The Cosmology of Split SUSY

Aaron Pierce (SLAC) February, 2005

> Discussions with S. Dimopoulos, N. Arkani-Hamed, J. Wacker

The Proving Ground

- LHC Start 2007
- Should Discover source of Electroweak
 Symmetry
 Breaking—Higgs
 Boson
- But what Else?



Naturalness

- A full blown industry—twenty years of thinking what lies at the weak scale and beyond.
- Technicolor, Supersymmetry, Large Extra Dimensions, Warped Compactifications, Little Higgs
- None are perfect. Model Builders are busy.

TeV Scale Supersymmetry

 $e \otimes e \text{ (selectron)}$ $q \otimes \tilde{q} \text{ (squark)}$ $\gamma \otimes \tilde{\gamma} \text{ (photino)}$ $h \otimes \tilde{h} \text{ (higgsino)}$

Weak Scale Supersymmetry Scorecard

SUCCESSES

- Gauge Coupling Unification
- Dark Matter (χ)

FAILURES

- Where are the superpartners?
- Higgs Mass
- Flavor Changing Neutral Currents (b→s γ), (μ→e γ), K-K mixing...
- Electric Dipole Moments
- Dimension 5 proton decay
- Comsological Problems

 (e.g. gravitino problem)

Cosmological Constant

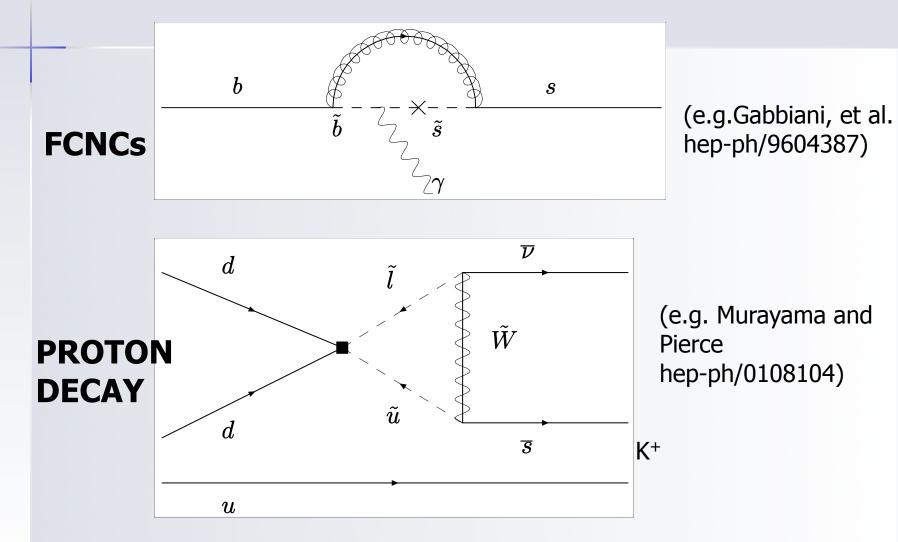
Higgs Fine tune

Stringy Landscape

- The existence of many vacua may liberate us from traditional ideas of naturalness.
- Weinberg C.C. Bound
- Friendly neighborhood?

(Arkani-Hamed, Dimopoulos, Kachru hep-th/0501082)

The Trouble with Scalars

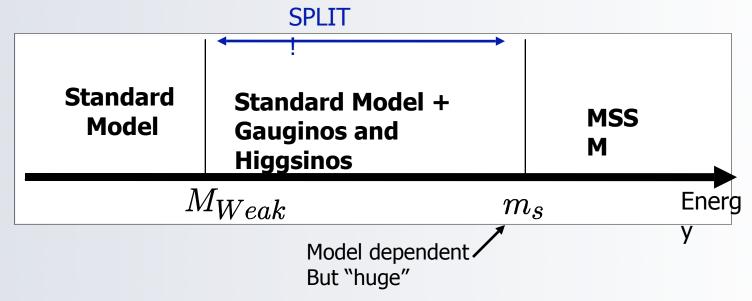


Gauge Coupling Unification

- Scalars (and SM fermions) come in complete GUT multiplets→ don't contribute to relative running
- Higgsinos, Gauginos, Gauge Bosons, Higgs Bosons are responsible for unification.

