 and seed 2 must be larger than the kinematic edge value
Cut 3. Discard all singlet-triplet events
Cut 4. The maximum seed-object invariant mass must be less than or equal to the kinematic edge value

## The Hemisphere Method

- Our cuts do not include realistic detector cuts ( $\eta$ acceptance, minimum $p_{\mathrm{T}}$, isolation requirements)

| PDR 1 Cut <br> Performance | Cut 1 | Cuts 1-2 | Cuts 1-3 | Cuts 1-4 |
| :--- | :--- | :--- | :--- | :--- |
| Model A - 7 TeV | $78.8 \%$ | $25.2 \%$ | $12.4 \%$ | $12.2 \%$ |
| Model A - 14 TeV | $81.7 \%$ | $35.8 \%$ | $18.5 \%$ | $18.2 \%$ |
| Model B-7 TeV | $81.8 \%$ | $27.1 \%$ | $13.4 \%$ | $13.3 \%$ |
| Model B-14 TeV | $83.9 \%$ | $37.5 \%$ | $19.2 \%$ | $18.7 \%$ |
| PDR 2 Cut | Cut 1 | Cuts 1-2 | Cuts 1-3 | Cuts 1-4 |
| Performance |  | $51.4 \%$ | $26.1 \%$ | $25.7 \%$ |
| Model A - 7 TeV | $78.8 \%$ | $58.1 \%$ | $30.5 \%$ | $30.1 \%$ |
| Model A -14 TeV | $81.7 \%$ | $38.5 \%$ | $19.6 \%$ | $19.6 \%$ |
| Model B-7 TeV | $81.8 \%$ | $46.1 \%$ | $24.4 \%$ | $24.4 \%$ |
| Model B-14 TeV | $83.9 \%$ |  |  |  |

## The New $p_{T}$ v. M Method

- Plot $p_{T} \mathrm{v} . \mathrm{M}$ for each qq pair of the event
- Each event has 6 unique qq pairs; 2 pairs are correct, 4 pairs are wrong
- The 6 qq pairs can be grouped into 3 unique pair-pair combinations, one of which is correct
- We want to isolate the correct pair-pair combination
- Observe excesses at high invariant mass (wrong diquark pairs) and high $\mathrm{p}_{\mathrm{T}}$ (correct diquark pairs)


## The $p_{T} v$. M method - Model A 7 TeV <br> 



## The $\mathrm{p}_{\mathrm{T}} \mathrm{v}$. M method - Model A 7 TeV - <br> The $p_{T} v . M$ method - Mod Correct Diquark Pairs Only




## The $\mathrm{p}_{\mathrm{T}} \mathrm{v}$. M method - Model A 7 TeV - <br> The $p_{\mathrm{T}}$ v. M method - Moc Wrong Diquark Pairs Only

## The $p_{\mathrm{T}} v$. M Implementation

- Plot $p_{T} \mathrm{v} . \mathrm{M}$ for each qq pair of the event
- If possible, observe the invariant mass edge
- For the (wrong) pairs with invariant masses larger than the edge value, gradually increase the $p_{T}$ cut such that the survival rate of pairs drops below $5 \%$
- Extrapolate this cut to the upper left region with high $p_{T}$ and low invariant mass
- This region will characteristically have high purity, i.e. be dominated by correct diquark pairs
- Only use pair-pair combinations where both diquark pairs lie in the boxed region


## The $p_{T} v$. M method - Model A 7 TeV <br> 

## Survival Probability

| Cut on min $p_{\mathbf{T}}$ Survival $\%$ |  |
| ---: | ---: | ---: |
| 0 | $100.000 \%$ |
| 25 | $98.525 \%$ |
| 50 | $94.402 \%$ |
| 75 | $88.347 \%$ |
| 100 | $80.857 \%$ |
| 125 | $72.162 \%$ |
| 150 | $62.948 \%$ |
| 175 | $53.796 \%$ |
| 200 | $44.963 \%$ |
| 225 | $36.548 \%$ |
| 250 | $29.014 \%$ |
| 275 | $22.481 \%$ |
| 300 | $16.960 \%$ |
| 325 | $12.398 \%$ |
| 350 | $8.982 \%$ |
| 375 | $6.275 \%$ |
| 400 | $4.378 \%$ |
| 425 | $3.007 \%$ |
| 450 | $2.077 \%$ |
| 475 | $1.427 \%$ |
| 500 | $0.994 \%$ |



## The $p_{T} v$. M method - Model B 7 TeV  <br>  <br> Shading reflects the fraction of diquark pairs at each ( $\left.\mathrm{m}_{\mathrm{qq}}, \mathrm{p}_{\mathrm{T}(\mathrm{qq})}\right)$ point that are correct

## The $p_{T}$ v. M method - Model B 14 TeV



## Comparison between the hemisphere

 method and the $p_{T} v . M$ method- Will use event efficiency v. event sample purity
- Event efficiency is the percentage of events that pass cuts
- Event purity is the percentage of remaining events that are fully correctly assigned
- Variable cut
- Hemisphere method: Vary the object-seed pdR difference
- Example: For a given object, its pdR w.r.t seed 1 is 500 GeV , and its pdR w.r.t seed 2 is 505 GeV . By the pdR measure, it could equally well be clustered with seed 1 or 2 . By imposing an increasing pdR difference cut (a minimum difference in pdR between an object and the two seeds), we can gradually eliminate these ambiguous assignment scenarios.
$-p_{T} v . M$ method: Vary the survival probability in the (wrong) diquark pairs region at high invariant mass
- Equivalently, require a higher $p_{T}$ cut


## Model A-7 TeV and 14 TeV - Event Efficiency v. Event Purity




Model B-7 TeV and 14 TeV - Event Efficiency v. Event Purity



## Discussion

- On-shell decay chains
$-p_{T} v . M$ is significantly better than the hemisphere method in retaining more event efficiency for a given event purity
- Off-shell decay chains
$-p_{T} v . M$ is marginally to moderately worse than the hemisphere method
- Possibly because of the flexibility in choosing second hemisphere axis


## Discussion

- $p_{\mathrm{T}} \mathrm{v}$. M is more flexible
- No distinction between choosing seeds and clustering
- If seed 1 and seed 2 are incorrect, the hemisphere method fails
- Requires strict cuts to ensure seed 1 and seed 2 are from different decay chains
- $p_{\mathrm{T}} \mathrm{v} . \mathrm{M}$ and the hemisphere method are readily generalized
- Multi-jets, leptons, complicated decay chains including W and Z bosons


## Future Work

- Apply $p_{T} v . M$ at reconstruction level
- SPS1a including ISR/FSR, detector simulation
- Perform shape analysis of $p_{T} v . M$
- Optimize the $\mathrm{p}_{\mathrm{T}} \mathrm{v}$. M cut
- Reorganize $p_{T} v . M$ to be an event-by-event variable
- Perform a detailed study of $p_{\mathrm{T}} \mathrm{v} . \mathrm{M}$ and the hemisphere method in off-shell cases
- Should use both in parallel since it is not known a priori whether the decay chain is on-shell or off-shell


## Conclusions

- Distinguishing combinatorial ambiguities is important for new physics searches at the LHC
- The $p_{\mathrm{T}} \mathrm{v}$. M method is better than the hemisphere method for on-shell decay chains in delivering high purity event samples
- The hemisphere method is better suited for off-shell decay chains
- The $p_{T} v . M$ method is easy to implement and flexible
- A "robustness" study in a simulated collider environment is underway


## Kinematic edge in invariant mass On-shell squark scenario




## Kinematic edge in invariant mass Off-shell squark scenario




## The Hemisphere Method - Cut 1

## $\mathrm{p}_{\mathrm{T}}($ of seed 1$) \geq 200 \mathrm{GeV}$



## The Hemisphere Method - PDR 1 Cut 2

## pdR (of seed 1 and seed 2 ) $\geq 1800 \mathrm{GeV}$

Seed 1 and seed 2 on opposite decay chains


## The Hemisphere Method - PDR 2 Cut 2

## $M$ (of seed 1 and seed 2 ) $>\mathrm{M}_{\text {edge }}$



## The Hemisphere Method - PDR 1 Cut 4

## $M$ (of seed and object) $\leq M_{\text {edge }}$

Seed and object on the same decay chain



## The Hemisphere Method - PDR 2 Cut 4

## M (of seed and object) $\leq \mathrm{M}_{\text {edge }}$

Seed and object on the same decay chain



## Alternative cut: maximum dR cut

- Constraining dR differences does not work