Take Home Message

- Fermions Good, Scalars Bad
- Fermion masses can be protected by a chiral symmetry. Can be much lighter than scalars.



Why Weak Scale?

Lost the link to the Higgs Mass... Superpartners could be anywhere?

■ Existence of Dark Matter → Weak Scale. (Lee-Weinberg)

Split Sosrestfecard

SUCCESSES

- Gauge Coupling Unification
- Dark Matter (χ)

PREDICTIONS

- Superpartners→Dark Matter
- Higgs Mass < 170 GeV
- Electric Dipole Moment near bound

(Arkani-Hamed, Dimopoulos, Giudice, Romanino)

FAILURES

- Where are the superpartners?
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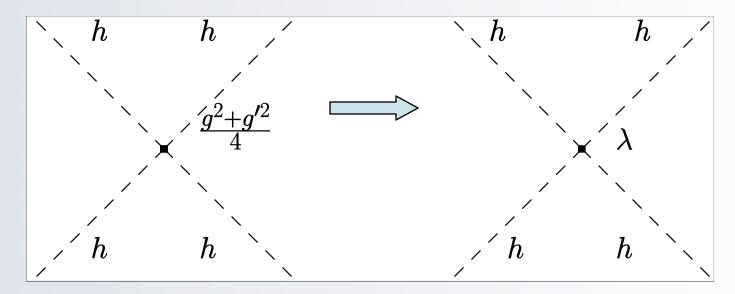
Higgs Fine tune

The Remaining Plan

Dark Matter and Split SupersymmetryGluino Cosmology

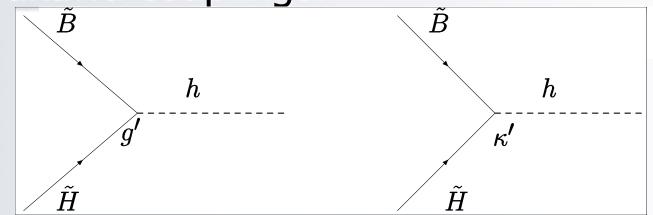
Novel Features

- Only one Higgs doublet at low energies.
- Supersymmetric boundary conditions are set at m_s



New RGEs for couplings

- Run Higgs Quartic → Predicts the Higgs Mass.
- Gaugino Couplings are generalized Yukawa couplings.



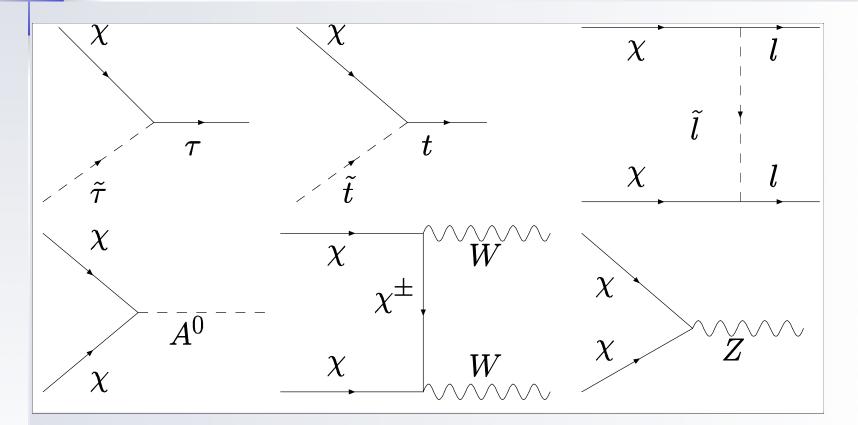
Gauge Couplings? Yukawas.

$$\kappa_1' = \sqrt{\frac{3}{10}} g_1 \sin \beta \qquad \kappa_2' = \sqrt{\frac{3}{10}} g_1 \cos \beta$$

$$\kappa_1 = \sqrt{2} g_2 \sin \beta \qquad \kappa_2 = \sqrt{2} g_2 \cos \beta.$$

$$M_{\chi^{0}} = \begin{pmatrix} M_{1} & 0 & -\frac{\kappa_{2}'v}{\sqrt{2}} & \frac{\kappa_{1}'v}{\sqrt{2}} \\ 0 & M_{2} & \frac{\kappa_{2}v}{\sqrt{8}} & -\frac{\kappa_{1}v}{\sqrt{8}} \\ -\frac{\kappa_{2}'v}{\sqrt{2}} & \frac{\kappa_{2}v}{\sqrt{8}} & 0 & -\mu \\ \frac{\kappa_{1}'v}{\sqrt{2}} & -\frac{\kappa_{1}v}{\sqrt{8}} & -\mu & 0 \end{pmatrix}, \qquad M_{\chi^{\pm}} = \begin{pmatrix} M_{2} & \frac{v\kappa_{1}}{2} \\ \frac{v\kappa_{2}}{2} & \mu \end{pmatrix}$$

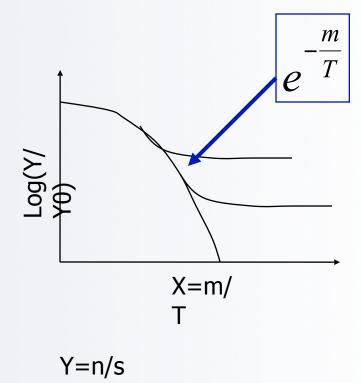
MSSM: Many Particles, Many Diagrams



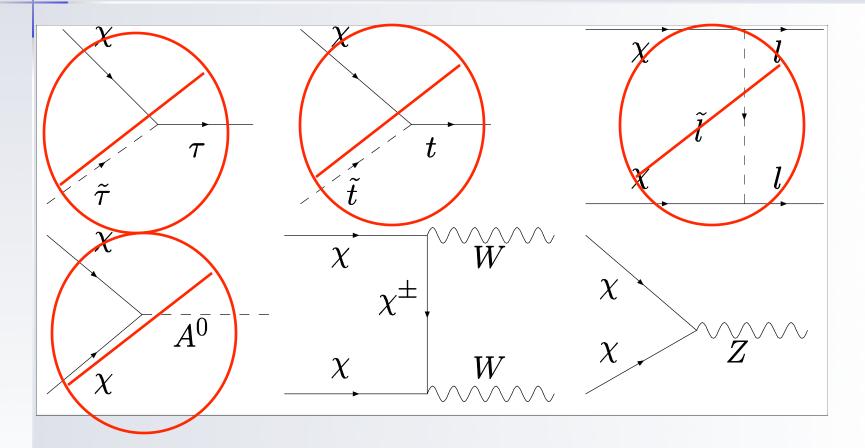
Dark Matter Calculation

Solve the Boltzman equation:

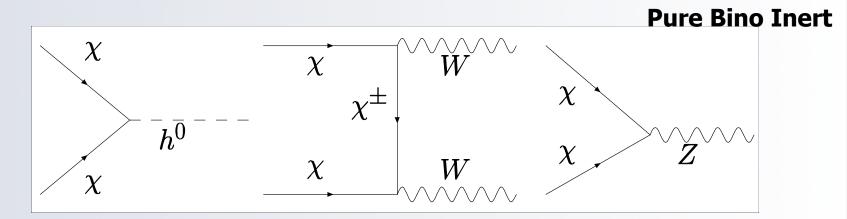
$$\frac{dn}{dt} + 3Hn = -\langle \sigma_A | v \rangle [n^2 - n_{EQ}^2]$$



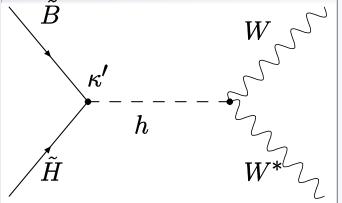
SPLIT SUSY FEWER MSSM: Many Particles, FEWE<mark>Many Diagrams</mark>



Split SUSY: Fewer Particles, Fewer Diagrams



■ Heavier Higgs → Wider Higgs → Easier Resonance $\setminus \tilde{B}$



Split SUSY: Three regions

- Pure Wino (AMSB-like)
- Bulk Region" (mixed Bino—Wino)
- Pure Higgsino (radiative generation of mu term)

$$16\pi^{2}\beta_{\mu} = -\mu(\frac{9}{10}g_{1}^{2} + \frac{9}{2}g_{2}^{2}) \\ +\frac{3}{2}\kappa_{1}\kappa_{2}M_{2} + 2\kappa_{1}'\kappa_{2}'M_{1} \\ +\frac{3}{8}\mu(\kappa_{1}^{2} + \kappa_{2}^{2}) + \frac{1}{2}\mu(\kappa_{1}'^{2} + \kappa_{2}'^{2})$$

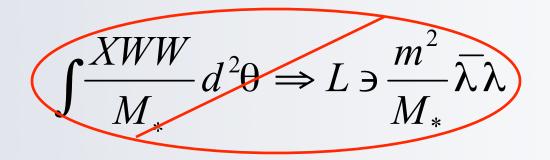
Pure Wino Case

 Naturally Arises from Anomaly Mediation.

(Randall and Sundrum; **Giudice, Luty, Murayama, Rattazzi; Wells**).

$$\int \frac{X^* X Q^* Q d^4 \theta}{M_*^2} \Rightarrow L \Rightarrow \frac{m^4}{M_*^2} q q^*$$

 $X = m^2 \theta^2$



Anomaly Mediation

Superconformal anomaly can communicate supersymmetry breaking.

$$m_{\lambda} = \frac{\beta m_{3/2}}{g}$$

$$12: M1: M3 = 1:3:7$$

$$M2 = 2.2 \text{ TeV}$$

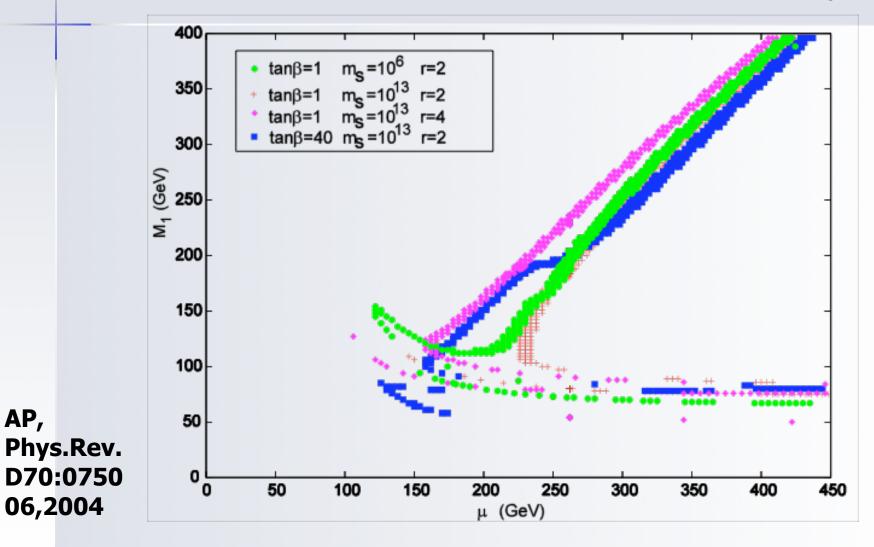
Μ

$$\int \frac{X^* X Q^* Q d^4 \theta}{M_*^2} \Rightarrow \quad 100 \text{ TeV Scalar Masses}$$

"Bulk" Allowed Region r =

$$r = M_2 / M_1$$

 m_s SUSY breaking scale



Mu problem?

Crucial that mu~M1. (NMSSM?) $W = \lambda XY^2 + \kappa XH_U H_D \Rightarrow$ $V \supset \lambda^2 |Y|^4$ $m_{\chi}^2 > 0, m_{\chi}^2 < 0 \Rightarrow$ $< Y > \sim m_{S}, < X > = 0.$

Direct Detection of Dark Matter

Look for WIMP colliding with a nucleus.

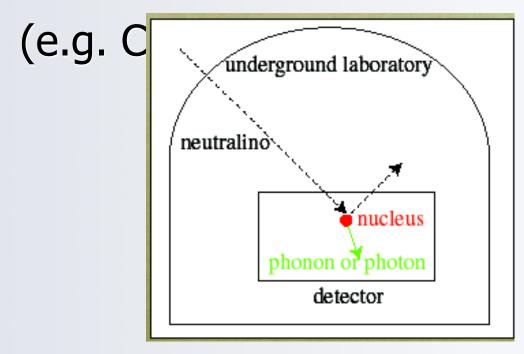
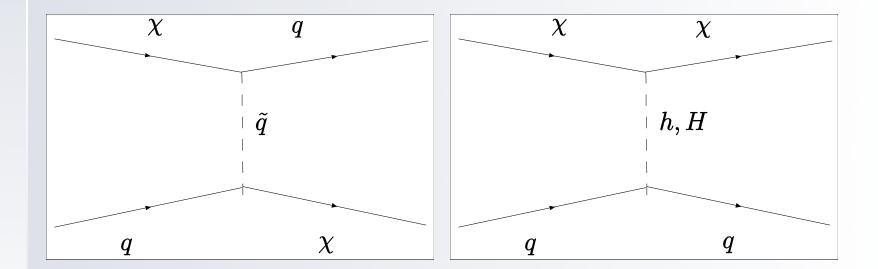
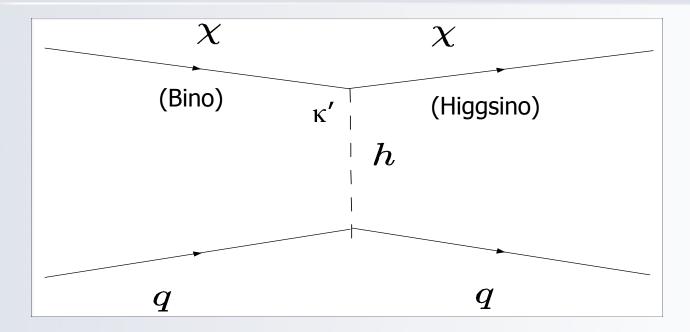


Figure from H. Murayama

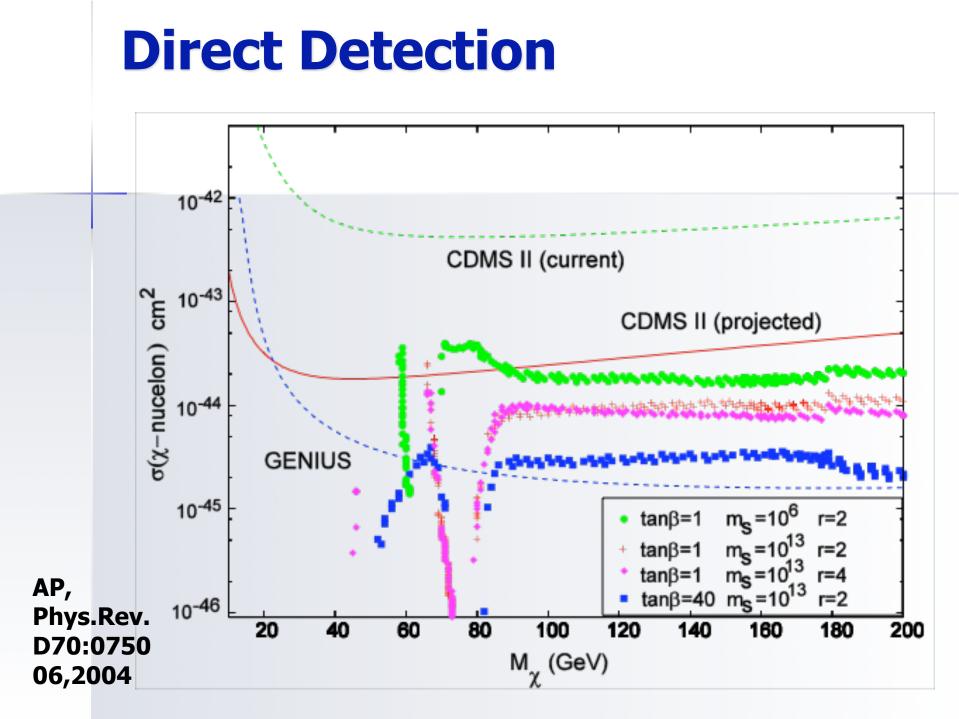
Direct Detection (MSSM)



Direct Detection Split SUSY



- Mixed Dark Matter
- Higgs Boson is somewhat heavier
- Use Modified DarkSUSY



Gluino Cosmology

- Heavy Squarks → Long Lived Gluino
- A Possible Collider Signal
- A Possible constraint...
 - Exotic Isotopes (heavy hydrogen)
 - BBN
 - CMBR/Gamma Rays

Ongoing discussions with S. Dimopoulos, J. Wacker, Graham, Arvanitaki, Davis

Constraints from BBN

- Radiative Decays Ellis et al. (1992)
- Hadronic Decays Dimopoulos et al. (1988)

Scatter off of electrons/photons or **nuclei.**

Compute the Relic Abundance

Again, solve the Boltzman equation. Compute freeze-out temperature.

$$\ln\left\{\frac{\langle\sigma^{\operatorname{ann}}v\rangle}{4\pi^{3}}\sqrt{\frac{45}{2g^{*}(T_{F})G_{N}}}m_{\widetilde{g}}g_{\widetilde{g}}x_{F}^{-1/2}\right\} = x_{F} \qquad \qquad \mathsf{X=m/}$$

$$\langle \sigma^{\rm ann} v \rangle = \frac{1}{8m_{\widetilde{g}}^4 T K_2^2(m_{\widetilde{g}}/T)} \int_{4m_{\widetilde{g}}^2}^{\infty} \sigma^{\rm ann}(s) s^{3/2} \beta^2 K_1(\sqrt{s}/T) ds$$

$$\frac{1}{Y_0} - \frac{1}{Y_F} = \left[\frac{45G_N}{\pi}\right]^{-1/2} \int_{x_F}^{x_0} \frac{h^*(T)}{\sqrt{g^*(T)}} \frac{m_{\widetilde{g}}}{x^2} \langle \sigma^{\rm ann} v \rangle dx \;, \qquad {\sf Y} = {\sf n/s}$$

See Gondolo and Gelmini (1991)

How does annihilation proceed?

Perturbative Annihilation

Baer, Cheung and Gunion (hep-ph/9806361)

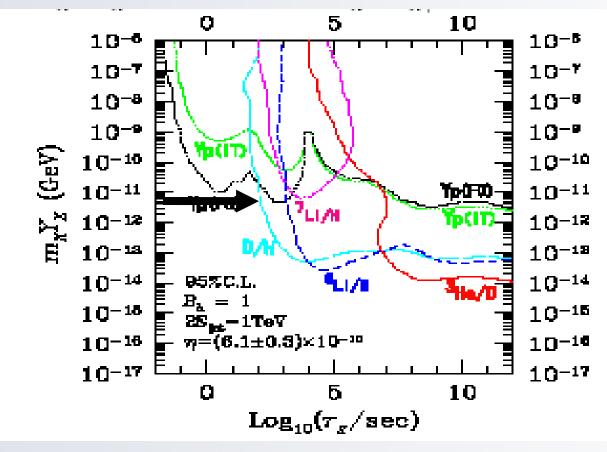
$$\begin{split} &\sigma(\widetilde{g}\widetilde{g} \to gg) = \frac{3\pi\alpha_s^2}{16\beta^2 s} \left\{ \log \frac{1+\beta}{1-\beta} \left[21 - 6\beta^2 - 3\beta^4 \right] - 33\beta + 17\beta^3 \right\} \\ &\sigma(\widetilde{g}\widetilde{g} \to q\overline{q}) = \frac{\pi\alpha_s^2\overline{\beta}}{16\beta s} (3-\beta^2)(3-\overline{\beta}^2) \,. \end{split}$$

Sommerfeld Enhancement

- (due to low velocity)

Gluinonia? R-Hadrons?

BBN Constraints



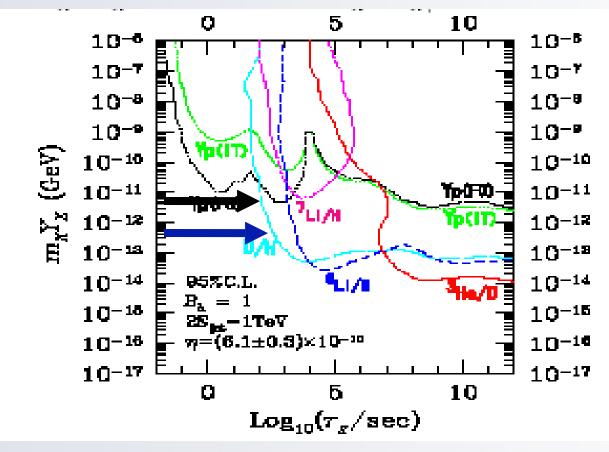
Kawasaki, Kahri and Moroi (hep-ph/0402490)

Additional Annihilation

- Extreme possibility \rightarrow gluinos annihilate with $\sigma \approx 1 / \Lambda_{QCD}^2$
- This evades all BBN constrains by ~factor 30.
- More likely: use "heavy gluino effective theory". Assume saturates s-wave unitarity.

$$\sigma \approx 4\pi / k^2 \approx 4\pi / (m_g^2 \beta^2)$$

BBN Constraints



Kawasaki, Kahri and Moroi (hep-ph/0402490)

Conclusions

- Dark Matter provides the link to the weak scale in the split story.
- Split SUSY Dark Matter is simpler
- Mixed case accessible at Direct Detection and LHC→ confirm DM!
- Tells us something about MSSM without resonances and coannihilations.

Gluino Lifetime

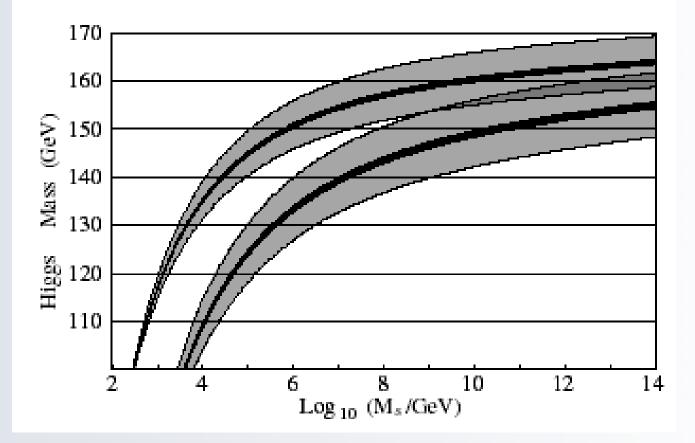
$$\tau \approx 8 \left(\frac{m_S}{10^9 \text{ GeV}} \right)^4 \left(\frac{1 \text{ TeV}}{m_{\tilde{g}}} \frac{1}{j} \text{ seconds} \right)^5$$

Indirect Detection

Dependent on model of galactic halo. Can see gamma annihilation? Some optimism for pure wino region.

> Arvanitaki and Graham hep-ph/0411376 Masiero et al. hep-ph/0412058

Higgs Mass



Wacker, et al. (2004)

Other Possibilities for Dark Matter

- Gravitino Decay could provide additional source of Dark Matter.
- \rightarrow lighter DM (all other things equal)
- Depends on reheat temp, and mass of gravitino. (Arkani-Hamed, Dimopoulos, Giudice, Romanino)